Regular time-reproductive solutions for generalized Boussinesq model with Neumann boundary conditions for temperature

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- 1. Statement of the problem. The main result
- 2. The Galerkin Initial-Boundary Problem
- 3. Differential Inequalities in regular norms
- 4. Proof of theorem
- 5. Some comments and open problems

## 1. Statement of the problem

 $\Omega \subset \mathbb{R}^{N} \text{ regular bounded domain, } N = 2, 3, T > 0.$   $\begin{cases} \partial_{t} \boldsymbol{u} - \nabla \cdot (\boldsymbol{\nu}(\theta) \nabla \boldsymbol{u}) + (\boldsymbol{u} \cdot \nabla) \boldsymbol{u} - \alpha \boldsymbol{g} \theta + \nabla \boldsymbol{p} = \boldsymbol{f}, \\ \nabla \cdot \boldsymbol{u} = 0, \\ \partial_{t} \theta - \nabla \cdot (\boldsymbol{k}(\theta) \nabla \theta) + (\boldsymbol{u} \cdot \nabla) \theta = 0, \end{cases}$ (1)

in  $\Omega \times [0,\infty)$ 

- $\boldsymbol{u}(x,t) \in \mathbb{R}^N$  velocity of the fluid at point  $x \in \Omega$  and time  $t \in [0,T)$
- $p(x,t) \in \mathbb{R}$  (hydrostatic) pressure.  $\theta(x,t) \in \mathbb{R}$  temperature.
- $g(x,t) \in \mathbb{R}^N$  gravitational field  $f(x,t) \in \mathbb{R}^N$  resulting of external forces.
- $\alpha > 0$  constant associated to the coefficient of volume expansion.
- $\nu(\cdot) : \mathbb{R} \to \mathbb{R}$  kinematic viscosity,  $k(\cdot) : \mathbb{R} \to \mathbb{R}$  thermal conductivity.

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in  $\Omega \times [0,\infty)$ 

• Dirichlet-Neumann boundary conditions:

$$\boldsymbol{u} = 0, \qquad \partial_n \boldsymbol{\theta} = \boldsymbol{0} \qquad \text{on } [0, \infty) \times \partial \Omega,$$

• Time reproductive condition:

$$\boldsymbol{u}(0) = \boldsymbol{u}(T), \qquad \boldsymbol{\theta}(0) = \boldsymbol{\theta}(T) \qquad \text{in } \Omega.$$

## **Previous results:** Existence and uniqueness of initial value problem with Dirichlet's boundary conditions

S.A. Lorca, J.L. Boldrini, *Stationary solutions for generalized Boussinesq models*, J. Diferential Equations 124 (2), (1996).

**Goal:** To obtain strong estimates:  $H^2$  for the velocity,  $H^3$  for the temperature and consequently, existence of regular reproductive solution

**Difficulties:** Time reproductive condition Nonlinear diffusion

## Let us consider the following spaces

$$H_N^k = \left\{ \theta \in H^k; \ \frac{\partial \theta}{\partial n} = 0 \text{ on } \partial \Omega, \int_{\Omega} \theta = 0 \right\}$$

where k = 2, 3.

 $|\Delta \theta|_2 \approx ||\theta||_2 \text{ in } H_N^2$  $|\nabla \Delta \theta|_2 \approx ||\theta||_3 \text{ in } H_N^3$ 

## Definition:

 $(u, p, \theta)$  is a regular solution in (0, T), if  $u \in L^2(H^2) \cap L^{\infty}(H^1)$  and  $\partial_t u \in L^2(L^2)$ ,  $p \in L^2(H^1)$ ,  $\theta \in L^2(H^3_N) \cap L^{\infty}(H^2_N)$  and  $\partial_t \theta \in L^2(H^1_N)$ ,

satisfying (1) a.e. in  $\Omega \times (0, T)$ , boundary conditions and time reproductivity conditions in the sense of spaces V and  $H_N^2$  respectively.

### Theorem

Let T > 0,  $\Omega \subset \mathbb{R}^N$  a bounded domain (N = 2 or 3),  $\partial \Omega \in C^{2,1}$ .

- $\nu \in C^1(\mathbb{R}), \quad 0 < \nu_{min} \le \nu(s) \le \nu_{max}, \quad |\nu'(s)| \le \nu'_{max}$
- $k \in C^2(\mathbb{R}), \quad 0 < k_{min} \le k(s) \le k_{max}, \quad |k'(s)| \le k'_{max}, \quad |k''(s)| \le k''_{max}.$
- $\boldsymbol{f} \in L^2(\boldsymbol{L}^2), \quad \boldsymbol{g} \in L^\infty(\boldsymbol{L}^2), \quad \|\boldsymbol{f}\|_{L^2(L^2)} \leq \delta, \quad \delta \text{ small enough}$

then there exists a regular (and small) reproductive solution of (1) in (0,T). Moreover, this solution also verifies  $\partial_t \theta(0) = \partial_t \theta(T)$ . Let  $\{\phi_i\}_{i\geq 1}$  and  $\{\varphi_i\}_{i\geq 1}$  "special" basis of V and  $H_0^1(\Omega)$ , respectively, formed by eigenfunctions of the Stokes and the Poisson problems following:

$$\begin{cases} -\Delta \phi_i = \lambda_i \phi_i & \text{in } \Omega \\ \phi_i = 0 & \text{on } \partial \Omega \end{cases} \begin{cases} -\Delta \varphi_i = \mu_i \varphi_i & \text{in } \Omega \\ \partial_n \varphi_i = 0 & \text{on } \partial \Omega \end{cases}$$
$$\|\phi_i\|_1 = 1, \|\varphi_i\|_1 = 1 \text{ for all } i \text{ and } \int_{\Omega} \varphi_i = 0. \end{cases}$$

$$\boldsymbol{u}_m(t) = \sum_{j=1}^m \xi_{i,m}(t)\phi_i \qquad \qquad \boldsymbol{\theta}_m(t) = \sum_{j=1}^m \zeta_{i,m}(t)\varphi_i,$$

For each  $m \ge 1$ , given  $u_{0m} \in V^m$  and  $\theta_{0m} \in W^m$ , there exits a unique solution  $(u_m, \theta_m)$ , with  $u_m : [0, T] \mapsto V^m$  and  $\theta_m : [0, T] \mapsto W^m$ , verifying the following variational formulation a.e. in  $t \in (0, T)$ :

$$\begin{cases} (\partial_t \boldsymbol{u}_m(t), \boldsymbol{v}_m) + ((\boldsymbol{u}_m(t) \cdot \nabla) \boldsymbol{u}_m(t), \boldsymbol{v}_m) + (\boldsymbol{\nu}(\theta_m(t)) \nabla \boldsymbol{u}_m(t), \nabla \boldsymbol{v}_m) \\ -(\alpha \theta_m(t) \boldsymbol{g}, \boldsymbol{v}_m) - (\boldsymbol{f}, \boldsymbol{v}_m) = 0 \qquad \forall \, \boldsymbol{v}_m \in \boldsymbol{V}^m \\ (\partial_t \theta_m(t), e_m) + ((\boldsymbol{u}_m(t) \cdot \nabla) \theta_m(t), e_m) \\ + (\boldsymbol{k}(\theta_m(t)) \nabla \theta_m(t), \nabla e_m) = 0 \qquad \forall \, e_m \in W^m \\ \boldsymbol{u}_m(0) = \boldsymbol{u}_{0m}, \quad \theta_m(0) = \theta_{0m}, \end{cases}$$

 $(\boldsymbol{u}_m \text{system}, A \boldsymbol{u}_m) + (\boldsymbol{u}_m \text{system}, \partial_t \boldsymbol{u}_m)$ 

 $\int_{\Omega} \frac{d}{dt} \int_{\Omega} (\boldsymbol{\nu}(\theta_m) + 1) |\nabla \boldsymbol{u}_m|^2 + \nu_{min} ||\boldsymbol{u}_m||_2^2 + |\partial_t \boldsymbol{u}_m||_2^2 \le \delta ||\partial_t \theta_m||_1^2$  $+ \varepsilon ||\boldsymbol{u}_m||_2^2 ||\theta_m||_2 + K(||\boldsymbol{u}_m||_1^6 + ||\boldsymbol{u}_m||_1^2 ||\theta_m||_2^4 + ||\boldsymbol{g}||_{L^{\infty}(L^2)}^2 ||\theta_m||_2^2 + |\boldsymbol{f}|_2^2)$  $for <math>\delta, \varepsilon > 0$  small enough, and  $K = K(\delta, \varepsilon) > 0$ 

From viscosity term

## 3. Differential inequalities in regular norms

 $(\partial_t(\theta_m \text{eq.}), \partial_t \theta_m) + (\theta_m \text{eq.}, \Delta^2 \theta_m)$ 

# $\frac{d}{dt} (\|\theta_m\|_2^2 + |\partial_t \theta_m|_2^2) + k_{min} (\|\theta_m\|_3^2 + \|\partial_t \theta_m\|_1^2)$ $\leq \delta |\partial_t \mathbf{u}_m|_2^2 + C_\delta (\|\theta_m\|_2^6 + \|\theta_m\|_2^4 |\partial_t \theta_m|_2^2 + \|\theta_m\|_2^2 \|\mathbf{u}_m\|_1^4)$

for  $\delta > 0$  small enough, and  $C_{\delta} > 0$ 

- $\Delta^2 \theta_m \in W^m$  thanks to the election of spectral basis
- Integrating by parts in all terms, boundary terms vanish since

 $(\nabla \Delta \theta_m \cdot \boldsymbol{n})|_{\partial \Omega} = 0$ 

One obtains:

 $-(\partial_t \nabla \theta_m, \nabla \Delta \theta_m) + (\nabla [\nabla \cdot (\mathbf{k}(\theta_m) \nabla \theta_m)], \nabla \Delta \theta_m) - (\nabla (\mathbf{u} \cdot \nabla \theta_m), \nabla \Delta \theta_m) = 0.$ 



$$\frac{d}{dt} \int_{\Omega} (\boldsymbol{\nu}(\theta_m) + 1) |\nabla \boldsymbol{u}_m|^2 + \nu_{min} \|\boldsymbol{u}_m\|_2^2 + |\partial_t \boldsymbol{u}_m|_2^2 \le \delta \|\partial_t \theta_m\|_1^2 + \varepsilon \|\boldsymbol{u}_m\|_2^2 \|\theta_m\|_2 + K(\|\boldsymbol{u}_m\|_1^6 + \|\boldsymbol{u}_m\|_1^2 \|\theta_m\|_2^4 + \|\boldsymbol{g}\|_{L^{\infty}(L^2)}^2 \|\theta_m\|_2^2 + |\boldsymbol{f}|_2^2)$$

+

$$\frac{d}{dt} (\|\theta_m\|_2^2 + |\partial_t \theta_m|_2^2) + k_{min} (\|\theta_m\|_3^2 + \|\partial_t \theta_m\|_1^2) 
\leq \delta |\partial_t u_m|_2^2 + C_\delta (\|\theta_m\|_2^6 + \|\theta_m\|_2^4 |\partial_t \theta_m|_2^2 + \|\theta_m\|_2^2 \|u_m\|_1^4)$$

Adequate balance 
$$\Rightarrow ||\mathbf{g}||_{L^{\infty}(L^2)}^2 ||\theta_m||_2^2$$

We denote:

$$\Phi_m(t) = \int_{\Omega} (\boldsymbol{\nu}(\theta_m) + 1) |\nabla \boldsymbol{u}_m|^2 + ||\theta_m||_2^2 + |\partial_t \theta_m|_2^2$$
$$\Psi_m(t) = ||\boldsymbol{u}_m||_2^2 + |\partial_t \boldsymbol{u}_m|_2^2 + ||\theta_m||_3^2 + ||\partial_t \theta_m||_1^2$$

We obtain:

$$\begin{cases} \Phi'_m + C\Psi_m \leq \varepsilon \Psi_m \Phi_m^{1/2} + C_0(t) + D\Phi_m^3 \\ \Phi_m(0) = \Phi_{m0} \\ C_0(t) = C_0 |\mathbf{f}|_2^2 \end{cases}$$

First step: If  $\Phi_m(0) \leq \delta$  and  $\|\mathbf{f}\|_{L^2(L^2)} \leq \delta$ , then  $\Phi_m(t) < 2\delta \ \forall t \in [0, T]$ .

Absurd argument:  $T^*$  such that  $\Phi_m(T^*) = 2\delta \quad \Phi_m(s) < 2\delta \quad \forall s \in [0, T^*).$ 

+

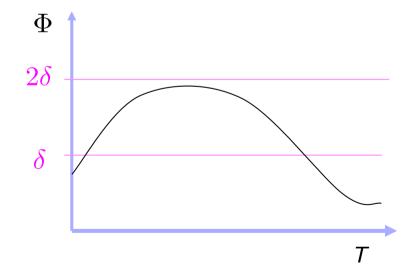
Poincaré inequality:  $C_p > 0$  such that  $\Phi_m(t) \leq C_p \Psi_m(t)$  ( $\varepsilon$  small enough)

$$\Rightarrow \Phi'_m + C\Psi_m \le \varepsilon \Psi_m \Phi_m^{1/2} + C_0(t) + D\Phi_m^3$$
$$\Phi'_m + \tilde{C}\Phi_m \le C_0(t) + D\Phi_m^3$$

Integrating in  $[0, T^*]$ :  $\Phi_m(T^*) \le \delta e^{-\bar{C}T^*} + \int_0^{T^*} C_0(t) < 2\delta. \longrightarrow \longleftarrow$ 

Second step: If  $\Phi_m(0)$  and  $\|\mathbf{f}\|_{L^2(L^2)}$  are small enough, then  $\Phi_m(T) < \Phi_m(0)$ . Similarly, integrating in [0, T],

$$\Phi_m(T) \le \Phi_m(0)e^{-\bar{C}T} + \int_0^T C_0(s).$$



Third step: Existence of approximate reproductive solution.

Given  $(\boldsymbol{u}_{m0}, \theta_{m0}) \in V^m \times W^m$ ,

$$L^{m}: [0,T] \mapsto \mathbb{R}^{m} \times \mathbb{R}^{m}$$
$$t \mapsto (\xi_{1m}(t), ..., \xi_{mm}(t), \zeta_{1m}(t), ..., \zeta_{mm}(t))$$

 $(\xi_{1m}(t),...,\xi_{mm}(t)), (\zeta_{1m}(t),...,\zeta_{mm}(t))$  coefficients of  $\boldsymbol{u}_m(t)$  and  $\theta_m(t)$ 

Given 
$$L_0^m = L^m(0)$$
  
 $\bar{B} = \{(\xi_{1m}, ..., \xi_{mm}, \zeta_{1m}, ..., \zeta_{mm}) := L_0^m : \Phi_m(0) \le \delta\}.$   
 $\mathcal{R}^m : \bar{B} \subset \mathbb{R}^m \times \mathbb{R}^m \quad \mapsto \quad \mathbb{R}^m \times \mathbb{R}^m$   
 $L_0^m \quad \mapsto \quad \mathcal{R}^m(L_0^m) = L^m(T)$ 

Brouwer Theorem.

Four step: Pass to the limit in reproductive approximate solutions

$$\Phi_m(t) = \int_{\Omega} (\nu(\theta_m) + 1) |\nabla u_m|^2 + ||\theta_m||_2^2 + |\partial_t \theta_m|_2^2 \le 2\delta$$

(independent of m) for small data.

 $(\boldsymbol{u}_m)$  uniformly bounded in  $L^{\infty}(H^1) \cap L^2(H^2)$ ,  $(\theta_m)$  uniformly bounded in  $L^{\infty}(H_N^2) \cap L^2(H_N^3)$ ,  $(\partial_t \boldsymbol{u}_m)$  uniformly bounded in  $L^2(L^2)$ ,  $(\partial_t \theta_m)$  uniformly bounded in  $L^{\infty}(L^2) \cap L^2(H^1)$ .

and

$$(\boldsymbol{u}_m)$$
 is relatively compact in  $L^2(H^1)$   
 $(\theta_m)$  is relatively compact in  $L^2(H^2)$ .

Sufficient to pass to the limit in equations.

 $\theta_m$  is relatively compact in  $C([0,T]; H^1)$ 

$$\begin{array}{c} \theta_m(T) \longrightarrow \theta(T) \text{ in } H^1(\Omega) \\ \\ \\ \theta_m(0) \longrightarrow \theta(0) \text{ in } H^1(\Omega) \end{array} \end{array}$$

 $\theta_m(T)$  and  $\theta_m(0)$  are bounded in  $H^2(\Omega)$ 

$$\Rightarrow \quad \theta(T) = \theta(0) \text{ in } H^2(\Omega)$$

 $\partial_{tt}\theta_m$  is uniformly bounded in  $L^2((H^1)')$ 

 $\partial_t \theta_m$  is uniformly bounded in  $L^{\infty}(L^2)$ 

 $\Rightarrow \partial_t \theta_m$  is relatively compact in  $C([0,T]; (H^1)')$ 

$$\partial_t \theta_m(T) \longrightarrow \partial_t \theta(T) \text{ in } (H^1)'(\Omega)$$

$$\stackrel{\parallel}{\partial_t \theta_m(0)} \longrightarrow \partial_t \theta(0) \text{ in } (H^1)'(\Omega)$$

$$\Rightarrow \partial_t \theta(0) = \partial_t \theta(T)$$

- Dirichlet boundary condition is imposed for the temperature  $\theta$ , the boundary terms do not vanish in the integration by parts.
- The uniqueness remains open. Higher regularity for the velocity is necessary.  $H^3$ -regularity for u and Dirichlet condition ?