

## Setup for the measurement of the $^{235}\text{U}(n,f)$ cross section relative to n-p scattering up to 1 GeV

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**Abstract.** The neutron induced fission of  $^{235}\text{U}$  is extensively used as a reference for neutron fluence measurements in various applications, ranging from the investigation of the biological effectiveness of high energy neutrons, to the measurement of high energy neutron cross sections of relevance for accelerator driven nuclear systems. Despite its widespread use, no data exist on neutron induced fission of  $^{235}\text{U}$  above 200 MeV. The neutron facility n\_TOF offers the possibility to improve the situation. The measurement of  $^{235}\text{U}(n,f)$  relative to the differential n-p scattering cross-section, was carried out in September 2018 with the aim of providing accurate and precise cross section data in the energy range from 10 MeV up to 1 GeV. In such measurements, Recoil Proton Telescopes (RPTs) are used to measure the neutron flux while the fission events are detected and counted with dedicated detectors. In this paper the measurement campaign and the experimental set-up are illustrated.

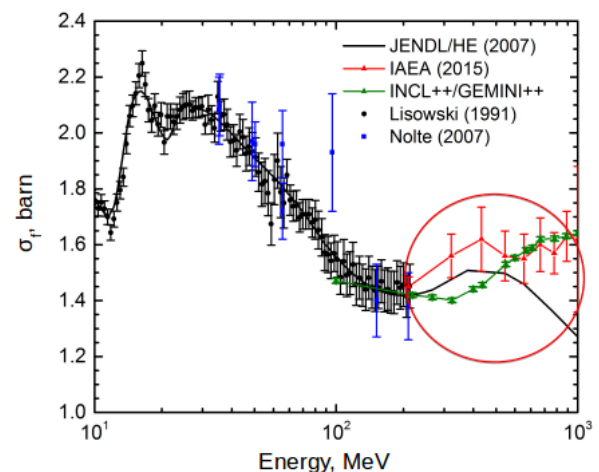
## 1 Introduction

Neutron cross section standards are fundamental ingredients for both measurements and evaluations of neutron-induced reaction cross sections. However, no cross section standard exists for neutron energies above 200 MeV. This led the International Atomic Energy Agency to issue a request for a new absolute measurement of neutron induced fission cross sections, relative to n-p scattering, to establish a fission cross section standard above 200 MeV [1]. An effective choice for the reaction to be studied is the  $^{235}\text{U}(n,f)$  reaction, already one of the most important standard cross section at thermal neutron energy and between 0.15 MeV and 200 MeV [2]. Despite its importance in fundamental nuclear physics and its widespread use as reference for neutron fluence measurements, only two data sets are available in the energy range between 20 and 200 MeV and two absolute experimental measurements have been performed in the high energy region, above 200 MeV [3, 4] (figure 1). Moreover the fission process of Uranium at high excitation energy is an important topic as it is related to fundamental quantities of excited nuclear matter, like the viscosity and transient effects [10, 11].

Thanks to the intense neutron beam, with a wide energy spectrum ranging from thermal to 1 GeV, the n\_TOF facility at CERN can answer to the request of the IAEA for new measurements. A dedicated measurement campaign was carried out with the aim of providing accurate and precise cross section data of the  $^{235}\text{U}(n,f)$  reaction in the high energy region.

## 2 Experimental set-up

Neutrons at n\_TOF (see [13] for an extensive description) are produced by spallation reaction: bunches of protons of 20 GeV/c momentum are shot every few seconds by the



**Figure 1.** The  $^{235}\text{U}(n,f)$  cross section from the JENDL/HE [5] and IAEA [1] evaluations, the experimental data [6, 7] measured relative to the n-p cross section and a new theoretical calculation [8] based on the intranuclear cascade model INCL++ [9] coupled to the deexcitation model GEMINI++ [12].

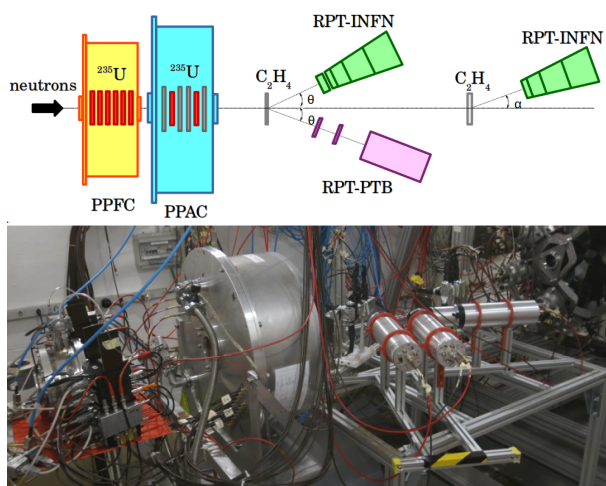
CERN Proton Synchrotron on a massive lead target; about 300 neutrons are produced per incident proton. The facility features two beam lines with a white neutrons spectrum, and two experimental areas at the nominal distance of 185 m for EAR-1 [14] and 20 m for EAR-2 [15] from the neutron-producing target.

The measurement of the  $^{235}\text{U}(n,f)$  cross section was performed in September 2018 in the first experimental area, where the neutron spectrum extends up to 1 GeV, and also to profit from the good energy resolution in the high energy region, thanks to the longer flight path.

The experimental setup included detectors to count the fission events from the Uranium target and, at the same time,

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detectors to measure the neutron fluence relative to the neutron-proton elastic scattering reaction, the primary reference for neutron cross sections. Fission events were detected using an ionisation chamber from the Physikalisch-Technische Bundesanstalt [16] and a reaction chamber based on Parallel Plate Avalanche Counters (PPACs) developed at Institut de Physique Nucléaire d'Orsay [17, 18]. Three recoil proton telescopes (RPTs), developed by the National Institute of Nuclear Physics (INFN, Italy) and the Physikalisch-Technische Bundesanstalt (PTB, Germany), are used to measure the neutron fluence. Figure 2 shows, the scheme of the experimental set up: the two fission chambers with a total of ten fissile deposits placed in the neutron beam, downstream two Polyethylene ( $C_2H_4$ ) targets and the three telescopes placed at a small angle with respect to the neutron beam.



**Figure 2.** Layout of the  $^{235}\text{U}(n,f)$  relative to  $^1\text{H}(n,n)^1\text{H}$  cross section measurement set-up carried out in EAR-1 at n\_TOF. Fission reactions are measured with a fission chamber and a set of Parallel Plate Avalanche Counters. Three telescopes pointing at two different  $C_2H_4$  radiators were used to measure the neutron fluence.

## 2.1 Fission reaction detection

The Parallel Plate Avalanche Counters consist of 3 electrodes, a central anode surrounded by two cathodes; the 3 mm gaps between electrodes are filled with forced flow of  $C_3F_8$  maintained at low-pressure (4 mbar). The electrodes are made of  $1.7 \mu\text{m}$  thick mylar foils, coated with aluminium or gold to make them conductive. The coating in the cathodes is divided into 2 mm wide strips to allow the localisation of the avalanche by using a delay line; the time difference between delay line outputs provides a one-dimensional position. By combining the signals from the two orthogonal cathodes strips, the fission fragment trajectory can be reconstructed. The fission reaction chamber includes 3 PPACs with two  $^{235}\text{U}$  samples in between. Each sample consisted in a 8 cm diameter,  $300 \mu\text{g} \cdot \text{cm}^{-2}$  thick layer of  $^{235}\text{U}$  (purity of 92.7%) that was placed on a  $600 \mu\text{g} \cdot \text{cm}^{-2}$  aluminium backing by electrodeposition.

The fission events are identified by coincident anode signals of two consecutive PPACs. This feature, combined with the fast timing properties of the device (9 ns FWHM), guarantees a very high signal-to-background ratio. In addition, since these detectors are quite insensitive to the  $\gamma$ -flash produced by high energy reactions at the spallation target, it is possible to recognize the events produced by 1 GeV neutrons. However, the PPACs can detect only fission fragments emitted into a forward cone with an opening angle of about  $60^\circ$  with a detection efficiency that is not easily evaluated.

The second system used to measure the fission reaction is an ionization chamber (PPFC) which contains a parallel plate electrode assembly with a thickness of  $100 \mu\text{m}$  each. The chamber is equipped with  $^{235}\text{U}$  (purity of 99.93%) deposits of 42 mm diameter and about  $300 \mu\text{g} \cdot \text{cm}^{-2}$  areal density on both sides of four aluminum electrodes. In total 32.660(3) mg of  $^{235}\text{U}$  are used. In addition, to study the background reactions due to the aluminium, there are two targets with only the backing. The chamber is operated at atmospheric pressure with a continuous gas flow of 90% argon and 10% methane mixture.

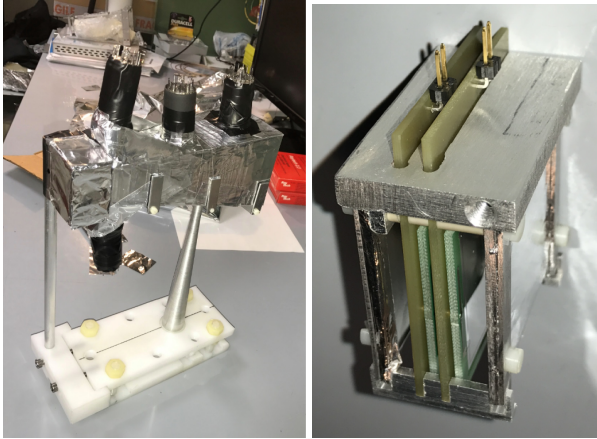
The detector is very well characterized in terms of detection uniformity of the fissile deposits and of fragment detection efficiency (about 95%). The use of the PPFC is however limited to up to a few hundreds MeV of neutron energy because of the background produced by neutron-induced reactions on the aluminium foils. It is perfectly suited for the measurement of the  $^{235}\text{U}(n,f)$  cross section in the energy range from thermal to ca. 150 MeV; therefore, it is used to study and calibrate the PPAC fragment-detection efficiency in the energy range below 100 MeV.

## 2.2 Flux measurement set-up

The number of neutrons impinging on the  $^{235}\text{U}$  samples is measured simultaneously to fission by detecting recoil protons emitted from the two polyethylene targets placed downstream of the fission chambers (figure 2). Due to the presence of carbon in the samples, it is necessary to use a telescope to discriminate the protons from other charged particles exploiting the  $\Delta E$ -E method. The energy deposited in the different layers of the telescope is different for different particles, due to the different stopping power. In non-relativistic approximation, the product between the energy loss in the transmission detectors ( $\Delta E$ ) and the remaining kinetic energy ( $E$ ) in the stop detector is proportional to the charge ( $z$ ) and the mass ( $M$ ) of the interacting particle. Graphing  $\Delta E$  against  $E$  yield a family of hyperbolas corresponding to the different values of  $z^2M$ . Two RPTs were built with the same design, the third follows a different concept.

The two detectors, indicated in figure 2 as RPT-INFN (shown in detail in figure 3) have a compact design. The two telescopes consist of a trapezoidal structure of four independent BC408 plastic scintillators with increasing thickness: 0.5, 3.0, 6.0 and 6.0 cm. Each scintillator is coupled to a 1" Hamamatsu R1924A Photomultiplier (PMT), except for the first scintillator, that, due to the low thickness, is coupled to two PMTs (one at each side). One





**Figure 3.** Left: structure of one of the two RPT-INFN: four plastic scintillators and two Silicon detectors. Right: detail of the case with the Silicon detectors inside.

of the two RPT-INFN has two silicon detectors placed before the first thin scintillator. The silicon detectors,  $300\ \mu\text{m}$  thick inside a aluminium case (in figure 3, right), allow to reduce the minimum detectable neutron energy from 30 MeV to 10 MeV.

The coincidence between the different detectors guarantees the suppression of the  $\gamma$  background and the events that do not come from the  $\text{C}_2\text{H}_4$  target. Moreover, thanks to the multi-stage structure it is possible to select, for each neutron energy, the best configuration in terms of number of detectors in coincidence, to reduce the background coming from  $n+^{12}\text{C}$  reactions. The two telescopes were installed at different angles: the upstream detector was placed at  $\theta = 25^\circ$  to the neutron beam direction, and the downstream one at  $\alpha = 20^\circ$ . This allowed to increase the statistics, as well as to minimise systematic uncertainties related to the angular position.

Considering the very fast response of the scintillators, the telescope is able to identify events generated by neutrons arriving as close as 80 ns after the  $\gamma$ -flash, corresponding to an energy of  $\sim 1$  GeV.

The third telescope, indicated in figure 2 as RPT-PTB (shown in detail in figure 4), consists of three discrete detectors.

This triple-stage RPT has two transmission detectors that consist of a square EJ204 plastic scintillator,  $45 \times 45\ \text{mm}^2$  ( $\Delta\text{E1}$ ) and  $38 \times 38\ \text{mm}^2$  ( $\Delta\text{E2}$ ) in size, respectively. The thicknesses of the detectors are optimised to the neutron energy range to be covered and vary between 0.5 mm and 5 mm. Moreover the dimension of the  $\Delta\text{E2}$  detector, which has the smallest transverse dimensions of all the three detectors, defines the solid angle of the RPT-PTB. The stop detector (E) is a cylindrical EJ204 scintillator of 80 mm diameter. Also this detector thickness is optimized with respect to the energy range: 50 mm for neutron energy from 25 MeV to 100 MeV, 100 mm from 50 MeV to 150 MeV. The  $\Delta\text{E}$  scintillator detectors are inserted in a rooflike housings with 0.1 mm aluminium walls covered with a diffuse white reflector on the inside and coupled to a XP2020Q PMTs (figure 4, right). Due to the diffuse reflec-



**Figure 4.** Left: structure of the triple-stage RPT-PTB: three plastic scintillators. Right: detail of the interior of the cover for the  $\Delta\text{E}$  detectors with a diffuse white reflector to reduce the inhomogeneity of the light collection.

tor the inhomogeneity of the light collection is less than 10%. The envelope and the front side of the cylindrical E detectors are also covered with a white diffuse reflector and coupled to XP2020 PMTs.

The main design goal for this instrument is to achieve directional sensitivity for background suppression and reduced influence of angular straggling for an accurate definition of the solid angle, requiring, during the data analysis, a triple coincidence.

### 3 The measurement campaign

The  $^{235}\text{U}(n,f)$  measurement at n\_TOF required a five-week beamtime in order to collect enough statistics. It started on September 17<sup>th</sup> and ended on October 29<sup>th</sup>, for a total of  $3.8 \cdot 10^{18}$  protons on target.

During the measurement campaign different  $\text{C}_2\text{H}_4$  targets were used; their thicknesses were optimized to match the neutron-energy range of interest. In particular for the low neutron energy range (from 10 to 200 MeV) targets with 1 and 2 mm of thickness were used, whereas the corresponding carbon samples to subtract the background reaction were of 0.5 and 1 mm thickness. For higher neutron energies, up to 1 GeV, polyethylene and carbon samples with 5 and 2.5 mm thickness were used.

The RTP-PTB and the RPT-INFN with the Silicon detectors were at the same angle with respect to the neutron beam,  $25^\circ$ , and pointing at the same sample. While the third telescope, at  $20^\circ$ , was pointing always at the thicker target, to increase the statistics for the high energy part. In the table 1 the protons on target for each detector and configuration are summarized.

### 4 Conclusion

The  $^{235}\text{U}(n,f)$  reaction is one of the most important standard cross section both in nuclear physics and for its

**Table 1.** Summary of statistics (protons on target, "pot") collected for each configuration used during the measurement campaign at n\_TOF.

| Sample                        | Thickness | RPTs at 25°<br>(pot) | RPT at 20°<br>(pot)  |
|-------------------------------|-----------|----------------------|----------------------|
| C <sub>2</sub> H <sub>4</sub> | 1 mm      | $7.08 \cdot 10^{17}$ | -                    |
| C <sub>2</sub> H <sub>4</sub> | 2 mm      | $8.41 \cdot 10^{17}$ | -                    |
| C <sub>2</sub> H <sub>4</sub> | 5 mm      | $7.44 \cdot 10^{17}$ | $2.06 \cdot 10^{18}$ |
| C                             | 0.5 mm    | $3.27 \cdot 10^{17}$ | -                    |
| C                             | 1 mm      | $3.86 \cdot 10^{17}$ | -                    |
| C                             | 2.5 mm    | $3.14 \cdot 10^{17}$ | $1.55 \cdot 10^{18}$ |

widespread use as reference for neutron flux measurements. Nevertheless few measurements are available in the neutron energy range between 20 and 200 MeV and none above 200 MeV. To fulfil this lack of experimental data a campaign to measure the <sup>235</sup>U fission reaction, relative to the <sup>1</sup>H(n,n)<sup>1</sup>H reaction, which is considered the primary reference in this energy region, has been carried on at the n\_TOF facility at CERN in 2018. The setup was designed to measure simultaneously the fission reactions by means of two fission fragment detectors and the neutron flux using three different recoil proton telescopes. During the experimental campaign data have been acquired using different C<sub>2</sub>H<sub>4</sub> and C targets for the evaluation of the neutron flux and systematics. The data analysis is in progress together with an intense MC simulation campaign to evaluate the detectors efficiencies using Geant-4 toolkit [19].

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