

The Research Program at RIBRAS (Radioactive Ion Beams in Brasil)-III

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The Research Program at RIBRAS (Radioactive Ion Beams in Brasil)-III

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Abstract. A part of the research program developed in the RIBRAS facility over the last four years is presented. Experiments using radioactive secondary beams of light exotic nuclei such as ⁶He, ⁷Be, ⁸Li on several targets have been performed. Elastic angular distributions have been analysed by the Optical Model and four body Continuous Discretized Coupled Channels Calculations (4b-CDCC) and the total reaction cross sections have been obtained. A comparison between the reaction cross sections of ⁶He and other stable projectiles with medium-heavy targets was performed. Measurements of the proton transfer reaction ¹²C(⁸Li, ⁹Be)¹¹B are also presented.

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THE RIBRAS FACILITY

The RIBRAS (Radioactive Ion Beams in Brasil) [1] facility consists of two superconducting solenoids capable of producing secondary beams of exotic nuclei. The solenoids are presently installed in one of the beam lines of the 8 MV Pelletron Tandem accelerator of the University of Sao Paulo. The solenoids make an in-flight selection of the reaction products emerging from the primary target and focus them on a secondary target where the reaction of interest takes place. The ⁷Li primary beam and a ⁹Be (~12 μm) target are the standard to produce the ⁶He and ⁸Li beams via the one proton stripping ⁹Be(⁷Li, ⁶He) and one neutron pickup ⁹Be(⁷Li, ⁸Li) reactions. A ³He gas primary target is used to produce the ⁸B and ⁷Be beams by the ³He(⁶Li, ⁸B) and ³He(⁷Li, ⁷Be) reactions. The primary beam traverses the primary target (see label (1) in Figure 1) and stops in the Faraday cup (3) which measures its current. The secondary

particles are collected by the first solenoid within a cone between $2^\circ \leq \theta \leq 6^\circ$ defined by the Faraday cup and the collimator (2). The acceptance solid angle of the system is 30 msr. The first solenoid makes a magnetic-rigidity selection to focus the secondary beam in the secondary target position (7). Unwanted secondary particles are stopped in the blocker (5) and collimator (6). A second identical solenoid is mounted after the target (7) and will allow an additional filtering improving the purity of the secondary beams. Presently we operate only with the first solenoid.

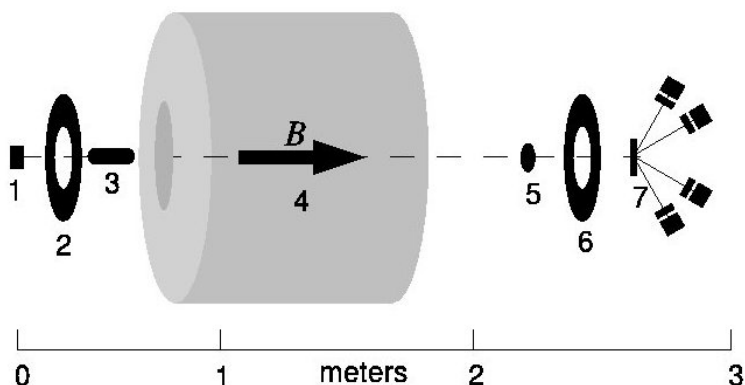


FIGURE 1. Scheme of the first solenoid RIBRAS system. –1-primary target, 2-collimator, 3-Faraday cup, 4-solenoid, 5-unwanted-beam blocker, 6-collimator, 7-secondary target and detectors.

This system allows the production of beams of the light nuclei mentioned above with intensities of 10^4 to 10^6 pps and energies in the range 1-5 MeV/A.

THE RESEARCH PROGRAM AT RIBRAS

Elastic scattering and Total Reaction Cross Section.

Most of the experiments performed at RIBRAS consist of elastic scattering and nucleon transfer reactions measurements using light exotic beams such as ${}^6\text{He}$, ${}^8\text{Li}$ and ${}^7\text{Be}$. The elastic scattering allows the study of the interacting potential and the total

reaction cross section. We have been developing this program using the medium-heavy target ^{120}Sn [2], medium mass targets such as ^{58}Ni and ^{51}V [3] and lighter targets such as ^{27}Al [4], ^{12}C [5,6] and ^9Be [7].

The angular distributions have been analyzed by optical model and four-body CDCC calculations. In Figure 2 we present the elastic angular distributions for the $^6\text{He}+^{120}\text{Sn}$ system compared to Optical Model and 4b-CDCC calculations[8].

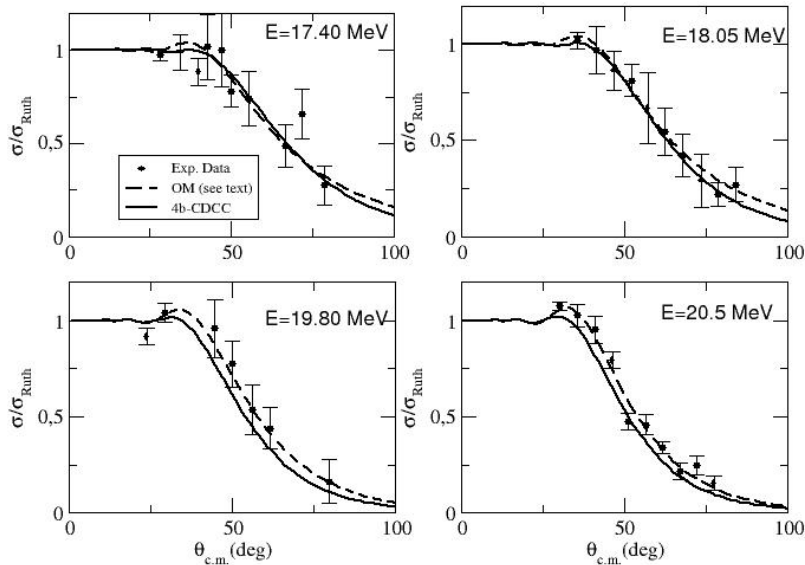


FIGURE 2. Elastic scattering angular distributions for the $^6\text{He}+^{120}\text{Sn}$. Dashed line is Optical Model and solid line is the 4b-CDCC calculation [8].

To compare the reaction cross sections of different systems at different energies, it is convenient to rescale the cross sections and energies in order to overcome trivial effects due to the different sizes and energies with respect to the Coulomb barrier. We apply two methods to reduce the cross sections and energies [9,10]: the first is to calculate the ratio to the square of the system radius $\sigma_{red} = \sigma / (A_p^{1/3} + A_t^{1/3})^2$ and $E_{red} = E_{cm} (A_p^{1/3} + A_t^{1/3}) / Z_p Z_t$, the second is based on the Wong's [11] formula which depends on the quantity $\chi = (E - E_{CB}) / \eta w$ where ηw is the curvature of the barrier.

With both methods (see Figure 3) we observe an increase of the total reaction cross section for the system $^6\text{He} + ^{120}\text{Sn}$ in comparison with the weakly-bound stable projectiles $^6,7\text{Li}$, ^9Be and tightly bound projectiles as ^{16}O and alpha particles. A similar

behavior was observed also for medium mass systems [9] and it was initially attributed to the larger breakup probability of ${}^6\text{He}$. However, there are strong indications that the neutron transfer stripping reactions to bound states are very important reaction channels contributing to the total reaction cross section of ${}^6\text{He} + \text{X}$. A large yield of alpha particles was observed in the ${}^6\text{He} + {}^{120}\text{Sn}$ collision, with energies matching the ones expected from 2n-stripping reaction kinematics[2].

Reduced total reaction cross sections for medium-heavy mass systems

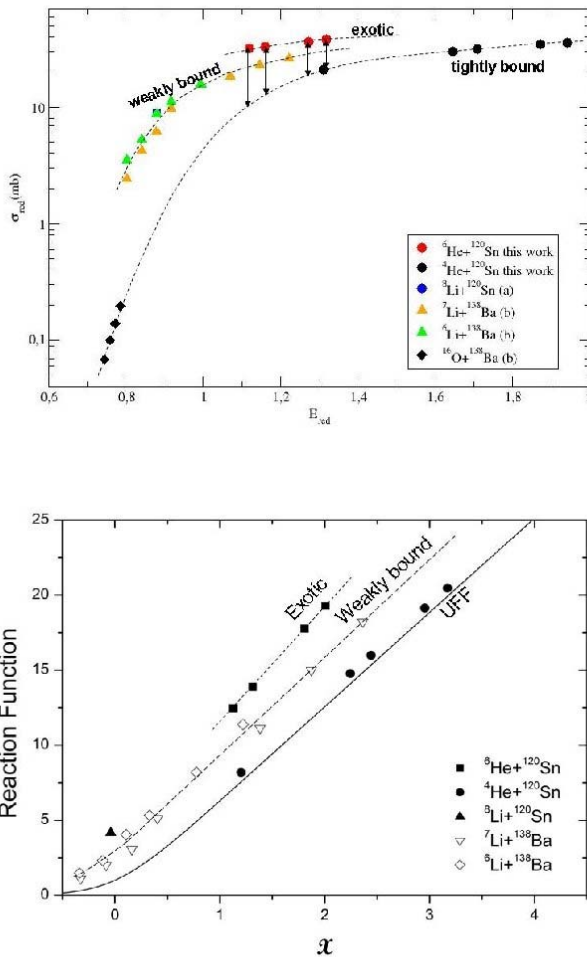


FIGURE 3. Total reduced reaction cross sections for medium-heavy mass systems with exotic, weakly bound and tightly bound projectiles. Dashed lines are guides to the eyes. Two reduction methods are used in the top and bottom figures (see text for more details). In the top figure, the arrows display the difference of the cross section between the exotic projectile ${}^6\text{He}$ and the tightly bound ${}^4\text{He}$. In the bottom figure, the solid line is the *Universal Fusion Function* (UFF) derived from the Wong formula [see Ref. 9]. (a) stands for Ref. [15] and (b) for Ref.[9]

For lighter systems such as ${}^6\text{He}+{}^{27}\text{Al}$, ${}^6\text{He}+{}^{12}\text{C}$ and ${}^6\text{He}+{}^9\text{Be}$ the situation is still inconclusive. There are some indications that the enhancement observed in ${}^6\text{He}+$ heavy targets could be smaller for lighter systems such as the ${}^6\text{He}+{}^{27}\text{Al}$ [4] and ${}^8\text{B}+{}^{12}\text{C}$ [6], although a larger reduced reaction cross section has been obtained in some analysis [9].

Nucleon Transfer Reactions

Nucleon transfer reactions induced by exotic projectiles provide, in many cases, important spectroscopic information about states of nuclei not accessible by other means. In particular, the unstable lithium isotopes such as ${}^8,9\text{Li}$ [6,12,13,14] are important under the nuclear structure point of view and also for astrophysics as they could help to bridge the $A = 8$ gap. In Figure 4 we present an angular distribution of the proton transfer ${}^{12}\text{C}({}^8\text{Li}, {}^9\text{Be}){}^{11}\text{B}$ reaction measured at RIBRAS. As the vertex $\langle {}^{11}\text{B}|{}^{12}\text{C} \rangle$ is known, this reaction provides information of the $\langle {}^8\text{Li}|{}^9\text{Be} \rangle$ spectroscopic factor, which normalizes the radiative capture ${}^8\text{Li} + p \rightarrow {}^9\text{Be}$ cross section. In Figure 5 we present the reaction rates for the above capture reaction derived using the spectroscopic factor $C^2S({}^9\text{Be}) = 1.10(25)$ obtained from this work. We compare our results with other measurements.

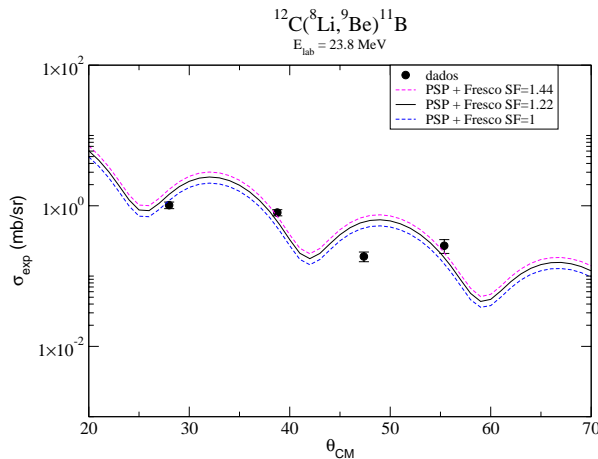


Figure 4 shows the angular distribution of the proton transfer reaction ${}^{12}\text{C}({}^8\text{Li}, {}^9\text{Be}){}^{11}\text{B}$ at $E_{\text{lab}} = 23.8$ MeV. The lines are DWBA calculations.

$E_{\text{lab}} = 23.8$ MeV. The lines are

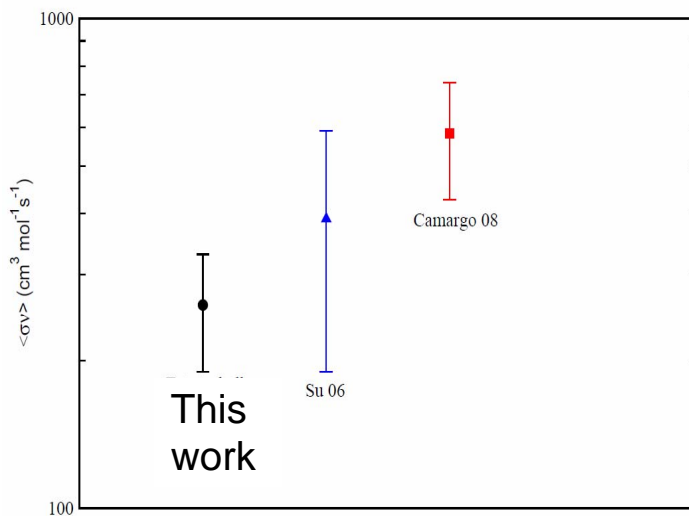


FIGURE 5. ${}^8\text{Li} + p \rightarrow {}^9\text{Be}$ reaction rate calculated at $T_9 = 1$. Su06 stands for Ref. [13] and Camargo 08 for Ref. [14].

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