# LEIBNIZ ALGEBRAS OF HEISENBERG TYPE 

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#### Abstract

We introduce and provide a classification theorem for the class of Heisenberg-Fock Leibniz algebras. This category of algebras is formed by those Leibniz algebras $L$ whose corresponding Lie algebras are Heisenberg algebras $H_{n}$ and whose $H_{n}$-modules $I$, where $I$ denotes the ideal generated by the squares of elements of $L$, are isomorphic to Fock modules. We also consider the three-dimensional Heisenberg algebra $H_{3}$ and study three classes of Leibniz algebras with $H_{3}$ as corresponding Lie algebra, by taking certain generalizations of the Fock module. Moreover, we describe the class of Leibniz algebras with $H_{n}$ as corresponding Lie algebra and such that the action $I \times H_{n} \rightarrow I$ gives rise to a minimal faithful representation of $H_{n}$. The classification of this family of Leibniz algebras for the case of $n=3$ is given.


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## 1. Introduction

The term Leibniz algebra was introduced in the study of a non-antisymmetric analogue of Lie algebras by Loday [35], being so the class of Leibniz algebras an extension of the one of Lie algebras. However this kind of algebras was previously studied under the name of $D$-algebras by D. Bloh [10, 11, 12]. Since the 1993 Loday's work many researchers have been attracted by this category of algebras, being remarkable the great activity in this field developed in the last years. This activity has been mainly focussed in the frameworks of low dimensional algebras, nilpotence and physics applications (see [2, 5, 6, 14, 15, 16, 17, 21, 22, 23, 25, 28, 33, 40, 41|).

Definition 1. A Leibniz algebra $L$ is a linear space over a base field $\mathbb{F}$ endowed with a bilinear product $[\cdot, \cdot]$ satisfying the Leibniz identity

$$
[[y, z], x]=[[y, x], z]+[y,[z, x]],
$$

for all $x, y, z \in L$.
In presence of anti-commutativity, Jacobi identity becomes Leibniz identity and therefore Lie algebras are examples of Leibniz algebras. Throughout this paper $\mathbb{F}$ will be algebraically closed and with zero characteristic.

Let $L$ be a Leibniz algebra. The ideal $I$ generated by the squares of elements of the algebra $L$, that is $I$ is generated by the set $\{[x, x]: x \in L\}$, plays an important role in the theory since it determines the (possible) non-Lie character of $L$. From the Leibniz identity, this ideal satisfies

$$
[L, I]=0
$$

The quotient algebra $L / I$ is a Lie algebra, called the corresponding Lie algebra of $L$, and the map $I \times L / I \rightarrow I$, $(i,[x]) \mapsto[i, x]$, endows $I$ of a structure of $L / I$-module (see [4, 37]). Observe that we can write

$$
\begin{equation*}
L=V \oplus I \tag{1}
\end{equation*}
$$

where $V$ is a linear complement of $I$ in $L$ and $V$ is isomorphic as linear space to $L / I$. From here, Leibniz algebras give us the opportunity of treating in an unifying way a Lie algebra together with a module over it.

On the other hand, we recall that Heisenberg (Lie) algebras play an important role in mathematical physics and geometry, in particular in Quantum Mechanics (see for instance [1, 8, 9, 19, 20, 24, 26, 27, 29, 30, 31, 32, 36, 42, 44]). Indeed, the Heisenberg Principle of Uncertainty implies the non-compatibility of position and momentum observables acting on fermions. This non-compatibility reduces to non-commutativity of the corresponding operators. If we represent by $\bar{x}$ the operator associated to position and by $\frac{\bar{\partial}}{\partial x}$ the one associated to momentum (acting for instance on a space $V$ of differentiable functions of a single variable), then $\left[\bar{x}, \frac{\bar{\partial}}{\partial x}\right]=\overline{1}_{V}$ which is non-zero. Thus we can identify the subalgebra generated by $\overline{1}, \bar{x}$ and $\frac{\bar{\partial}}{\partial x}$ with the three-dimensional

Heisenberg algebra whose multiplication table in the basis $\left\{\overline{1}, \bar{x}, \frac{\bar{\partial}}{\partial x}\right\}$ has as unique non-zero product $\left[\bar{x}, \frac{\bar{\partial}}{\partial x}\right]=$ $\overline{1}$.

For any non-negative integer $k$ the Heisenberg algebra of dimension $n=2 k+1$ (denoted further by $H_{n}$ ) is characterized by the existence of a basis

$$
\begin{equation*}
B=\left\{\overline{1}, \bar{x}_{1}, \frac{\bar{\delta}}{\delta x_{1}}, \ldots, \bar{x}_{k}, \frac{\bar{\delta}}{\delta x_{k}}\right\} \tag{2}
\end{equation*}
$$

in which the multiplicative non-zero relations are

$$
\left[\bar{x}_{i}, \frac{\bar{\delta}}{\delta x_{i}}\right]=-\left[\frac{\bar{\delta}}{\delta x_{i}}, \bar{x}_{i}\right]=\overline{1}
$$

for $1 \leq i \leq k$.
In the present paper we are focusing in introducing and studying several classes of Leibniz algebras whose corresponding Lie algebras are Heisenberg algebras $H_{n}$. Recall that there is a unique irreducible representation of the Heisenberg algebra (at least a unique one that can be exponentiated). This is why physicists are able to use the Heisenberg commutation relations to do calculations, without worry about what they are being represented on. This representation is called the Fock (or Bargmann-Fock) representation (see [3, 7, 34, 38, 39, 43]). Physically this representation corresponds to an harmonic oscillator, with the vector $\overline{1} \in \mathbb{C}[x]$ as the vacuum state and $\bar{x}$ the operator that adds one quantum to the vacuum state. This representation is also sometimes known as the oscillator representation. For a given Heisenberg algebra $H_{n}, n=2 k+1$, this representation gives rise to the so-called Fock module on $H_{n}$, the linear space $\mathbb{F}\left[x_{1}, \ldots, x_{k}\right]$ with the action induced by

$$
\begin{array}{ll}
\left(p\left(x_{1}, \ldots, x_{k}\right), \overline{1}\right) & \mapsto p\left(x_{1}, \ldots, x_{k}\right) \\
\left(p\left(x_{1}, \ldots, x_{k}\right), \bar{x}_{i}\right) & \mapsto x_{i} p\left(x_{1}, \ldots, x_{k}\right)  \tag{3}\\
\left(p\left(x_{1}, \ldots, x_{k}\right), \frac{\bar{\delta}}{\delta x_{i}}\right) & \mapsto
\end{array} \frac{\delta}{\delta x_{i}}\left(p\left(x_{1}, \ldots, x_{k}\right)\right) \text {. }
$$

for any $p\left(x_{1}, \ldots, x_{k}\right) \in \mathbb{F}\left[x_{1}, \ldots, x_{k}\right]$ and $i=1, \ldots, k$.
Taking now into account the above comments, we introduce in Section 2 the class of Heisenberg-Fock Leibniz algebras as those Leibniz algebras whose corresponding Lie algebras are $H_{n}$ and whose $H_{n}$-modules $I$ are isomorphic to Fock modules, and provide a classification theorem. Thus, we have the opportunity of considering Heisenberg Lie algebras together with their Fock representations in a unifying viewpoint. In this section we also consider a generalization of this class of algebras by means of a direct sum of Heisenberg algebras as corresponding Lie algebras, and provide also a classification theorem.

In Section 3, we center in the three-dimensional Heisenberg algebra $H_{3}$ and study three classes of Leibniz algebras with $H_{3}$ as corresponding Lie algebra by taking certain generalizations of the Fock module. We also note that Sections 2 and 3 allow us to introduce several new classes of infinite-dimensional Leibniz algebras.

Finally, in Section 4, we deal with the category of Leibniz algebras with $H_{n}$ as corresponding algebra and such that the action $I \times H_{n} \rightarrow I$ gives rise to a minimal faithful representation of $H_{n}$. A description of this category of algebras is given and also a classification theorem when $n=3$.

## 2. Classification of Heisenberg-Fock type Leibniz algebras

2.1. Classification of $H F L_{n}$. Consider a Heisenberg algebra $H_{n}$, with $n=2 k+1$, and its Fock module $\mathbb{F}\left[x_{1}, \ldots, x_{k}\right]$ under the action (3). The Heisenberg-Fock Leibniz algebra $H F L_{n}$ is defined as the Leibniz algebra with corresponding Lie algebra $H_{n}$ and such that the action $I \times H_{n} \rightarrow I$ makes of $I$ the Fock module. Since $\mathbb{F}\left[x_{1}, \ldots, x_{k}\right]$ is infinite-dimensional we get a family of infinite-dimensional Leibniz algebras.
Theorem 1. The Heisenberg-Fock Leibniz algebra HF $L_{n}$ admits a basis

$$
\left\{\overline{1}, \bar{x}_{i}, \frac{\bar{\delta}}{\delta x_{i}}, x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}} \mid t_{i} \in \mathbb{N} \cup\{0\}, 1 \leq i \leq k\right\}
$$

in such a way that the multiplication table on this basis has the form:

$$
\begin{array}{ll}
{\left[\overline{x_{i}}, \frac{\bar{\delta}}{\delta x_{i}}\right]=\overline{1},} & 1 \leq i \leq k \\
{\left[\frac{\bar{\delta}}{\delta x_{i}}, \overline{x_{i}}\right]=-\overline{1},} & 1 \leq i \leq k
\end{array}
$$

$$
\begin{array}{ll}
\left.x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \overline{1}\right]=x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \\
{\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{k_{k}}, \bar{x}_{i}\right]=x_{1}^{t_{1}} \ldots x_{i-1}^{i_{i-1}} x_{i}^{t_{i}+1} x_{i+1}^{t_{i+1}} \ldots x_{k}^{t_{k}},} & 1 \leq i \leq k, \\
{\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{k_{k}}, \frac{\delta}{\delta x_{i}}\right]=t_{i} x_{1}^{t_{1}} \ldots x_{i-1}^{t_{i-1}} x_{i}^{t_{i}-1} x_{i+1}^{t_{i+1}} \ldots x_{k}^{t_{k}}, \quad 1 \leq i \leq k,}
\end{array}
$$

where the omitted products are equal to zero.
Proof. Taking into account Equations (1) and (3) we conclude that

$$
\left\{\overline{1}, \bar{x}_{i}, \frac{\bar{\delta}}{\delta x_{i}}, x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}} \mid t_{i} \in \mathbb{N} \cup\{0\}, 1 \leq i \leq k\right\}
$$

is a basis of $H F L_{n}$ and

$$
\begin{aligned}
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \overline{1}\right]=x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}},} \\
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \bar{x}_{i}\right]=x_{1}^{t_{1}} \ldots x_{i-1}^{t_{i}} x_{i}^{t_{i+1}} x_{i+1}^{t_{i+1}} \ldots x_{k}^{t_{k}},} \\
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \frac{\bar{\delta}}{\delta x_{i}}\right]=t_{i} x_{1}^{t_{1}} \ldots x_{i-1}^{t_{i-1}} x_{i}^{i_{i}-1} x_{i+1}^{t_{i+1}} \ldots x_{k}^{t_{k}},}
\end{aligned}
$$

for $1 \leq i \leq k$.
Observe that we can write

$$
\begin{array}{ll}
{\left[\overline{x_{i}}, \overline{1}\right]=p_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & 1 \leq i \leq k, \\
{\left[\bar{\delta}, \overline{\delta x}, \overline{1}=q_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),\right.} & 1 \leq i \leq k, \\
{[\overline{1}, \overline{1}]=r\left(x_{1}, x_{2}, \ldots, x_{k}\right),} &
\end{array}
$$

where $p_{i}, q_{i}, r \in \mathbb{F}\left[x_{1}, \ldots, x_{k}\right]$.
Taking the following change of basis,

$$
\begin{gathered}
{\overline{x_{i}}}^{\prime}=\overline{x_{i}}-p_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right), \quad 1 \leq i \leq k, \\
\frac{\bar{\delta}^{\prime}}{\delta x_{i}}=\frac{\bar{\delta}}{\delta x_{i}}-q_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right), \quad 1 \leq i \leq k, \\
\overline{1}^{\prime}=\overline{1}-r\left(x_{1}, x_{2}, \ldots, x_{k}\right),
\end{gathered}
$$

we derive

$$
\left[\overline{x_{i}}, \overline{1}\right]=0, \quad\left[\frac{\bar{\delta}}{\delta x_{i}}, \overline{1}\right]=0, \quad[\overline{1}, \overline{1}]=0, \quad 1 \leq i \leq k .
$$

Now denote

$$
\begin{array}{lll}
{\left[\overline{x_{i}}, \overline{x_{j}}\right]=a_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & {\left[\frac{\bar{\delta}}{\overline{\delta x_{i}}}, \frac{\bar{\delta}}{\frac{\delta}{x_{j}}}\right]=b_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & 1 \leq i, j \leq k, \\
{\left[\frac{\bar{\delta}}{\delta \delta_{i}}, \overline{x_{j}}\right]=c_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & {\left[\overline{x_{i}}, \frac{\delta}{\delta x_{j}}\right]=d_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & 1 \leq i, j \leq k, i \neq j, \\
{\left[\overline{x_{i}}, \frac{\bar{\delta}}{\delta x_{i}}\right]=\overline{1}+e_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & {\left[\frac{\bar{\delta}}{\delta x_{i}}, \overline{x_{i}}\right]=-\overline{1}+f_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & 1 \leq i \leq k, \\
{\left[\overline{1}, \overline{x_{i}}\right]=h_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & {\left[\overline{1}, \frac{\bar{\delta}}{\delta x_{i}}\right]=g_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right),} & 1 \leq i \leq k .
\end{array}
$$

The Leibniz identity on the following triples imposes further constraints on the products.

| Leibniz identity | Constraint |
| :---: | :---: |
| $\left\{\overline{x_{i}}, \overline{x_{j}}, \overline{1}\right\}$ | $\Rightarrow \quad a_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i, j \leq k$, |
| $\left\{\frac{\bar{\delta}}{\bar{\delta} x_{i}}, \frac{\bar{\delta}}{\delta x_{j}}, \overline{1}\right\}$ | $\Rightarrow \quad b_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i, j \leq k$, |
| $\left\{\frac{\bar{\delta}}{\bar{\delta} x_{i}}, \overline{j_{j}}, \overline{1}\right\}$ | $\Rightarrow \quad c_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i, j \leq k, i \neq j$, |
| $\left\{\overline{x_{i}}, \frac{\bar{\delta}}{\delta x_{j}}, \overline{1}\right\}$ | $\Rightarrow \quad d_{i, j}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i, j \leq k, i \neq j$, |
| $\left\{\overline{x_{i}}, \frac{\bar{\delta}}{\delta x_{i}}, \overline{1}\right\}$ | $\Rightarrow \quad e_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i \leq k$, |
| $\left\{\frac{\bar{\delta}}{\overline{\delta x}}, \overline{x_{i}}, \overline{1}\right\}$ | $f_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i \leq k$, |
| $\left\{\overline{1}, \overline{x_{i}}, \overline{1}\right\}$ | $h_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i \leq k$, |
| $\{\overline{1}, \bar{\delta}, \bar{\delta}\}$ | $\Rightarrow \quad g_{i}\left(x_{1}, x_{2}, \ldots, x_{k}\right)=0, \quad 1 \leq i \leq k$. |

The proof is complete.
2.2. Classification of generalized Heisenberg-Fock Leibniz algebras. In this subsection we are interested in classifying the class of (infinite-dimensional) Leibniz algebras formed by those Leibniz algebras $L$ satisfying that their corresponding Lie algebras are finite direct sums of Heisenberg algebras and that the actions on $I$ are induced by Fock representations.

Since

$$
\begin{equation*}
L / I \cong H_{2 k_{1}+1} \oplus H_{2 k_{2}+1} \oplus H_{2 k_{3}+1} \oplus \cdots \oplus H_{2 k_{s}+1} \tag{4}
\end{equation*}
$$

we easily get

$$
\begin{equation*}
\mathcal{B}_{i}:=\left\{\overline{1_{i}}, \overline{x_{1, i}}, \overline{x_{2, i}}, \ldots, \overline{x_{k_{i}, i}}, \frac{\bar{\delta}}{\delta x_{1, i}}, \frac{\bar{\delta}}{\delta x_{2, i}}, \ldots, \frac{\bar{\delta}}{\delta x_{k_{i}, i}}\right\} \tag{5}
\end{equation*}
$$

for the standard basis of $H_{2 k_{i}+1}, i \in\{1,2, \ldots, s\}$.
We put

$$
\begin{equation*}
I=\mathbb{F}\left[x_{1}, \ldots, x_{n}\right] \tag{6}
\end{equation*}
$$

where $n=k_{1}+k_{2}+\cdots+k_{s}$.
The action

$$
I \times L / I \rightarrow I
$$

given by

$$
\begin{array}{ll}
\left(p\left(x_{1}, \ldots, x_{n}\right), \overline{1_{i}}\right) & \mapsto p\left(x_{1}, \ldots, x_{n}\right) \\
\left(p\left(x_{1}, \ldots, x_{n}\right), \overline{x_{j, i}}\right) & \mapsto p\left(x_{1}, \ldots, x_{n}\right) x_{k_{1}+k_{2}+\cdots+k_{i-1}+j} \\
\left(p\left(x_{1}, \ldots, x_{n}\right), \frac{\bar{\delta}}{\delta x_{j, i}}\right) & \mapsto
\end{array} \frac{\delta}{\delta x_{k_{1}+k_{2}+\cdots+k_{i-1}+j}} p\left(x_{1}, \ldots, x_{n}\right)
$$

for any $p\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{F}\left[x_{1}, \ldots, x_{n}\right]$ and $(i, j)$ with $i \in\{1,2, \ldots, s\}, j \in\left\{1, \ldots, k_{i}\right\}$, endows $I$ of a structure of $L / I$-module. Hence, we get a new family of Heisenberg-Fock type Leibniz algebras which generalize the previous ones considered in $\S 2.1$ (case $s=1$ ), that we call generalized Heisenberg-Fock Leibniz algebras, by introducing the algebras $L=L / I \oplus I$ with $L / I$ and $I$ as in Equations (4) and (6). We will denote them as

$$
H F L_{2 k_{1}+1,2 k_{2}+1, \ldots, 2 k_{s}+1}
$$

Our aim is to classify this class of Leibniz algebras.
By taking into account the previous arguments, it is clear that for any $i \in\{1,2, \ldots, s\}$ we have [ $\left.H_{2 k_{i}+1}, H_{2 k_{i}+1}\right] \subset H_{2 k_{i}+1}$ being the multiplication table among the elements in the basis $\mathcal{B}_{i}$ as in Theorem 1 . Therefore, we only need to study the products $\left[H_{2 k_{i}+1}, H_{2 k_{j}+1}\right]$ with $i, j \in\{1,2, \ldots, s\}$ and $i \neq j$.
Lemma 1. Let $a \in \mathcal{B}_{i}$ and $b \in \mathcal{B}_{j}, i, j \in\{1,2, \ldots, s\}$ with $i \neq j$. Then $[a, b]=0$.
Proof. For $i \neq j$ we have $[a, b]=p$ and $\left[b, \overline{1_{i}}\right]=q$ for some $p, q \in \mathbb{F}\left[x_{1}, \ldots, x_{n}\right]$. Taking now into account Theorem 1 we derive $\left[a, \overline{1_{i}}\right]=0$ and so

$$
p=\left[[a, b], \overline{1_{i}}\right]=\left[\left[a, \overline{1_{i}}\right], b\right]+\left[a,\left[b, \overline{1_{i}}\right]\right]=0
$$

The next theorem is now consequence of Theorem 1 and Lemma 1
Theorem 2. The Leibniz algebra HFL $L_{2 k_{1}+1,2 k_{2}+1, \ldots, 2 k_{s}+1}$ admits a basis (see Equations (5) and (6))

$$
\mathcal{B}_{1} \dot{\cup} \mathcal{B}_{2} \dot{\cup} \cdots \dot{\cup} \mathcal{B}_{s} \dot{\cup}\left\{x_{1}^{t_{1}} x_{2}^{t_{2}} \cdots x_{n}^{t_{n}} \mid t_{i} \in \mathbb{N} \cup\{0\}, 1 \leq i \leq n\right\}
$$

where $n=k_{1}+k_{2}+\cdots+k_{s}$, and in such a way that the multiplication table on this basis has the form:

$$
\begin{aligned}
& {\left[\overline{x_{j, i}}, \frac{\bar{\delta}}{\delta x_{j, i}}\right]=\overline{1_{i}}, \quad\left[\frac{\bar{\delta}}{\delta x_{j, i}}, \overline{x_{j, i}}\right]=-\overline{1_{i}},} \\
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{n}^{t_{n}}, \overline{1_{i}}\right]=x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{n}^{t_{n}},} \\
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{n}^{t_{n}}, \bar{x}_{j, i}\right]=x_{1}^{t_{1}} \ldots x_{k_{1}+\cdots+k_{k_{-1}+j-1}^{t_{k_{1}}+\cdots+k_{i-1}+j-1}}^{t_{k_{1}+\cdots+k_{i-1}+j+1}^{t_{k_{1}+\cdots+k_{i-1}+j}^{t_{1}}} x_{k_{k_{1}+\cdots+k_{i-1}+j+1}^{t_{k_{1}+\cdots+k_{i-1}+j+1}} \ldots x_{n}^{t_{n}},}^{t_{k_{1}+\cdots+k_{i-1}+j-1}^{t_{k_{1}}}, \cdots+k_{i-1}+j-1} x_{t_{k_{1}+\cdots+k_{i-1}+j+1}}^{t_{k_{1}+\cdots}} \ldots x_{n}^{t_{n}},}} \\
& {\left[x_{1}^{t_{1}} x_{2}^{t_{2}} \ldots x_{k}^{t_{k}}, \frac{\bar{\delta}}{\delta x_{j, i}}\right]=t_{k_{1}+\cdots+k_{i-1}+j} x_{1}^{t_{1}} \ldots x_{k_{1}+\cdots+k_{i-1}+j-1} x_{k_{1}+\cdots+k_{i-1}+j} x_{k_{1}+\cdots+k_{i-1}+j+1},}
\end{aligned}
$$

for $1 \leq i \leq s, 1 \leq j \leq k_{i}$ and where the omitted products are equal to zero.

## 3. Several degenerations of the Fock representation for the 3-dimensional Heisenberg algebra

In this section we consider several degenerations of the Fock representation of the Heisenberg algebra $H_{3}$. First, we study when an extension of the Fock action $\mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x]$, (see Equation (3)), by allowing arbitrary polynomials as results of the action of a fixed element in the basis $\left\{\overline{1}, \bar{x}, \frac{\bar{\delta}}{\delta x}\right\}$ of $H_{3}$ over the elements of $\mathbb{F}[x]$, makes of $\mathbb{F}[x]$ an $H_{3}$-module. Second, the new $H_{3}$-modules obtained in this way give rise to new classes of Leibniz algebras that will be described.

For any linear mapping $\Omega: \mathbb{F}[x] \rightarrow \mathbb{F}[x]$, consider the linear space $\mathbb{F}[x]$ with the action induced by the following applications:

$$
\begin{aligned}
& \psi_{1}: \mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x] \quad \psi_{2}: \mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x] \\
& (p(x), \overline{1}) \quad \mapsto \Omega(p(x)) \quad(p(x), \overline{1}) \quad \mapsto \quad p(x) \\
& (p(x), \bar{x}) \quad \mapsto x p(x) \quad(p(x), \overline{\bar{x}}) \quad \mapsto \quad \Omega(p(x)) \\
& \left(p(x), \frac{\bar{\delta}}{\delta x}\right) \quad \mapsto \frac{\delta}{\delta x} p(x) . \quad\left(p(x), \overline{\frac{\delta}{\delta x}}\right) \mapsto \frac{\delta}{\delta x} p(x) . \\
& \psi_{3}: \mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x] \\
& (p(x), \overline{1}) \quad \mapsto \quad p(x) \\
& (p(x), \bar{x}) \quad \mapsto \quad x p(x) \\
& \left(p(x), \frac{\bar{\delta}}{\delta x}\right) \mapsto \Omega(p(x))
\end{aligned}
$$

for any $p(x) \in \mathbb{F}[x]$.
From now on, let us denote by $\left\{x^{i}\right\}_{i \in \mathbb{N} \cup\{0\}}$ the canonical basis of $\mathbb{F}[x]$. By considering $\psi_{1}\left(p(x),\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]\right)$, it is immediate to get that the first action $\psi_{1}$ makes of $\mathbb{F}[x]$ an $H_{3}$-module if and only if $\Omega=1_{\mathbb{F}[x]}$. As consequence we have.

Proposition 1. The Leibniz algebras obtained from the first action $\psi_{1}$ are the same as those obtained in Theorem [1]

Consider now the second action $\psi_{2}: \mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x]$.
Proposition 2. The action $\psi_{2}$ makes of $\mathbb{F}[x]$ an $H_{3}$-module if and only if

$$
\begin{equation*}
\Omega\left(x^{i}\right)=x^{i+1}+\sum_{k=0}^{i} c_{k}\binom{i}{k} x^{i-k} \tag{7}
\end{equation*}
$$

where $\left\{c_{k}\right\}_{k \in \mathbb{N} \cup\{0\}}$ is a fixed sequence in $\mathbb{F}$ and $\binom{i}{k}$ are binomial coefficients.
Proof. Suppose $\mathbb{F}[x]$ is an $H_{3}$-module through the action $\psi_{2}$. Then we have

$$
x^{i}=\left[x^{i}, \overline{1}\right]=\left[x^{i},\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]\right]=\left[\left[x^{i}, \bar{x}\right], \frac{\bar{\delta}}{\delta x}\right]-\left[\left[x^{i}, \frac{\bar{\delta}}{\delta x}\right], \bar{x}\right]=\left[\left[x^{i}, \bar{x}\right], \frac{\bar{\delta}}{\delta x}\right]-\left[i x^{i-1}, \bar{x}\right]
$$

and so

$$
\begin{equation*}
\left[\left[x^{i}, \bar{x}\right], \frac{\bar{\delta}}{\delta x}\right]=x^{i}+\left[i x^{i-1}, \bar{x}\right] . \tag{8}
\end{equation*}
$$

Taking into account Equation (8), we can easily prove by induction (7). Indeed, for $i=0$ we get from (8) that $\left[[1, \bar{x}], \frac{\bar{\delta}}{\delta x}\right]=1$, which implies $[1, \bar{x}]=x+c_{0}=\Omega(1)$. For $i=1$ the same equation allows us to get $\left[[x, \bar{x}], \frac{\bar{\delta}}{\delta x}\right]=x+[1, \bar{x}]=2 x+c_{0}$ and so $[x, \bar{x}]=x^{2}+c_{0} x+c_{1}=\Omega(x)$.

Let the induction hypothesis true for $i=j$ and we will show it for $i=j+1$. Taking into account (8) we have

$$
\begin{aligned}
{\left[\left[x^{j+1}, \bar{x}\right], \frac{\bar{\delta}}{\delta x}\right] } & =x^{j+1}+\left[(j+1) x^{j}, \bar{x}\right]=x^{j+1}+(j+1)\left(x^{j+1}+\sum_{k=0}^{j} c_{k}\binom{j}{k} x^{j-k}\right)= \\
& =(j+2) x^{j+1}+\sum_{k=0}^{j} c_{k}(j+1)\binom{j}{k} x^{j-k}=
\end{aligned}
$$

$$
\begin{aligned}
& =(j+2) x^{j+1}+\sum_{k=0}^{j} c_{k}(j+1) \frac{j!}{k!(j-k)!} x^{j-k}= \\
& =(j+2) x^{j+1}+\sum_{k=0}^{j} c_{k} \frac{(j+1)!}{k!(j+1-k)!}(j+1-k) x^{j-k}
\end{aligned}
$$

From here

$$
\left[x^{j+1}, \bar{x}\right]=x^{j+2}+\sum_{k=0}^{j} c_{k} \frac{(j+1)!}{k!(j+1-k)!} x^{j+1-k}+c_{j+1}=x^{j+2}+\sum_{k=0}^{j+1} c_{k}\binom{j+1}{k} x^{j+1-k}
$$

that is,

$$
\Omega\left(x^{j+1}\right)=x^{j+2}+\sum_{k=0}^{j+1} c_{k}\binom{j+1}{k} x^{j+1-k}
$$

The converse is of immediate verification.
Proposition 3. Any Leibniz algebra obtained from the second action $\psi_{2}$ admits a basis

$$
\left\{\overline{1}, \bar{x}, \frac{\bar{\delta}}{\delta x}\right\} \dot{\cup}\left\{x^{i}: i \in \mathbb{N} \cup\{0\}\right\}
$$

in such a way that the multiplication table on this basis has the form:

$$
\begin{array}{ll}
{\left[x^{i}, \overline{1}\right]=x^{i},} & {\left[x^{i}, \bar{x}\right]=\Omega\left(x^{i}\right),} \\
{\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]=\overline{1},} & {\left[\frac{\left.x^{i}, \frac{\bar{\delta}}{\delta x}\right]=i x^{i-1}}{\delta x}, \bar{x}\right]=-\overline{1},}
\end{array}
$$

where the omitted products are equal to zero and $\Omega\left(x^{i}\right)$ satisfies Equation (7).
Proof. By Proposition 2 we have the restriction on $\Omega\left(x^{i}\right)$. On the other hand, we know

$$
\begin{array}{lll}
{\left[x^{i}, \overline{1}\right]=x^{i},} & {\left[x^{i}, \bar{x}\right]=\Omega\left(x^{i}\right),} & {\left[x^{i}, \frac{\bar{\delta}}{\delta x}\right]=i x^{i-1},} \\
{[\bar{x}, \overline{1}]=p(x),} & {\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]=\overline{1}+q(x),} & {[\bar{x}, \bar{x}]=a(x),} \\
{\left[\frac{\bar{\delta}}{\delta x}, \overline{1}\right]=r(x),} & {\left[\frac{\bar{\delta}}{\delta x}, \frac{\bar{\delta}}{\delta x}\right]=b(x),} & {\left[\frac{\bar{\delta}}{\delta x}, \bar{x}\right]=-\overline{1}+s(x),} \\
{[\overline{1}, \bar{x}]=c(x),} & {[\overline{1}, \overline{1}]=d(x),} & {\left[\overline{1}, \frac{\bar{\delta}}{\delta x}\right]=e(x) .}
\end{array}
$$

By making the change of basis $\overline{1}^{\prime}=\overline{1}+q(x)$ we can suppose that $\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]=\overline{1}$.
Now, from Leibniz identity we obtain the following equations:

| Leibniz identity | Constraint |  |
| :--- | :--- | :--- |
|  | $\Rightarrow$ | $c(x)=[d(x), \bar{x}]$, |
| $\{\overline{1}, \overline{1}, \overline{1}\}$ | $\Rightarrow$ | $e(x)=\frac{\delta}{\delta x}(d(x))$, |
| $\left\{\overline{1}, \overline{1}, \frac{\bar{\delta}}{\delta x}\right\}$ | $\Rightarrow$ | $[e(x), \bar{x}]=\frac{\delta}{\delta x}(c(x))-d(x)$, |
| $\left\{\overline{1}, \bar{x}, \frac{\delta}{\delta x}\right\}$ | $\Rightarrow$ | $a(x)=[p(x), \bar{x}]$, |
| $\{\bar{x}, \overline{1}, \bar{x}\}$ | $\Rightarrow$ | $d(x)=\frac{\delta}{\delta x}(p(x))$, |
| $\left\{\bar{x}, \overline{1}, \frac{\bar{\delta}}{\delta x}\right\}$ | $\Rightarrow$ | $p(x)+c(x)=\frac{\delta}{\delta x}(a(x))$, |
| $\left\{\bar{x}, \bar{x}, \frac{\delta}{\delta x}\right\}$ | $\Rightarrow$ | $s(x)=d(x)+[r(x), \bar{x}]$, |
| $\left\{\frac{\bar{\delta}}{\delta x}, \overline{1}, \bar{x}\right\}$ | $\Rightarrow$ | $b(x)=\frac{\delta}{\delta x}(a(x))$, |
| $\left\{\frac{\delta}{\delta x}, \overline{1}, \frac{\bar{\delta}}{\delta x}\right\}$ | $\Rightarrow$ | $[b(x), \bar{x}]=-e(x)-r(x)+\frac{\delta}{\delta x}(a(x))$. |

By making the next change of basis:

$$
\begin{aligned}
& \overline{1}^{\prime}=\overline{1}-\frac{\delta}{\delta x}(p(x)), \\
& \bar{x}^{\prime}=\bar{x}-p(x), \\
& \frac{\bar{\delta}}{\delta x}=\frac{\bar{\delta}}{\delta x}-r(x),
\end{aligned}
$$

we obtain the family of the proposition.

Finally we consider the third action $\psi_{3}: \mathbb{F}[x] \times H_{3} \rightarrow \mathbb{F}[x]$, being then

$$
\begin{aligned}
& {\left[x^{i}, \overline{1}\right]=x^{i}} \\
& {\left[x^{i}, \bar{x}\right]=x^{i+1}} \\
& {\left[x^{i}, \frac{\bar{\delta}}{\delta x}\right]=\Omega\left(x^{i}\right), \quad i \in \mathbb{N} \cup\{0\}}
\end{aligned}
$$

By arguing in a similar way to Propositions 2 and 3 we can prove the next results.
Proposition 4. The action $\psi_{3}$ makes of $\mathbb{F}[x]$ an $H_{3}$-module if and only if

$$
\begin{equation*}
\Omega\left(x^{i}\right)=i x^{i-1}+x^{i} c(x) \tag{9}
\end{equation*}
$$

for a fixed $c(x) \in \mathbb{F}[x]$ and $i \in \mathbb{N} \cup\{0\}$.
Proposition 5. Any Leibniz algebra obtained from the third action $\psi_{3}$ admits a basis

$$
\left\{\overline{1}, \bar{x}, \frac{\bar{\delta}}{\delta x}\right\} \dot{\cup}\left\{x^{i}: i \in \mathbb{N} \cup\{0\}\right\}
$$

in such a way that the multiplication table on this basis has the form:

$$
\begin{array}{ll}
{\left[x^{i}, \overline{1}\right]=x^{i},} & {\left[x^{i}, \bar{x}\right]=x^{i+1},} \\
{\left[\bar{x}, \frac{\bar{\delta}}{\delta x}\right]=\overline{1},} & {\left[x^{i}, \frac{\bar{\delta}}{\delta x}\right]=\Omega\left(x^{i}\right),} \\
\delta x \\
\bar{\delta} & \bar{x}]=-\overline{1}
\end{array}
$$

where the omitted products are equal to zero and $\Omega\left(x^{i}\right)$ satisfies Equation (9).

## 4. Leibniz algebras of minimal faithful representation-Heisenberg type

4.1. General case. Let $H_{2 m+1}$ be a Heisenberg algebra of dimension $2 m+1$, then it is well-known that its minimal faithful representations have dimension $m+2$, (see [13]). From now on, for a more comfortable notation, we will denote by

$$
\left\{x_{1}, x_{2}, \ldots, x_{m}, y_{1}, y_{2}, \ldots, y_{m}, z\right\}
$$

the standard basis of $H_{2 m+1}$, (see Equation (2)), where the non-zero products are

$$
\left[y_{i}, x_{i}\right]=-\left[x_{i}, y_{i}\right]=z
$$

By [18], we can take as minimal faithful representation the linear mapping

$$
\varphi: H_{2 m+1} \rightarrow \operatorname{End}(I)
$$

where $I$ is an $(m+2)$-dimensional linear space with a fixed basis $\left\{e_{1}, e_{2}, \ldots, e_{m+2}\right\}$, determined by

$$
\begin{array}{ll}
\varphi\left(x_{i}\right)=E_{1, i+1} & 1 \leq i \leq m \\
\varphi\left(y_{i}\right)=E_{i+1, m+2} & 1 \leq i \leq m \\
\varphi(z)=E_{1, m+2} &
\end{array}
$$

Here $E_{i, j}$ denotes the elemental matrix with 1 in the $(i, j)$ slot and 0 in the remaining places and we have $\varphi([x, y])(e)=\varphi(y)(\varphi(x)(e))-\varphi(x)(\varphi(y)(e))$ for any $x, y \in H_{2 m+1}$ and $e \in I$. Observe that $H_{2 m+1}$ corresponds to the $(m+2) \times(m+2)$ matrices

$$
\left(\begin{array}{cccccc}
0 & a_{2} & a_{3} & \ldots & a_{m+1} & c \\
0 & 0 & 0 & \ldots & 0 & b_{2} \\
0 & 0 & 0 & \ldots & 0 & b_{3} \\
\vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\
0 & 0 & 0 & \ldots & 0 & b_{m+1} \\
0 & 0 & 0 & \ldots & 0 & 0
\end{array}\right)
$$

This representation makes of $I$ an $H_{2 m+1}$-module under the action

$$
\begin{array}{llll}
\phi: & I \times H_{2 m+1} & \rightarrow I \\
& \left(e_{i+1}, x_{i}\right) & \mapsto &  \tag{10}\\
& \left(e_{m+2}, y_{i}\right) & \mapsto & 1 \leq i \leq m, \\
& \left(e_{m+2}, z\right) & \mapsto e_{1}, &
\end{array}
$$

being zero the remaining products among the bases elements in the action.

In this section we are going to study the Leibniz algebras $(L,[\cdot, \cdot])$ satisfying that $L / I \cong H_{2 m+1}$ and where the $H_{2 m+1}$-module $I$ is isomorphic to the minimal faithful representation $(I, \phi)$. From the above, $\operatorname{dim} L=$ $3 m+3$ and $\left\{x_{1}, x_{2}, \ldots, x_{m}, y_{1}, y_{2}, \ldots, y_{m}, z, e_{1}, e_{2}, \ldots, e_{m+2}\right\}$ is a basis of $L$. We also have

$$
\begin{array}{ll}
{\left[e_{i+1}, x_{i}\right]=e_{1},} & 1 \leq i \leq m \\
{\left[e_{m+2}, y_{i}\right]=e_{i+1},} & 1 \leq i \leq m \\
{\left[e_{m+2}, z\right]=e_{1}}
\end{array}
$$

Theorem 3. Let $L$ be a Leibniz algebra such that $L / I \cong H_{2 m+1}(m \neq 1)$ and $I$ is the $L / I$-module with the minimal faithful representation given by Equation (10). Then $L$ admits a basis

$$
\left\{x_{1}, x_{2}, \ldots, x_{m}, y_{1}, y_{2}, \ldots, y_{m}, z, e_{1}, e_{2}, \ldots, e_{m+2}\right\}
$$

in such a way that the multiplications table on this basis has the form

$$
\begin{array}{ll}
{\left[e_{i+1}, x_{i}\right]=e_{1},} & {\left[e_{m+2}, y_{i}\right]=e_{i+1},} \\
{\left[e_{m+2}, z\right]=e_{1},} & {\left[x_{i}, x_{j}\right]=\sum_{s=1}^{m+1} \alpha_{i, j}^{s} e_{s}} \\
{\left[x_{i}, y_{j}\right]=\gamma_{i, j} e_{1}, i \neq j,} & {\left[x_{i}, y_{i}\right]=-z+\delta_{i} e_{1}+\tau e_{2}+\sum_{s=2}^{m} \nu_{1, s}^{2} e_{s+1}} \\
{\left[y_{i}, y_{j}\right]=\beta_{i, j} e_{1},} & {\left[y_{1}, x_{1}\right]=z,} \\
{\left[y_{i}, x_{j}\right]=\sum_{s=1}^{m+1} \nu_{i, j}^{s} e_{s}, i \neq j,} & {\left[y_{i}, x_{i}\right]=z+\left(\nu_{i, 1}^{i+1}-\tau\right) e_{2}+\varepsilon_{i}^{i+1} e_{i+1}+\sum_{s=2}^{m}\left(\nu_{i, s}^{i+1}-\nu_{1, s}^{2}\right) e_{s+1}, i \neq 1,} \\
{\left[z, x_{1}\right]=\tau e_{1},} & {\left[z, x_{i}\right]=\nu_{1, i}^{2} e_{1}, i \neq 1,}
\end{array}
$$

for $1 \leq i, j \leq m$, where any $\alpha_{p, q}^{r}, \gamma_{p, q}, \delta_{p}, \tau, \nu_{p, q}^{r}, \beta_{p, q}, \varepsilon_{p}^{r} \in \mathbb{F}$ and where the omitted products are equal to zero.

Proof. We consider the following products:

$$
\left[y_{i}, x_{i}\right]=z+\sum_{k=1}^{m+2} \varepsilon_{i}^{k} e_{k}, \quad 1 \leq i \leq m
$$

Putting $z^{\prime}=z+\sum_{k=1}^{m+2} \varepsilon_{1}^{k} e_{k}$ we can assume $\left[y_{1}, x_{1}\right]=z$. Thus, we have

$$
\begin{array}{lll}
{\left[e_{i+1}, x_{i}\right]=e_{1},} & {\left[e_{m+2}, y_{i}\right]=e_{i+1},} & {\left[e_{m+2}, z\right]=e_{1},} \\
{\left[x_{i}, x_{j}\right]=\sum_{k=1}^{m+2} \alpha_{i, j}^{k} e_{k},} & {\left[x_{i}, y_{j}\right]=\sum_{k=1}^{m+2} \gamma_{i, j}^{k} e_{k}, \quad i \neq j} & {\left[x_{i}, y_{i}\right]=-z+\sum_{k=1}^{m+2} \delta_{i}^{k} e_{k},} \\
{\left[x_{i}, z\right]=\sum_{k=1}^{m+2} \eta_{i}^{k} e_{k},} & {\left[y_{i}, y_{j}\right]=\sum_{k=1}^{m+2} \beta_{i, j}^{k} e_{k},} & {\left[y_{i}, x_{j}\right]=\sum_{k=1}^{m+2} \nu_{i, j}^{k} e_{k}, i \neq j,} \\
{\left[y_{i}, z\right]=\sum_{k=1}^{m+2} \theta_{i}^{k} e_{k},} & {\left[y_{1}, x_{1}\right]=z,} & {\left[y_{i}, x_{i}\right]=z+\sum_{k=1}^{m+2} \varepsilon_{i}^{k} e_{k}, i \neq 1,} \\
{\left[z, x_{i}\right]=\sum_{k=1}^{m+2} \tau_{i}^{k} e_{k},} & {\left[z, y_{i}\right]=\sum_{k=1}^{m+2} \lambda_{i}^{k} e_{k},} & {[z, z]=\sum_{k=1}^{m+2} \mu^{k} e_{k},}
\end{array}
$$

with $1 \leq i, j \leq m$.

We compute all Leibniz identities using the software Mathematica and we get the following restrictions:

Leibniz identity


## Constraint

From here,

$$
\begin{array}{ll}
{\left[e_{i+1}, x_{i}\right]=e_{1},} & 1 \leq i \leq m \\
{\left[e_{m+2}, y_{i}\right]=e_{i+1},} & 1 \leq i \leq m \\
{\left[e_{m+2}, z\right]=e_{1},} & 1 \leq i, j \leq m, \\
{\left[x_{i}, x_{j}\right]=\sum_{s=1}^{m+1} \alpha_{i, j}^{s} e_{s},} & 1 \leq i, j \leq m, \\
{\left[y_{i}, y_{j}\right]=\beta_{i, j}^{1} e_{1}+\theta_{i}^{1} e_{j+1},} & 1 \leq
\end{array}
$$

$$
\begin{array}{ll}
{\left[x_{i}, y_{j}\right]=\gamma_{i, j}^{1} e_{1}+\eta_{i}^{1} e_{j+1},} & 1 \leq i, j \leq m, i \neq j \\
{\left[x_{1}, y_{1}\right]=-z+\delta_{1}^{1} e_{1}+\left(\eta_{1}^{1}+\tau_{1}^{1}\right) e_{2}+\sum_{s=2}^{m} \nu_{1, s}^{2} e_{s+1},} & 1 \leq i \leq m, \\
{\left[x_{i}, y_{i}\right]=-z+\delta_{i}^{1} e_{1}+\tau_{1}^{1} e_{2}+\left(\eta_{i}^{1}+\nu_{1, i}^{2}\right) e_{i+1}+\sum_{s=2}^{m} \nu_{s \neq i}^{2} e_{1, s} e_{s+1},} & 2 \leq i \leq m, \\
{\left[y_{1}, x_{1}\right]=z,} & 1 \leq i, j \leq m, i \neq j \\
{\left[y_{i}, x_{i}\right]=z+\varepsilon_{i}^{1} e_{1}+\left(\nu_{i, 1}^{i+1}-\tau_{1}^{1}\right) e_{2}+\varepsilon_{i}^{i+1} e_{i+1}+\sum_{s=2}^{m}\left(\nu_{i, s}^{i+1}-\nu_{1, s}^{2}\right) e_{s+1},} & 2 \leq i \leq m, \\
{\left[y_{i}, x_{j}\right]=\sum_{s=1}^{m+1} \nu_{i, j}^{s} e_{s},} & 1 \leq i \leq m, \\
{\left[x_{i}, z\right]=\eta_{i}^{1} e_{1},} & 1 \leq i \leq m, \\
{\left[y_{i}, z\right]=\theta_{i}^{1} e_{1},} & 2 \leq i \leq m, \\
{\left[z, x_{1}\right]=\tau_{1}^{1} e_{1},} & 1 \leq 2
\end{array}
$$

with the following restrictions

$$
\begin{array}{ll}
\alpha_{i, j}^{k+1}=\alpha_{i, k}^{j+1}, & 1 \leq i, j, k \leq m \\
\nu_{i, j}^{k+1}=\nu_{i, k}^{j+1}, & 1 \leq i, j, k \leq m, j \neq i \neq k
\end{array}
$$

Only rest to make the next change of basis

$$
\begin{cases}x_{i}^{\prime}=x_{i}-\eta_{i}^{1} e_{m+2}, & 1 \leq i \leq m \\ y_{1}^{\prime}=y_{1}-\theta_{1}^{1} e_{m+2} & \\ y_{j}^{\prime}=y_{j}-\varepsilon_{j}^{1} e_{j+1}-\theta_{j}^{1} e_{m+2}, & 2 \leq j \leq m\end{cases}
$$

and we obtain the family of the theorem (renaming the parameters).
4.2. Particular case: Classification of Leibniz algebras when $m=1$. In this subsection we classify the Leibniz algebras such that $L / I \cong H_{3}$ and $I$ is the $L / I$-module with the minimal faithful representation given by Equation (10). Let us fix $\left\{x, y, z, e_{1}, e_{2}, e_{3}\right\}$ as basis of $L$. All computations have been made by using the software Mathematica.

We have the following products:

$$
\begin{array}{lll}
{\left[e_{2}, x\right]=e_{1},} & {\left[e_{3}, y\right]=e_{2},} & {\left[e_{3}, z\right]=e_{1},} \\
{[x, x]=\alpha_{1} e_{1}+\alpha_{2} e_{2}+\alpha_{3} e_{3},} & {[x, y]=-z+\delta_{1} e_{1}+\delta_{2} e_{2}+\delta_{3} e_{3},} & {[x, z]=\eta_{1} e_{1}+\eta_{2} e_{2}+\eta_{3} e_{3}} \\
{[y, y]=\beta_{1} e_{1}+\beta_{2} e_{2}+\beta_{3} e_{3},} & {[y, x]=z,} & {[y, z]=\theta_{1} e_{1}+\theta_{2} e_{2}+\theta_{3} e_{3}} \\
{[z, x]=\tau_{1} e_{1}+\tau_{2} e_{2}+\tau_{3} e_{3},} & {[z, y]=\lambda_{1} e_{1}+\lambda_{2} e_{2}+\lambda_{3} e_{3},} & {[z, z]=\mu_{1} e_{1}+\mu_{2} e_{2}+\mu_{3} e_{3}}
\end{array}
$$

The Leibniz identity on the following triples imposes further constraints on the products.

| Leibniz identity | Constraint |
| :--- | :--- |
|  | $\Rightarrow \quad-\eta_{1}=\tau_{1}-\delta_{2}, \alpha_{3}-\eta_{2}=\tau_{2}, \quad-\eta_{3}=\tau_{3}$, |
| $\{x, x, y\}$ | $\Rightarrow \alpha_{3}=\eta_{2}$, |
| $\{x, x, z\}$ | $\Rightarrow \mu_{1}=\delta_{3}, \mu_{2}=-\eta_{3}, \quad \mu_{3}=0$, |
| $\{x, y, z\}$ | $\Rightarrow \beta_{3}=\theta_{3}=0$, |
| $\{y, y, z\}$ | $\Rightarrow-\theta_{1}=\lambda_{1}-\beta_{2},-\theta_{2}=\lambda_{2},-\theta_{3}=\lambda_{3}$, |
| $\{y, x, y\}$ | $\Rightarrow \mu_{1}=\theta_{2}, \quad \mu_{2}=0$, |
| $\{y, x, z\}$ | $\Rightarrow \mu_{1}=\lambda_{2}, \quad \mu_{2}=\tau_{3}$, |
| $\{z, x, y\}$ | $\Rightarrow \mu_{2}=\tau_{3}$, |
| $\{z, x, z\}$ | $\Rightarrow \lambda_{3}=0$, |
| $\{z, y, z\}$ | $\Rightarrow \eta_{2}=\alpha_{3}$, |
| $\{x, z, x\}$ | $\Rightarrow \mu_{1}=\delta_{3}, \quad \mu_{2}=-\eta_{3}$. |

Thus, we get the following family of algebras, $L\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \beta_{1}, \beta_{2}, \delta_{1}, \delta_{2}, \eta_{1}, \theta_{1}\right)$ :

$$
\left\{\begin{array}{lll}
{\left[e_{2}, x\right]=e_{1},} & {\left[e_{3}, y\right]=e_{2},} & {\left[e_{3}, z\right]=e_{1}} \\
{[x, x]=\alpha_{1} e_{1}+\alpha_{2} e_{2}+\alpha_{3} e_{3},} & {[x, y]=-z+\delta_{1} e_{1}+\delta_{2} e_{2},} & {[x, z]=\eta_{1} e_{1}+\alpha_{3} e_{2}} \\
{[y, y]=\beta_{1} e_{1}+\beta_{2} e_{2},} & {[y, x]=z,} & {[y, z]=\theta_{1} e_{1}} \\
{[z, x]=\left(\delta_{2}-\eta_{1}\right) e_{1}-2 \alpha_{3} e_{2},} & {[z, y]=\left(\beta_{2}-\theta_{1}\right) e_{1}} &
\end{array}\right.
$$

Theorem 4. Let $L$ be a Leibniz algebra such that $L / I \cong H_{3}$ and $I$ is the $L / I$-module with the minimal faithful representation given by Equation (10). Then $L$ is isomorphic to one of the following pairwise non-isomorphic algebras:

$$
\begin{array}{lll}
L(0,1,0,1,0,0,0,1, \lambda), \lambda \in \mathbb{F}, & L(0,1,0,1,0,0,0,0,1), & L(0,1,0,1,0,0,0,0,0), \\
L(0,1,0,0,0,0,0,1, \lambda), \lambda \in \mathbb{F}, & L(0,1,0,0,0,0,0,0,1), & L(0,1,0,0,0,0,0,0,0), \\
L(0,0,0,1,0,0,0,1,1), & L(0,0,0,1,0,0,0,1,0), & L(0,0,0,1,0,0,0,0,1), \\
L(0,0,0,1,0,0,0,0,0), & L(0,0,0,0,0,0,0,1,1), & L(0,0,0,0,0,0,0,1,0), \\
L(0,0,0,0,0,0,0,0,1), & L(0,0,0,0,0,0,0,0,0), & L(0,0,1,1,0,0,0,1, \lambda), \lambda \in \mathbb{F}, \\
L(0,0,1,1,0,0,0,0,1), & L(0,0,1,1,0,0,0,0,0), & L(0,0,1,0,0,0,0,1,1), \\
L(0,0,1,0,0,0,0,1,0), & L(0,0,1,0,0,0,0,0,1), & L(0,0,1,0,0,0,0,0,0)
\end{array}
$$

Proof. We can distinguish two cases:
Case 1: $e_{3} \in[L, L]$. Then $\alpha_{3}=0$.
Applying the general change of basis generators:
$x^{\prime}=A_{1} x+A_{2} y+A_{3} z+\sum_{k=1}^{3} P_{i} e_{i}, y^{\prime}=B_{1} x+B_{2} y+B_{3} z+\sum_{k=1}^{3} Q_{i} e_{i}, e_{3}^{\prime}=C_{1} x+C_{2} y+C_{3} z+\sum_{k=1}^{3} R_{i} e_{i}$ we derive the expressions of the new parameters in the new basis:

$$
\begin{array}{ll}
\alpha_{1}^{\prime}=\frac{\alpha_{1} A_{1}^{2} B_{2}-\alpha_{2} A_{1}^{2} B_{3}+\delta_{2} A_{1} A_{3} B_{2}+A_{1} B_{2} P_{2}+A_{3} B_{2} P_{3}}{A_{1} B_{2}^{2}}, & \alpha_{2}^{\prime}=\frac{\alpha_{2} A_{1}^{2}}{B_{2} R_{3}} \\
\beta_{1}^{\prime}=\frac{\beta_{1} B_{2}}{A_{1} R_{3}}, & \beta_{2}^{\prime}=\frac{\beta_{2} B_{2}+Q_{3}}{R_{3}} \\
\delta_{1}^{\prime}=\frac{\beta_{2} A_{3} B_{2}+\delta_{1} A_{1} B_{2}+A_{1} Q_{2}+A_{3} Q_{3}}{A_{1} B_{2} R_{3}}, & \delta_{2}^{\prime}=\frac{\delta_{2} A_{1}+P_{3}}{R_{3}} \\
\eta_{1}^{\prime}=\frac{\eta_{1} A_{1}+P_{3}}{R_{3}}, & \theta_{1}^{\prime}=\frac{\theta_{1} B_{2}+Q_{3}}{R_{3}}
\end{array}
$$

and the following restrictions:

$$
\left\{\begin{array}{l}
C_{1}=C_{2}=C_{3}=B_{1}=A_{2}=0 \\
R_{5}=-\frac{A_{3} R_{3}}{A_{1}} \\
A_{1} B_{2} R_{3} \neq 0
\end{array}\right.
$$

We set

$$
\begin{array}{ll}
P_{3}=-\delta_{2} A_{1} & \Rightarrow \delta_{2}^{\prime}=0 \\
Q_{3}=-\beta_{2} B_{2} & \Rightarrow \beta_{2}^{\prime}=0 \\
Q_{2}=-\delta_{1} B_{2} & \Rightarrow \delta_{1}^{\prime}=0 \\
P_{2}=-\frac{\left(\alpha_{1} B_{2}-\alpha_{2} B_{3}\right) A_{1}}{B_{2}} & \Rightarrow \alpha_{1}^{\prime}=0
\end{array}
$$

then we get

$$
\begin{array}{lll}
{\left[e_{2}, x\right]=e_{1},} & {\left[e_{3}, y\right]=e_{2},} & {\left[e_{3}, z\right]=e_{1},} \\
{[x, x]=\alpha_{2}^{\prime} e_{2},} & {[x, y]=-z,} & {[x, z]=\eta_{1}^{\prime} e_{1}} \\
{[y, y]=\beta_{1}^{\prime} e_{1},} & {[y, x]=z,} & {[y, z]=\theta_{1}^{\prime} e_{1},} \\
{[z, x]=-\eta_{1}^{\prime} e_{1},} & {[z, y]=-\theta_{1}^{\prime} e_{1},} &
\end{array}
$$

where

$$
\alpha_{2}^{\prime}=\frac{\alpha_{2} A_{1}^{2}}{B_{2} R_{3}}, \quad \beta_{1}^{\prime}=\frac{\beta_{1} B_{2}}{A_{1} R_{3}}, \quad \eta_{1}^{\prime}=\frac{\left(\eta_{1}-\delta_{2}\right) A_{1}}{R_{3}}, \quad \theta_{1}^{\prime}=\frac{\left(\theta_{1}-\beta_{2}\right) B_{2}}{R_{3}}
$$

We observe that the nullities of $\alpha_{2}, \beta_{1}, \eta_{1}, \theta_{1}$ are invariant. Thus, we can distinguish the following nonisomorphic cases. An appropriate choice of the parameter values $\left(A_{1}, B_{2}\right.$ and $\left.R_{3}\right)$ allows us to obtain the following algebras or families of algebras.

| Case | Algebra |
| :--- | :--- |
| $\alpha_{2} \neq 0, \beta_{1} \neq 0, \eta_{1} \neq 0$, | $L(0,1,0,1,0,0,0,1, \lambda), \lambda \in \mathbb{F}$, |
| $\alpha_{2} \neq 0, \beta_{1} \neq 0, \eta_{1}=0, \theta_{1} \neq 0$, | $L(0,1,0,1,0,0,0,0,1)$, |
| $\alpha_{2} \neq 0, \beta_{1} \neq 0, \eta_{1}=0, \theta_{1}=0$, | $L(0,1,0,1,0,0,0,0,0)$, |
| $\alpha_{2} \neq 0, \beta_{1}=0, \eta_{1} \neq 0$, | $L(0,1,0,0,0,0,0,1, \lambda), \lambda \in \mathbb{F}$, |
| $\alpha_{2} \neq 0, \beta_{1}=0, \eta_{1}=0, \theta_{1} \neq 0$, | $L(0,1,0,0,0,0,0,0,1)$, |
| $\alpha_{2} \neq 0, \beta_{1}=0, \eta_{1}=0, \theta_{1}=0$, | $L(0,1,0,0,0,0,0,0,0)$, |
| $\alpha_{2}=0, \beta_{1} \neq 0, \eta_{1} \neq 0, \theta_{1} \neq 0$, | $L(0,0,0,1,0,0,0,1,1)$, |
| $\alpha_{2}=0, \beta_{1} \neq 0, \eta_{1} \neq 0, \theta_{1}=0$, | $L(0,0,0,1,0,0,0,1,0)$, |
| $\alpha_{2}=0, \beta_{1} \neq 0, \eta_{1}=0, \theta_{1} \neq 0$, | $L(0,0,0,1,0,0,0,0,1)$, |
| $\alpha_{2}=0, \beta_{1} \neq 0, \eta_{1}=0, \theta_{1}=0$, | $L(0,0,0,1,0,0,0,0,0)$, |
| $\alpha_{2}=0, \beta_{1}=0, \eta_{1} \neq 0, \theta_{1} \neq 0$, | $L(0,0,0,0,0,0,0,1,1)$, |
| $\alpha_{2}=0, \beta_{1}=0, \eta_{1} \neq 0, \theta_{1}=0$, | $L(0,0,0,0,0,0,0,1,0)$, |
| $\alpha_{2}=0, \beta_{1}=0, \eta_{1}=0, \theta_{1} \neq 0$, | $L(0,0,0,0,0,0,0,0,1)$, |
| $\alpha_{2}=0, \beta_{1}=0, \eta_{1}=0, \theta_{1}=0$, | $L(0,0,0,0,0,0,0,0,0)$. |

Case 2: $e_{3} \notin[L, L]$. Then $\alpha_{3} \neq 0$. Making the following change of basis in $L\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \beta_{1}, \beta_{2}, \delta_{1}, \delta_{2}, \eta_{1}, \theta_{1}\right)$

$$
\left\{\begin{array}{l}
e_{3}^{\prime}=\alpha_{1} e_{1}+\alpha_{2} e_{2}+\alpha_{3} e_{3} \\
e_{2}^{\prime}=\alpha_{3} e_{2} \\
e_{1}^{\prime}=\alpha_{3} e_{1}
\end{array}\right.
$$

we obtain $L\left(0,0,1, \beta_{1}, \beta_{2}, \delta_{1}, \delta_{2}, \eta_{1}, \theta_{1}\right)$ :

$$
\left\{\begin{array}{lll}
{\left[e_{2}, x\right]=e_{1},} & {\left[e_{3}, y\right]=e_{2},} & {\left[e_{3}, z\right]=e_{1}} \\
{[x, x]=e_{3},} & {[x, y]=-z+\delta_{1} e_{1}+\delta_{2} e_{2},} & {[x, z]=\eta_{1} e_{1}+e_{2},} \\
{[y, y]=\beta_{1} e_{1}+\beta_{2} e_{2},} & {[y, x]=z,} & {[y, z]=\theta_{1} e_{1}} \\
{[z, x]=\left(\delta_{2}-\eta_{1}\right) e_{1}-2 e_{2},} & {[z, y]=\left(\beta_{2}-\theta_{1}\right) e_{1} .} &
\end{array}\right.
$$

Analogously to the previous case, by making the general change of basis of generators

$$
x^{\prime}=A_{1} x+A_{2} y+A_{3} z+\sum_{k=1}^{3} P_{i} e_{i}, \quad y^{\prime}=B_{1} x+B_{2} y+B_{3} z+\sum_{k=1}^{3} Q_{i} e_{i}
$$

we derive the expressions of the new parameters in the new basis:

$$
\begin{array}{ll}
\beta_{1}^{\prime}=\frac{\beta_{1} B_{2}}{A_{1}^{3}}, & \beta_{2}^{\prime}=\frac{\beta_{2} B_{2}+Q_{3}}{A_{1}^{2}}, \\
\delta_{1}^{\prime}=\frac{\beta_{2} A_{3} B_{2}^{2}+A_{1} B_{3}^{2}+\delta_{1} A_{1} B_{2}^{2}+A_{1} B_{2} Q_{2}+A_{3} B_{2} Q_{3}}{A_{1}^{3} B_{2}^{2}}, & \delta_{2}^{\prime}=\frac{-A_{1} B_{3}+\delta_{2} A_{1} B_{2}+B_{2} P_{3}}{A_{1}^{2} B_{2}} \\
\eta_{1}^{\prime}=\frac{-A_{1} B_{3}+\eta_{1} A_{1} B_{2}+B_{2} P_{3}}{A_{1}^{2} B_{2}}, & \theta_{1}^{\prime}=\frac{\theta_{1} B_{2}+Q_{3}}{A_{1}^{2}},
\end{array}
$$

with the restriction:

$$
\left\{\begin{array}{l}
A_{2}=B_{1}=0 \\
A_{1} B_{2} \neq 0
\end{array}\right.
$$

By putting

$$
\begin{array}{lll}
P_{3}=\frac{A_{1}\left(B_{3}-\delta_{2} B_{2}\right)}{B_{2}} & \Rightarrow & \delta_{2}^{\prime}=0 \\
Q_{3}=-\beta_{2} B_{2} & \Rightarrow & \beta_{2}^{\prime}=0 \\
Q_{2}=-\frac{B_{3}^{2}+\delta_{1} B_{2}^{2}}{B_{2}} & \Rightarrow & \delta_{1}^{\prime}=0
\end{array}
$$

we deduce

$$
\begin{array}{lll}
{\left[e_{2}, x\right]=e_{1},} & {\left[e_{3}, y\right]=e_{2},} & {\left[e_{3}, z\right]=e_{1},} \\
{[x, x]=e_{3},} & {[x, y]=-z,} & {[x, z]=\eta_{1}^{\prime} e_{1}+e_{2}} \\
{[y, y]=\beta_{1}^{\prime} e_{1},} & {[y, x]=z,} & {[y, z]=\theta_{1}^{\prime} e_{1}} \\
{[z, x]=-\eta_{1}^{\prime} e_{1}-2 e_{2},} & {[z, y]=-\theta_{1}^{\prime} e_{1},} &
\end{array}
$$

where

$$
\beta_{1}^{\prime}=\frac{\beta_{1} B_{2}}{A_{1}^{3}}, \quad \eta_{1}^{\prime}=\frac{\eta_{1}-\delta_{2}}{A_{1}}, \quad \theta_{1}^{\prime}=\frac{\left(\theta_{1}-\beta_{2}\right) B_{2}}{A_{1}^{2}} .
$$

We observe that the nullities of $\beta_{1}, \eta_{1}, \theta_{1}$ are invariant. Thus, we can distinguish the following nonisomorphic cases. An appropriate choice of the parameter values ( $A_{1}$ and $B_{2}$ ) allows us to obtain the following algebras or families of algebras.

| Case | Algebra |
| :--- | :--- |
| $\beta_{1} \neq 0, \eta_{1} \neq 0$, | $L(0,0,1,1,0,0,0,1, \lambda), \lambda \in \mathbb{F}$, |
| $\beta_{1} \neq 0, \eta_{1}=0, \theta \neq 0$, | $L(0,0,1,1,0,0,0,0,1)$, |
| $\beta_{1} \neq 0, \eta_{1}=0, \theta=0$, | $L(0,0,1,1,0,0,0,0,0)$, |
| $\beta_{1}=0, \eta_{1} \neq 0, \theta \neq 0$, | $L(0,0,1,0,0,0,0,1,1)$, |
| $\beta_{1}=0, \eta_{1} \neq 0, \theta=0$, | $L(0,0,1,0,0,0,0,1,0)$, |
| $\beta_{1}=0, \eta_{1}=0, \theta \neq 0$, | $L(0,0,1,0,0,0,0,0,1)$, |
| $\beta_{1}=0, \eta_{1}=0, \theta=0$, | $L(0,0,1,0,0,0,0,0,0)$. |

The proof is complete.

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