ON THE PROPERTIES OF SPECIAL FUNCTIONS ON THE LINEAR-TYPE LATTICES

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ABSTRACT. We present a general theory for studying the difference analogues of special functions of hypergeometric type on the linear-type lattices, i.e., the solutions of the second order linear difference equation of hypergeometric type on a special kind of lattices: the linear type lattices. In particular, using the integral representation of the solutions we obtain several difference-recurrence relations for such functions. Finally, applications to q-classical polynomials are given.

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1. Introduction

The study of the so-called q-special functions has known an increasing interest in the last years due its connection with several problems in mathematics and mathematical-physics (see e.g. [3, 6, 8, 13, 17]). A systematic study starting from the second order linear difference equation that such functions satisfy was started by Nikiforov and Uvarov in 1983 and further developed by Atakishiyev and Suslov (for a very nice reviews see e.g. [7, 13, 16]). Of particular interest is the so-called q-classical polynomials (see e.g. [5]) introduced by Hahn in 1949 which are polynomials on the lattice q^s .

Our main aim in this paper is to present a constructive approach for generating recurrence relations and ladder-type operators for the difference analogues of special functions of hypergeometric type on the linear-type lattices. Here we will focus our attention on functions defined on the q-linear lattice (for the linear lattice x(s) = s see [4] and references therein, and for the continuous case see e.g. [18]). Therefore we will complete the work started in [16] where few recurrence relations where obtained. In fact we will prove, by using the q-analoge of the technique introduced in [4] for the discrete case (uniform lattice), that the solutions (not only the polynomial ones) of the difference equation on the q-linear lattice $x(s) = c_1 q^s + c_2$ satisfy a very general recurrent-difference relation from where several well known relations (such as the three-term recurrence relation and the ladder-type relations) follow.

The structure of the paper is as follows: In section 2 the needed results and notations from the q-special function theory are introduced. In sections 3 and 4 the general theorems for obtaining recurrences relations are presented. In section 5 the special case of classical q-polynomials are considered in details and some examples are worked out in details.

2. Some preliminar results

Here we collect the basic background [1, 13, 16] on q-hypergeometric functions needed in the rest of the work.

The hypergeometric functions on the non-uniform lattice x(s) are the solutions of the second order linear difference equation of hypergeometric type on non-uniform lattices

$$\sigma(s) \frac{\Delta}{\Delta x(s - \frac{1}{2})} \left[\frac{\nabla y(s)}{\nabla x(s)} \right] + \tau(s) \frac{\Delta y(s)}{\Delta x(s)} + \lambda y(s) = 0,$$

$$\sigma(s) = \widetilde{\sigma}(x(s)) - \frac{1}{2} \widetilde{\tau}(x(s)) \Delta x \left(s - \frac{1}{2} \right), \quad \tau(s) = \widetilde{\tau}(x(s)),$$
(1)

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where $\Delta y(s) := y(s+1) - y(s)$, $\nabla y(s) := y(s) - y(s-1)$, are the forward and backward difference operators, respectively; $\tilde{\sigma}(x(s))$ and $\tilde{\tau}(x(s))$ are polynomials in x(s) of degree at most 2 and 1, respectively, and λ is a constant. Here we will deal with the linear and q-linear lattices, i.e., lattices of the form

$$x(s) = c_1 s + c_2$$
 or $x(s) = c_1(q)q^s + c_2(q)$, (2)

respectively, with $c_1 \neq 0$ and $c_1(q) \neq 0$.

We will define the k-order difference derivative of a solution y(s) of (1) by

$$y^{(k)}(s) := \Delta^{(k)}[y(s)] = \frac{\Delta}{\Delta x_{k-1}(s)} \frac{\Delta}{\Delta x_{k-2}(s)} \dots \frac{\Delta}{\Delta x(s)}[y(s)],$$

where $x_{\nu}(s) = x(s + \frac{\nu}{2})$. It is known [13] that $y^{(k)}(s)$ also satisfy a difference equation of the same type. Moreover, for the solutions of the difference equation (1) the following theorem holds

Theorem 2.1. [12, 16] The difference equation (1) has a particular solution of the form

$$y_{\nu}(z) = \frac{C_{\nu}}{\rho(z)} \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\nu+1)}},$$
(3)

if the condition

$$\frac{\sigma(s)\rho_{\nu}(s)\nabla x_{\nu+1}(s)}{[x_{\nu-1}(s)-x_{\nu-1}(z+1)]^{(\nu+1)}}\bigg|_a^b=0,$$

is satisfied, and of the form

$$y_{\nu}(z) = \frac{C_{\nu}}{\rho(z)} \int_{C} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\nu+1)}} ds, \tag{4}$$

if the condition

$$\int_{C} \Delta_{s} \frac{\sigma(s)\rho_{\nu}(s)\nabla x_{\nu+1}(s)}{[x_{\nu-1}(s) - x_{\nu-1}(z+1)]^{(\nu+1)}} = 0,$$
(5)

is satisfied. Here C is a contour in the complex plane, C_{ν} is a constant, $\rho(s)$ and $\rho_{\nu}(s)$ are the solution of the Pearson-type equations

$$\frac{\rho(s+1)}{\rho(s)} = \frac{\sigma(s) + \tau(s)\Delta x(s - \frac{1}{2})}{\sigma(s+1)} = \frac{\phi(s)}{\sigma(s+1)},$$

$$\frac{\rho_{\nu}(s+1)}{\rho_{\nu}(s)} = \frac{\sigma(s) + \tau_{\nu}(s)\Delta x_{\nu}(s - \frac{1}{2})}{\sigma(s+1)} = \frac{\phi_{\nu}(s)}{\sigma(s+1)},$$
(6)

where

$$\tau_{\nu}(s) = \frac{\sigma(s+\nu) - \sigma(s) + \tau(s+\nu)\Delta x(s+\nu - \frac{1}{2})}{\Delta x_{\nu-1}(s)},$$
(7)

 ν is the root of the equation

$$\lambda_{\nu} + [\nu]_q \left\{ \alpha_q(\nu - 1)\widetilde{\tau}' + [\nu - 1]_q \frac{\widetilde{\sigma}''}{2} \right\} = 0, \tag{8}$$

and $[\nu]_q$ and $\alpha_q(\nu)$ are the q-numbers

$$[\nu]_q = \frac{q^{\nu/2} - q^{-\nu/2}}{q^{1/2} - q^{-1/2}}, \qquad \alpha_q(\nu) = \frac{q^{\nu/2} + q^{-\nu/2}}{2}, \quad \forall \nu \in \mathbb{C},$$
(9)

respectively. The generalized powers $[x_k(s) - x_k(z)]^{(\nu)}$ are defined by

$$[x_k(s) - x_k(z)]^{(\nu)} = (q-1)^{\nu} c_1^{\nu} q^{\nu(k-\nu+1)/2} q^{\nu z} \frac{\Gamma_q(s-z+\nu)}{\Gamma_q(s-z)}, \quad \nu \in \mathbb{R},$$
 (10)

for the q-linear (exponential) lattice $x(s) = c_1q^s + c_2$ and

$$[x_k(s) - x_k(z)]^{(\nu)} = c_1^{\nu} \frac{\Gamma(s - z + \mu)}{\Gamma(s - z)}, \quad \nu \in \mathbb{R},$$

for the linear lattice $x(s) = c_1 s + c_2$, respectively. For the definitions of the Gamma and the q-Gamma functions see, for instance, [6].

Remark 2.2. For the special case when $\nu \in \mathbb{N}$, the generalized powers become

$$[x_k(s) - x_k(z)]^{(n)} = (-1)^n c_1^n q^{-n(n-1)/2} q^{n(z+k/2)} (q^{s-z}; q)_n,$$

$$[x_k(s) - x_k(z)]^{(n)} = c_1^n (s-z)_n,$$

for q-linear and linear lattices, respectively.

We will need the following straightforward proposition which proof we omit here (see e.g. [1, 16])

Proposition 2.3. Let μ and ν be complex numbers and m and k be positive integers with $m \ge k$. For the q-linear lattice $x(s) = c_1 q^s + c_2$ we have

(1)
$$\frac{\left[x_{\mu}(s) - x_{\mu}(z)\right]^{(m)}}{\left[x_{\nu}(s) - x_{\nu}(z)\right]^{(m)}} = q^{\frac{m(\mu - \nu)}{2}},$$

(2)
$$\frac{[x_{\mu}(s) - x_{\mu}(z)]^{(m)}}{[x_{\mu}(s) - x_{\mu}(z)]^{(k)}} = [x_{\mu}(s) - x_{\mu}(z-k)]^{(m-k)},$$

(3)
$$\frac{\left[x_{\mu}(s) - x_{\mu}(z)\right]^{(m)}}{\left[x_{\nu}(s) - x_{\nu}(z)\right]^{(k)}} = q^{\frac{k(\mu-\nu)}{2}} \left[x_{\mu}(s) - x_{\mu}(z-k)\right]^{(m-k)},$$

(4)
$$\frac{\left[x_{\mu}(s) - x_{\nu}(z)\right]^{(m+1)}}{\left[x_{\mu-1}(s+1) - x_{\mu-1}(z)\right]^{(m)}} = x_{\mu-m}(s) - x_{\mu-m}(z),$$
(5)
$$\frac{\left[x_{\mu}(s) - x_{\mu}(z)\right]^{(m+1)}}{\left[x_{\mu-1}(s) - x_{\mu-1}(z)\right]^{(m)}} = x_{\mu-m}(s+m) - x_{\mu-m}(z).$$

(5)
$$\frac{\left[x_{\mu}(s) - x_{\mu}(z)\right]^{(m+1)}}{\left[x_{\mu-1}(s) - x_{\mu-1}(z)\right]^{(m)}} = x_{\mu-m}(s+m) - x_{\mu-m}(z)$$

To obtain the result for the linear lattice one only has to put in the above formulas q=1.

3. The general recurrence relation in the linear-type lattices

In this section we will obtain several recurrence relations for the solutions (3) and (4) of the difference equation (1) in the linear-type lattices (2). Since the equation (1) is linear we can restrict ourselves to the canonical cases $x(s) = q^s$ and x(s) = s.

Let us define the functions¹

$$\Phi_{\nu,\mu}(z) = \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\mu+1)}}$$
(11)

and

$$\Phi_{\nu,\mu}(z) = \int_C \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\nu+1)}} ds.$$
 (12)

Notice that the functions y_{ν} and the functions $\Phi_{\nu,\mu}$ are related by the formula

$$y_{\nu}(z) = \frac{C_{\nu}}{\rho(z)} \Phi_{\nu,\nu}(z). \tag{13}$$

Lemma 3.1. For the functions $\Phi_{\nu,\mu}(z)$ the following relation holds

$$\nabla_z \, \Phi_{\nu,\mu}(z) = [\mu + 1]_a \nabla x_{\nu-\mu}(z) \Phi_{\nu,\mu+1}(z), \tag{14}$$

where $[t]_q$ denotes the symmetric q-numbers (9).

¹Obviously the functions (3) correspond to the functions (11), whereas the functions y_{ν} given by (4) correspond to those of (12).

Proof. We will prove it for the functions (11). The other case is analogous. Using (10), one gets

$$\nabla_{z} \Phi_{\nu,\mu}(z) = \sum_{s=a}^{b-1} \nabla_{z} \left(\frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\mu+1)}} \right)$$

$$= \sum_{s=a}^{b-1} \left(\frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\mu+1)}} - \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z-1)]^{(\mu+1)}} \right)$$

$$= \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z-1)]^{(\mu)}} \left(\frac{1}{x_{\nu}(s) - x_{\nu}(z)} - \frac{1}{x_{\nu}(s) - x_{\nu}(z-1-\mu)} \right)$$

$$= \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z-1)]^{(\mu)}} \frac{x_{\nu}(z) - x_{\nu}(z-1-\mu)}{(x_{\nu}(s) - x_{\nu}(z))(x_{\nu}(s) - x_{\nu}(z-1-\mu))}$$

$$= \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{[x_{\nu}(s) - x_{\nu}(z)]^{(\mu+2)}} (x_{\nu}(z) - x_{\nu}(z-1-\mu))$$

Since $x(s) - x(s-t) = [t]_q \nabla x \left(s - \frac{t-1}{2}\right)$ we then have

$$\nabla_z \, \Phi_{\nu,\mu}(z) = \sum_{s=a}^{b-1} \frac{\rho_{\nu}(s) \nabla x_{\nu+1}(s)}{\left[x_{\nu}(s) - x_{\nu}(z) \right]^{(\mu+2)}} [\mu + 1]_q \nabla x_{\nu} \left(z - \frac{\mu}{2} \right)$$
$$= [\mu + 1]_q \nabla x_{\nu-\mu}(z) \Phi_{\nu,\mu+1}(z)$$

which is (14).

From (14) follows that

$$\Delta_z \Phi_{\nu,\mu}(z) = [\mu + 1]_q \Delta x_{\nu-\mu}(z) \Phi_{\nu,\mu+1}(z+1).$$

Next we prove the following lemma that is the discrete analog of the Lemma in [14, page 14].

Lemma 3.2. Let x(z) be $x(z) = q^z$ or x(z) = z. Then, any three functions $\Phi_{\nu_i,\mu_i}(z)$, i = 1, 2, 3, are connected by a linear relation

$$\sum_{i=1}^{3} A_i(z) \Phi_{\nu_i, \mu_i}(z) = 0, \tag{15}$$

with non-zero at the same time polynomial coefficients on x(z), $A_i(z)$, provided that the differences $\nu_i - \nu_j$ and $\mu_i - \mu_j$, i, j = 1, 2, 3, are integers and that the following condition holds²

$$\frac{x^k(s)\sigma(s)\rho_{\nu_0}(s)}{[x_{\nu_0-1}(s)-x_{\nu_0-1}(z)]^{(\mu_0)}}\bigg|_{s=a}^{s=b} = 0, \quad k = 0, 1, 2, \dots,$$
(16)

when the functions Φ_{ν_i,μ_i} are given by (11) and

$$\int_{C} \Delta_{s} \frac{x^{k}(s)\sigma(s)\rho_{\nu_{0}}(s) ds}{[x_{\nu_{0}-1}(s) - x_{\nu_{0}-1}(z)]^{(\mu_{0})}} = 0, \quad k = 0, 1, 2, \dots,$$
(17)

when Φ_{ν_i,μ_i} are given by (12). Here ν_0 is the ν_i , i = 1, 2, 3, with the smallest real part and μ_0 is the μ_i , i = 1, 2, 3, with the largest real part.

Proof. Since in [4] we have proved the case when x(s) = s (the uniform lattice) we will restrict here to the case of the q-linear lattice $x(s) = c_1q^s + c_2$). Moreover, we will give the proof for the case of functions of the form (11), the other case is completely equivalent. Using the identity

$$\nabla x_{\nu_i+1}(s) = q^{\frac{\nu_i - \nu_0}{2}} \nabla x_{\nu_0+1}(s),$$

²In some cases this condition is equivalent to the condition $x(s)^k \sigma(s) \rho_{\nu_0}(s)|_{s=a}^{s=b} = 0, k=0,1,2,\ldots$

as well as (3) of Proposition 2.3, we have

$$\begin{split} \sum_{i=1}^{3} A_{i}(z) \Phi_{\nu_{i},\mu_{i}}(z) &= \sum_{i=1}^{3} A_{i}(z) \sum_{s=a}^{b-1} \frac{\rho_{\nu_{i}}(s) \nabla x_{\nu_{i}+1}(s)}{[x_{\nu_{i}}(s) - x_{\nu_{i}}(z)]^{(\mu_{i}+1)}} \\ &= \sum_{s=a}^{b-1} \sum_{i=1}^{3} A_{i}(z) \frac{\rho_{\nu_{i}}(s) \nabla x_{\nu_{i}+1}(s)}{[x_{\nu_{i}}(s) - x_{\nu_{i}}(z)]^{(\mu_{i}+1)}} = \sum_{s=a}^{b-1} \frac{1}{[x_{\nu_{0}}(s) - x_{\nu_{0}}(z)]^{(\mu_{0}+1)}} \times \\ &\left(\sum_{i=1}^{3} A_{i}(z) q^{\frac{(\mu_{i}+1)(\nu_{0}-\nu_{i})}{2}} [x_{\nu_{0}}(s) - x_{\nu_{0}}(z-\mu_{i}-1)]^{(\mu_{0}-\mu_{i})} \rho_{\nu_{i}}(s) \nabla x_{\nu_{i}+1}(s)\right) \\ &= \sum_{s=a}^{b-1} \frac{\rho_{\nu_{0}}(s) \nabla x_{\nu_{0}+1}(s)}{[x_{\nu_{0}}(s) - x_{\nu_{0}}(z)]^{(\mu_{0}+1)}} \times \\ &\left(\sum_{i=1}^{3} A_{i}(z) q^{\frac{\mu_{i}(\nu_{0}-\nu_{i})}{2}} [x_{\nu_{0}}(s) - x_{\nu_{0}}(z-\mu_{i}-1)]^{(\mu_{0}-\mu_{i})} \frac{\rho_{\nu_{i}}(s)}{\rho_{\nu_{0}}(s)}\right). \end{split}$$

Using the Pearson-type equation (6) we obtain

$$\rho_{\nu_i}(s) = \phi(s + \nu_0)\phi(s + \nu_0 + 1)\dots\phi(s + \nu_i - 1)\rho_{\nu_0}(s), \tag{18}$$

so

$$\sum_{i=1}^{3} A_i(z) \Phi_{\nu_i, \mu_i}(z) = \sum_{s=a}^{b-1} \frac{\rho_{\nu_0}(s) \nabla x_{\nu_0 + 1}(s)}{\left[x_{\nu_0}(s) - x_{\nu_0}(z)\right]^{(\mu_0 + 1)}} \Pi(s)$$

where

$$\Pi(s) = \sum_{i=1}^{3} A_i(z) q^{\frac{\mu_i(\nu_0 - \nu_i)}{2}} \left[x_{\nu_0}(s) - x_{\nu_0}(z - \mu_i - 1) \right]^{(\mu_0 - \mu_i)} \times \phi(s + \nu_0) \phi(s + \nu_0 + 1) \cdots \phi(s + \nu_i - 1) .$$
(19)

Let us show that there exists a polynomial Q(s) in x(s) (in general, $Q \equiv Q(z,s)$ is a function of z and s) such that

$$\frac{\rho_{\nu_0}(s)\nabla x_{\nu_0+1}(s)}{\left[x_{\nu_0}(s) - x_{\nu_0}(z)\right]^{(\mu_0+1)}}\Pi(s) = \Delta \left[\frac{\rho_{\nu_0}(s-1)}{\left[x_{\nu_0-1}(s) - x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s)\right]
= \Delta \left[\frac{\sigma(s)\rho_{\nu_0}(s)}{\left[x_{\nu_0-1}(s) - x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s)\right].$$
(20)

If such polynomial exists, then, taking the sum in s from s = a to b - 1 and using the boundary conditions (16) we obtain (15).

To prove the existence of the polynomial Q(s) in the variable x(s) in (20) we write

$$\frac{\sigma(s+1)\rho_{\nu_0}(s+1)}{\left[x_{\nu_0-1}(s+1)-x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s+1) - \frac{\sigma(s)\rho_{\nu_0}(s)}{\left[x_{\nu_0-1}(s)-x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s) = \frac{\rho_{\nu_0}(s)}{\left[x_{\nu_0}(s)-x_{\nu_0}(z)\right]^{(\mu_0+1)}}\left[\sigma(s+1)\frac{\rho_{\nu_0}(s+1)}{\rho_{\nu_0}(s)}\frac{\left[x_{\nu_0}(s)-x_{\nu_0}(z)\right]^{(\mu_0+1)}}{\left[x_{\nu_0-1}(s+1)-x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s+1) - \sigma(s)\frac{\left[x_{\nu_0}(s)-x_{\nu_0}(z)\right]^{(\mu_0+1)}}{\left[x_{\nu_0-1}(s)-x_{\nu_0-1}(z)\right]^{(\mu_0)}}Q(s)\right].$$

From (4) and (5) of Proposition 2.3, and using (6), the above expression becomes

$$\frac{\rho_{\nu_0}(s)}{\left[x_{\nu_0}(s)-x_{\nu_0}(z)\right]^{(\mu_0+1)}} \left\{\phi_{\nu_0}(s)\left[x_{\nu_0-\mu_0}(s)-x_{\nu_0-\mu_0}(z)\right]Q(s+1) - \sigma(s)\left[x_{\nu_0-\mu_0}(s+\mu_0)-x_{\nu_0-\mu_0}(z)\right]Q(s)\right\}.$$

Thus

$$(\sigma(s) + \tau_{\nu_0}(s)\nabla x_{\nu_0+1}(s)) \left[x_{\nu_0-\mu_0}(s) - x_{\nu_0-\mu_0}(z)\right] Q(s+1) -$$

$$\sigma(s) \left[x_{\nu_0-\mu_0}(s+\mu_0) - x_{\nu_0-\mu_0}(z)\right] Q(s) = \nabla x_{\nu_0+1}(s)\Pi(s).$$
(21)

Since $\nabla x_{\nu_0+1}(s)$ is a polynomial of degree one in x(s), $x_k(s)$ and $\tau_{\nu_0}(s)$ are polynomials of degree at most one in x(s), and $\sigma(s)$ is a polynomial of degree at most two in x(s), we conclude that the degree of Q(s) is, at least, two less than the degree of $\Pi(s)$, i.e., $\deg Q \ge \deg \Pi - 2$. Moreover, equating the coefficients of the powers of $x(s) = q^s$ on the two sides of the above equation (21), we find a system of linear equations in the coefficients of Q(s) and the coefficients $A_i(z)$ which have at least one unknown more then the number of equations. Notice that the coefficients of the unknowns are polynomials in q^z , so that after one coefficient is selected the remaining coefficients are rational functions of q^z , therefore after multiplying by the common denominator of the $A_i(z)$ we obtain the linear relation with polynomial coefficients on $x \equiv x(z) = q^z$. This completes the proof.

The above Lemma when $q \to 1$ and x(s) = s leads to the corresponding result on the uniform lattice x(s) [4].

3.1. **Some representative examples.** In the following examples, and for the sake of simplicity, we will use the notation

$$\sigma(s) = aq^{2s} + bq^{s} + c, \ \tau(s) = dq^{s} + e, \ \phi_{\nu}(s) = \sigma(s) + \tau_{\nu-1}(s)\nabla x_{\nu}(s) = fq^{2s} + gq^{s} + h. \ (22)$$

Example 3.3. The following relation holds

$$A_1(z)\Phi_{\nu,\nu-1}(z) + A_2(z)\Phi_{\nu,\nu} + A_3(z)\Phi_{\nu+1,\nu}(z) = 0$$

where the coefficients A_1 , A_2 and A_3 , are polynomials in $x \equiv x(z) = q^z$, given by

$$A_1(z) = -eq^{\frac{\nu}{2}} + \frac{b + e\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)}{a + d\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)} \left(dq^{\frac{\nu}{2}} + a[\nu]_q\right) + \left(dq^{\nu} + a[2\nu]_q\right)q^{\frac{\nu}{2} + z},$$

$$A_2(z) = \frac{c \left(dq^{\nu} + a[2\nu]_q\right)}{a + d\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)} + \frac{b + e\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)}{q^{\frac{1}{2}} - q^{-\frac{1}{2}}} \left(q^{\nu} + \frac{a}{q^{\nu}\left(a + d\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)\right)}\right) q^z + \left(dq^{\nu} + a[2\nu]_q\right)q^{2z},$$

$$A_3(z) = -\frac{dq^{\frac{\nu}{2}} + a[\nu]_q}{a + d\left(q^{\frac{1}{2}} - q^{-\frac{1}{2}}\right)},$$

where a, b, c, d, and e, are the coefficients of σ and τ (22).

Proof. Using the notations of Lemma 3.2 we have $\nu_1 = \nu$, $\nu_2 = \nu$, $\nu_3 = \nu + 1$, $\mu_1 = \nu - 1$, $\mu_2 = \nu$ and $\mu_3 = \nu$, thus $\nu_0 = \nu$ and $\mu_0 = \nu$. By (19)

$$\Pi(s) = A_1 \left(q^{s + \frac{\nu}{2}} - q^{z - \frac{\nu}{2}} \right) + A_2 + A_3 q^{-\frac{\nu}{2}} \left[\left(a + d \left(q^{\frac{1}{2}} - q^{-\frac{1}{2}} \right) \right) q^{2\nu + 2s} + \left(b + e \left(q^{\frac{1}{2}} - q^{-\frac{1}{2}} \right) \right) q^{\nu + s} + c \right]. \tag{23}$$

On the other hand, from (21) and because Q(s) = k is a constant –notice that $deg(\Pi) = 2$ – we have

$$\nabla x_{\nu_0+1}(s)\Pi(s) = k \left\{ \left[\left(a + d \left(q^{\frac{1}{2}} - q^{-\frac{1}{2}} \right) \right) q^{2\nu+2s} + \left(b + e \left(q^{\frac{1}{2}} - q^{-\frac{1}{2}} \right) \right) q^{\nu+s} + c \right] (q^s - q^z) - \left(aq^{2s} + bq^s + c \right) (q^{\nu+s} - q^z) \right\}$$
(24)

where k is an arbitrary constant. Introducing (23) in (24), using the identity

$$\nabla x_{\nu_0+1}(s) = q^{\frac{\nu}{2}} \left(q^{\frac{1}{2}} - q^{-\frac{1}{2}} \right) q^s$$

and comparing the coefficients of the powers of $x(s) = q^s$ we get a linear system of three equations with four variables A_1 , A_2 , A_3 and k. Choosing k = 1 and solving the corresponding system we get, after some simplifications, the coefficients A_1 , A_2 and A_3 .

In the next examples, since the technique is similar to the previous one we will omit the details.

Example 3.4. The following relation holds

$$A_1(z)\Phi_{\nu,\nu}(z) + A_2(z)\Phi_{\nu,\nu+1}(z) + A_3(z)\Phi_{\nu+1,\nu+1}(z) = 0$$

where the coefficients A_1 , A_2 and A_3 , are polynomials in $x \equiv x(z) = q^z$, given by

$$A_1(z) = f \left(a - f q^{2\nu} \right) q^z + a g q - f b q^{\nu+1} ,$$

$$A_2(z) = q^{-\frac{\nu}{2} - 1} \left(a - f q^{2\nu} \right) \left(f q^{2z} + g q^{z+1} + h q^2 \right) ,$$

$$A_3(z) = \sqrt{q} \left(a q - f q^{\nu} \right) ,$$

where a, b, c, f, g and h, are the coefficients of σ and ϕ_{ν} (22).

Example 3.5. The following relation holds

$$A_1(z)\Phi_{\nu-1,\nu-1}(z) + A_2(z)\Phi_{\nu,\nu-1}(z) + A_3(z)\Phi_{\nu,\nu}(z) = 0$$

where the coefficients A_1 , A_2 and A_3 , are polynomials in $x \equiv x(z) = q^z$, given by

$$\begin{split} A_1(z) = &q^{-\frac{1}{2}-\nu} \Big\{ fq^{2z} \Big[-a^2hq^4 + agbq^{\nu+4} - q^{2\nu+2} \left(ag^2q - 2fah + fb^2 \right) \\ &- fgbq^{3\nu+1} \left(q^2 - q - 1 \right) + fq^{4\nu} \left(g^2 \left(q - 1 \right) q - fh \right) \Big] + \\ gq^{z+1} \Big[-a^2hq^5 + aq^{\nu+2} \left(gbq^3 + fhq^2 - fh \right) - q^{2\nu+2} \left(\left(fah + fgb + ag^2 \right) q^2 - f\left(2ah - b^2 + gb \right) q - fah \right) + fq^{3\nu} \left(q^2 \left(g^2q - fh + gb - g^2 \right) + fh \right) + f^2hq^{4\nu} \left(q^2 - q - 1 \right) \Big] \\ &- a^2h^2q^6 + aghq^{\nu+5} \left(bq + gq - g \right) + fghq^{3\nu+4} \left(gq + b - g \right) - f^2h^2q^{4\nu+2} \\ &- hq^{2\nu+3} \left(ag^2q^3 + fgbq^2 + fg^2q^2 - 2fahq + fb^2q - 2fg^2q - fgb + fg^2 \right) \Big\} \,, \\ A_2(z) = &\left(q^{-\frac{\nu}{2}} - q^{\frac{\nu}{2}} \right) \left(fq^{2z} + gq^{z+1} + hq^2 \right) \left(fq^z \left(fq^{2\nu} - aq^2 \right) + fq^{\nu+1} \left(gq + b - g \right) - agq^3 \right), \\ A_3(z) = &f\left(fq^{\nu} - aq \right) \left[\left(fq^{2z} + hq^2 \right) \left(fq^{2\nu} - aq^2 \right) + gq^{z+1} \left(fq^{\nu} \left(q^{\nu} + q - 1 \right) - aq^3 \right) \right], \end{split}$$

where a, b, c, f, g and h, are the coefficients of σ and ϕ_{ν} (22).

Example 3.6. The following relation holds

$$A_1(z)\Phi_{\nu-1,\nu-1}(z) + A_2(z)\Phi_{\nu,\nu}(z) + A_3(z)\Phi_{\nu,\nu+1}(z) = 0$$

where the coefficients A_1 , A_2 and A_3 , are polynomials in $x \equiv x(z) = q^z$, given by

$$\begin{split} A_1(z) &= a^2hq^4 - agbq^{\nu+3} + q^{2\nu+2} \left(fb^2 - 2fah + ag^2\right) - fgbq^{3\nu+1} + f^2hq^{4\nu}, \\ A_2(z) &= q^{-\frac{1}{2}} \left(fq^{\nu} - aq^2\right) \left(fq^{z+2\nu} - aq^{z+2} + gq^{2\nu+1} - bq^{\nu+2}\right), \end{split}$$

$$A_3(z) = -q^{\frac{v-3}{2}} \left(f q^{2v} - aq^2 \right) \left(q^{v+1} - 1 \right) \left(gq^{z+1} + fq^{2z} + hq^2 \right),$$

where a, b, c, f, g and h, are the coefficients of σ and ϕ_{ν} (22).

Example 3.7. The relation

$$A_1(z)\Phi_{\nu,\nu-1}(z) + A_2(z)\Phi_{\nu,\nu}(z) + A_3(z)\Phi_{\nu+1,\nu+1}(z) = 0,$$

is verified when the polynomial coefficients A_1 , A_2 and A_3 , in the variable $x \equiv x(z) = q^z$, are given by

$$\begin{split} A_1(z) &= q^{\frac{\nu+1}{2}} \Big(fq^{z+\nu} \left(fq^{2\nu} - gq^{\nu} + b - a \right) - f \left(h - b \right) q^{2\nu+1} - aq \left(gq^{\nu} - h \right) \Big) \,, \\ A_2(z) &= q^{-z+\nu+\frac{1}{2}} \Big(q^z \left(fq^{2\nu} - a \right) + q^{\nu} \left(gq^{\nu} - b \right) \Big) \left(fq^{2z} + gq^{z+1} + hq^2 \right) \,, \\ A_3(z) &= q^{2z} \left(fq^{\nu} - aq \right) + q^{z+\nu} \Big(q \left(gq^{\nu} - aq - b \right) - fq^{\nu} \left(q^{\nu+1} - q - 1 \right) \Big) + q^{\nu+1} \Big(\left(h - b \right) q^{\nu+1} + gq^{\nu} - h \Big), \end{split}$$

where a, b, c, f, g and h, are the coefficients of σ and ϕ_{ν} (22).

4. Recurrences involving the solutions y_{ν}

In [16] the following relevant relation was established

$$\Delta^{(k)} y_{\nu}(s) = \frac{C_{\nu}^{(k)}}{\rho_k(s)} \Phi_{\nu, \nu - k}(s), \tag{25}$$

where

$$C_{\nu}^{(k)} = C_{\nu} \prod_{m=0}^{k-1} \left[\alpha_q(\nu + m - 1) \tilde{\tau}' + [\nu + m - 1]_q \frac{\tilde{\sigma}''}{2} \right].$$

This relation is valid for solutions of the form (3) and (4) of the difference equation (1). In the following, $y_n^{(k)}(s)$ denotes the k-th differences $\Delta^{(k)}y_n(s)$.

Theorem 4.1. In the same conditions as in Lemma 3.2, any three functions $y_{\nu_i}^{(k_i)}(s)$, i = 1, 2, 3, are connected by a linear relation

$$\sum_{i=1}^{3} B_i(s) y_{\nu_i}^{(k_i)}(s) = 0, \tag{26}$$

where the $B_i(s)$, i = 1, 2, 3, are polynomials.

Proof. From Lemma 3.2 we know that there exists three polynomials $A_i(s)$, i = 1, 2, 3 such that

$$\sum_{i=1}^{3} A_i(s) \Phi_{\nu_i, \nu_i - k_i}(s) = 0,$$

then, using the relation (25), we find

$$\sum_{i=1}^{3} A_i(s) (C_{\nu}^{(k)})^{-1} \rho_{k_i}(s) y_{\nu_i}^{(k_i)}(s) = 0.$$

Now, dividing the last expression by $\rho_{k_0}(s)$, where $k_0 = \min\{k_1, k_2, k_3\}$, and using (18) we obtain

$$\sum_{i=1}^{3} B_i(s) y_{\nu_i}^{(k_i)}(s) = 0, \quad B_i(s) = A_i(s) (C_{\nu}^{(k)})^{-1} \phi(s+k_0) \cdots \phi(s+k_i-1),$$

which completes the proof.

Corollary 4.2. In the same conditions as in Lemma 3.2, the following three-term recurrence relation holds

$$A_1(s)y_{\nu}(s) + A_2(s)y_{\nu+1}(s) + A_3(s)y_{\nu-1}(s) = 0,$$

with polynomial coefficients $A_i(s)$, i = 1, 2, 3.

Proof. It is sufficient to put $k_1=k_2=k_3=0,\ \nu_1=\nu,\ \nu_2=\nu+1$ and $\nu_3=\nu-1$ in (26).

Corollary 4.3. In the same conditions as in Lemma 3.2, the following Δ -ladder-type relation holds

$$B_1(s)y_{\nu}(s) + B_2(s)\frac{\Delta y_{\nu}(s)}{\Delta x(s)} + B_3(s)y_{\nu+m}(s) = 0, \qquad m \in \mathbb{Z},$$
 (27)

with polynomial coefficients $B_i(s)$, i = 1, 2, 3.

Proof. It is sufficient to put $k_1=k_3=0,\ k_2=1,\ \nu_1=\nu_2=\nu$ and $\nu_3=\nu+m$ in (26).

Notice that for the case $m = \pm 1$ (27) becomes

$$B_1(s)y_{\nu}(s) + B_2(s)\frac{\Delta y_{\nu}(s)}{\Delta x(s)} + B_3(s)y_{\nu+1}(s) = 0,$$
(28)

$$\widetilde{B}_1(s)y_{\nu}(s) + \widetilde{B}_2(s)\frac{\Delta y_{\nu}(s)}{\Delta x(s)} + \widetilde{B}_3(s)y_{\nu-1}(s) = 0,$$
 (29)

with polynomial coefficients $B_i(s)$ and $B_i(s)$, i = 1, 2, 3. The above relations are usually called raising and lowering operators, respectively, for the functions y_{ν} .

Let us now obtain a raising and lowering operators for the functions y_{ν} but associated to the $\nabla/\nabla x(s)$ operators.

We start applying the operator $\nabla/\nabla x(s)$ to (13)

$$\begin{split} \frac{\nabla}{\nabla x(s)} y_{\nu}(s) = & \frac{\nabla}{\nabla x(s)} \left[\frac{C_{\nu}}{\rho(s)} \Phi_{\nu,\nu}(s) \right] \\ = & \frac{1}{\nabla x(s)} \left[C_{\nu} \Phi_{\nu\nu}(s) \left(\frac{1}{\rho(s)} - \frac{1}{\rho(s-1)} \right) + \frac{C_{\nu}}{\rho(s-1)} \nabla \Phi_{\nu\nu}(s) \right], \end{split}$$

or, equivalently,

$$\frac{\nabla \Phi_{\nu\nu}}{\nabla x(s)} = \frac{\rho(s-1)}{C_{\nu}} \frac{\nabla y_{\nu}(s)}{\nabla x(s)} - \frac{\Phi_{\nu\nu}(s)}{\nabla x(s)} \left[\frac{\rho(s-1)}{\rho(s)} - 1 \right].$$

By Lemma (3.2) with $\nu_1 = \mu_1 = \nu_2 = \nu$, $\mu_2 = \nu + 1$ and $\nu_3 = \mu_3 = \nu + m$, there exist polynomial coefficients on x(s), $A_i(s)$, i = 1, 2, 3, such that

$$A_1(s)\Phi_{\nu,\nu}(s) + A_2(s)\Phi_{\nu,\nu+1}(s) + A_3(s)\Phi_{\nu+m,\nu+m}(s) = 0.$$

From (14)

$$\Phi_{\nu,\nu+1}(s) = \frac{1}{[\nu+1]_a} \frac{\nabla \Phi_{\nu,\nu}}{\nabla x(z)} = \frac{1}{[\nu+1]_a} \frac{\nabla \Phi_{\nu,\nu}}{\nabla x(z)}.$$

Therefore

$$A_{1}(s)\Phi_{\nu,\nu} + \frac{A_{2}(s)}{[\nu+1]_{q}} \left[\frac{\rho(s-1)}{C_{\nu}} \frac{\nabla y_{\nu}}{\nabla x(s)} - \frac{\Phi_{\nu\nu}(s)}{\nabla x(s)} \left(\frac{\rho(s-1)}{\rho(s)} - 1 \right) \right] + A_{3}\Phi_{\nu+m} + \mu = 0.$$

Using now the Pearson equation (6) and dividing by $\rho(s)$ we get

$$\begin{split} A_1(s)y_{\nu}(s) + \frac{A_2(q)}{[\nu+1]_q} \left[\frac{\sigma(s)}{\phi(s-1)} \frac{\nabla y_{\nu}}{\nabla x(s)} - \frac{y_{\nu}(s)}{\nabla x(s)} \left(\frac{\sigma(s)}{\phi(s-1)} - 1 \right) \right] \\ + A_3 \frac{C_{\nu}}{C_{\nu+m}} y_{\nu+m}(s) = 0 \; . \end{split}$$

Multiplying both sides by $[\nu + 1]_q \phi(s-1)$,

$$A_1(s)[\nu+1]_q \phi(s-1)y_{\nu}(s) + A_2(s)\sigma(s) \frac{\nabla y_{\nu}}{\nabla x(s)} - A_2(s) \frac{\sigma(s) - \phi(s-1)}{\nabla x(s)} y_{\nu}(s) + [\mu+1]_q C_{\nu} C_{\nu+m}^{-1} A_3 \phi(s-1) y_{\nu+m}(s) = 0.$$

Thus we have proven the following

Theorem 4.4. In the same conditions as in Lemma 3.2, the following ∇ -ladder-type relation holds

$$C_1(s)y_{\nu}(s) + C_2(s)\frac{\nabla y_{\nu}(s)}{\nabla x(s)} + C_3(s)y_{\nu+m}(s) = 0, \qquad m \in \mathbb{Z},$$
 (30)

with polynomial coefficients $C_i(s)$, i = 1, 2, 3.

Notice that for the case $m = \pm 1$ (30) becomes

$$C_1(s)y_{\nu}(s) + C_2(s)\frac{\nabla y_{\nu}(s)}{\nabla x(s)} + C_3(s)y_{\nu+1}(s) = 0,$$
(31)

$$\widetilde{C}_1(s)y_{\nu}(s) + \widetilde{C}_2(s)\frac{\nabla y_{\nu}(s)}{\nabla x(s)}y_{\nu}(s) + \widetilde{C}_3(s)y_{\nu-1}(s) = 0,$$
 (32)

with polynomial coefficients $C_i(s)$ and $\widetilde{C}_i(s)$, i = 1, 2, 3. The above relation are usually called raising and lowering operators, respectively, for the functions y_n . Eq. (31) was firstly obtained in [16, Eq. (3.4)].

To conclude this section let us point that from formula (25) and the examples 3.3, 3.5, and 3.7 follow the relations

$$B_{1}(s)y_{\nu}^{(1)}(s) + B_{2}(s)y_{\nu}(s) + B_{3}(s)y_{\nu+1}^{(1)}(s) = 0,$$

$$B_{1}(s)y_{\nu}^{(1)}(s) + B_{2}(s)y_{\nu-1}(s) + B_{3}(s)y_{\nu}(s) = 0,$$

$$B_{1}(s)y_{\nu}^{(1)}(s) + B_{2}(s)y_{\nu}(s) + B_{3}(s)y_{\nu+1}(s) = 0,$$
(33)

respectively, being the last two expressions the lowering and raising operators for the functions y_{ν} . Moreover, combining the explicit values of A_1 , A_2 and A_3 with formula (25), one can obtain the explicit expressions for the coefficients B_1 , B_2 and B_3 in (33).

5. Applications to q-classical polynomials

In this section we will apply the previous results to the q-classical orthogonal polynomials [2, 10, 11] in order to show how the method works. We first notice that these polynomials are instances of the functions y_{ν} on the lattice $x(s) = q^s$ defined in (4). In fact we have [13, 16]

$$P_n(x(s)) = \frac{[n]_q! B_n}{\rho(s) \ 2\pi i} \int_C \frac{\rho_n(z) \nabla x_{n+1}(z)}{[x_n(z) - x_n(s)]^{(n+1)}} dz, \tag{34}$$

where B_n is a normalizing constant, C is a closed contour surrounding the points $x = s, s - 1, \ldots, s - n$ and it is assumed that $\rho_n(s) = \rho(s+n) \prod_{m=1}^n \sigma(s+m)$ and $\rho_n(s+1)$ are analytic inside C (ρ is the solution of the Pearson equation (6)), i.e., the condition (5) holds.

A detailed study of the q-classical polynomials, including several characterization theorems, was done in [2, 9, 11]. In particular, a comparative analysis of the q-Hahn tableau with the q-Askey tableau [9] and Nikiforov-Uvarov tableau [15] was done in [5]. In the following we use the standard notation for the q-calculus [8]. In particular by $(a;q)_k = \prod_{m=0}^{k-1} (1-aq^m)$, we denote the q-analogue of the Pochhammer symbol.

Since the q-classical polynomials are defined by (34) where the contour C is closed and ν is a non-negative integer, then the condition (17) is automatically fulfilled, so Lemma 3.2 holds for all of them. Moreover, the Theorem 4.1 holds and there exist the non vanishing polynomials B_1 , B_2 and B_3 of (26).

In the following we will assume that the three term recurrence relation is known, i.e.,

$$x(s)P_n(x(s)) = \alpha_n P_{n+1}(x(s)) + \beta_n P_n(x(s)) + \gamma_n P_{n-1}(x(s)) = 0, \quad n \ge 0$$

$$P_{-1}(x(s)) = 0, \quad P_0(x(s)) = 1, \quad x(s) = q^s.$$
(35)

where the coefficients α_n , β_n and γ_n can be computed using the coefficients σ , τ and $\lambda \equiv \lambda_n$ of (1), being λ_n given by (8) and (9) with $\nu = n$. For more details see, e.g., [1, 11].

Since the TTRR and the differentiation formulas for the q-polynomials are very well known (see e.g. [9, 11, 16]) we will obtain here two recurrent-difference relations involving the q-differences of the polynomials and the polynomials themselves.

5.1. The first difference-recurrece relation. If we choose $\nu_1 = n - 1$, $\nu_2 = n$, $\nu_3 = n + 1$, $k_1 = 1$, $k_2 = 1$ and $k_3 = 0$, in Theorem 4.1 one gets

$$A_1(s)\Delta^{(1)}P_{n-1}(x(s)) + A_2(s)\Delta^{(1)}P_n(x(s)) + A_3(s)P_{n+1}(x(s)) = 0.$$

Using [1, Eq. (6.14), page 193]

$$[\sigma(s) + \tau(s)\Delta x(s - 1/2)]\Delta^{(1)}P_n(x(s)) = \widehat{\alpha}_n P_{n+1}(x(s)) + \widehat{\beta}_n P_n(x(s)) + \widehat{\gamma}_n P_{n-1}(x(s)),$$

where

$$\widehat{\alpha}_n = \frac{\lambda_n}{[n]_q} \left[q^{-\frac{n}{2}} \alpha_n - \frac{B_n}{\tau_n' B_{n+1}} \right], \quad \widehat{\beta}_n = \frac{\lambda_n}{[n]_q} \left[q^{-\frac{n}{2}} \beta_n + \frac{\tau_n(0)}{\tau_n'} - c_3(q^{-\frac{n}{2}} - 1) \right],$$

$$\widehat{\gamma}_n = \frac{\lambda_n q^{-\frac{n}{2}} \gamma_n}{[n]_q},$$

to compute $\Delta^{(1)}P_n(x(s)) = \frac{\Delta P_n(x(s))}{\Delta x(s)}$ we get

$$\begin{split} & \left[A_2(s) \frac{\lambda_n}{[n]_q} \left(q^{-\frac{n}{2}} \alpha_n - \frac{B_n}{\tau_n' B_{n+1}} \right) + \left(\sigma(s) + \tau(s) \Delta x \left(s - \frac{1}{2} \right) \right) A_3(s) \right] P_{n+1} + \\ & \left[A_1(s) \frac{\lambda_{n-1}}{[n-1]_q} \left(q^{-\frac{n-1}{2}} \alpha_{n-1} - \frac{B_{n-1}}{\tau_{n-1}' B_n} \right) + A_2(s) \frac{\lambda_n}{[n]_q} \left(q^{-\frac{n}{2}} \beta_n + \frac{\tau_n(0)}{\tau_n'} \right) \right] P_n + \\ & \left[A_1(s) \frac{\lambda_{n-1}}{[n-1]_q} \left(q^{-\frac{n-1}{2}} \beta_{n-1} + \frac{\tau_{n-1}(0)}{\tau_{n-1}'} \right) + A_2(s) \frac{\lambda_n q^{-\frac{n}{2}} \gamma_n}{[n]_q} \right] P_{n-1} + \\ & A_1(s) \frac{\lambda_{n-1} q^{-\frac{n-1}{2}} \gamma_{n-1}}{[n-1]_q} P_{n-2} = 0 \,, \end{split}$$

By (35) we may write

$$P_{n-2}(x(s)) = \frac{x(s) - \beta_{n-1}}{\gamma_{n-1}} P_{n-1}(x(s)) - \frac{\alpha_{n-1}}{\gamma_{n-1}} P_n(x(s))$$

so the above equality becomes

$$\left[\frac{\lambda_{n}}{[n]_{q}}\left(q^{-\frac{n}{2}}\alpha_{n} - \frac{B_{n}}{\tau_{n}'B_{n+1}}\right)A_{2}(s) + \left(\sigma(s) + \tau(s)\Delta x\left(s - \frac{1}{2}\right)\right)A_{3}(s)\right]P_{n+1}(x(s)) + \\
\left[-\frac{\lambda_{n-1}}{[n-1]_{q}}\frac{B_{n-1}}{\tau_{n-1}'B_{n}}A_{1}(s) + \frac{\lambda_{n}}{[n]_{q}}\left(q^{-\frac{n}{2}}\beta_{n} + \frac{\tau_{n}(0)}{\tau_{n}'}\right)A_{2}(s)\right]P_{n}(x(s)) + \\
\left[\frac{\lambda_{n-1}}{[n-1]_{q}}\left(\frac{\tau_{n-1}(0)}{\tau_{n-1}'} + q^{-\frac{n-1}{2}}x\right)A_{1}(s) + \frac{\lambda_{n}}{[n]_{q}}q^{-\frac{n}{2}}\gamma_{n}A_{2}(s)\right]P_{n-1}(x(s)) = 0.$$
(36)

Comparing the above equation with the TTRR (35) one can obtain the explicit values of A_1 , A_2 , and A_3 .

5.1.1. Some examples. Since we are working in the q-linear lattice $x(s) = q^s$, for the sake of simplicity, we will use the letter x to denote the variable of the polynomials [9, 11]. We will consider monic polynomials, i.e., those with the leading coefficient equal to 1. In the following we need the value of $\tau_n(x)$ for each family, which can be computed using (7).

Al-Salam-Carlitz I q-polynomials. For the Al-Salam-Carlitz I monic polynomials $U_n^{(a)}(x;q)$ we have (see [1, see table 6.5, p.208] or [11])

$$\sigma(x) = (1-x)(a-x), \quad \tau_n(x) = \frac{q^{\frac{1-n}{2}}}{1-q} \left(x - (1+a) \right),$$

$$\tau(x) = \tau_0(x), \quad \lambda_n = -\frac{q^{\frac{3}{2}-n}(1-q^n)}{(1-q)^2},$$

and

$$\alpha_n = 1$$
, $\beta_n = (1+a)q^n$, $\gamma_n = -aq^{n-1}(1-q^n)$.

The constant B_n is given by [1, Eq. (5.57), p. 147], $B_n = q^{\frac{1}{4}n(3n-5)}(1-q)^n$. Introducing these values into the equation (36) it becomes

$$\left[q\left(q^{-\frac{n}{2}}-1\right)A_{2}(x)+a(1-q)q^{n}A_{3}(x)\right]U_{n+1}^{(a)}(x;q)+$$

$$\left[q^{-\frac{n}{2}-\frac{5}{2}}A_{1}(x)+q^{1+\frac{n}{2}}(1+a)\left(1-q^{\frac{n}{2}}\right)A_{2}(x)\right]U_{n}^{(a)}(x;q)+$$

$$\left[\left(q^{\frac{n+3}{2}}(1+a)-q^{2-n}x\right)A_{1}(x)+aq^{n}\left(1-q^{n}\right)A_{2}(x)\right]U_{n-1}^{(a)}(x;q)=0.$$

Comparing with the TTRR (35) for the Al-Salam I polynomials we obtain a linear system for getting the unknown coefficients A_1 , A_2 and A_3

$$q\left(q^{-\frac{n}{2}}-1\right)A_2(x) + a(1-q)q^n A_3(x) = 1,$$

$$q^{-\frac{n}{2}-\frac{5}{2}}A_1(x) + q^{1+\frac{n}{2}}(1+a)\left(1-q^{\frac{n}{2}}\right)A_2(x) = (1+a)q^n - x,$$

$$\left(q^{\frac{n+3}{2}}(1+a) - q^{2-n}x\right)A_1(x) + aq^n\left(1-q^n\right)A_2(x) = aq^{n-1}\left(q^n - 1\right).$$

The solution of the above system is

$$A_{1}(x) = \frac{aq^{n}\left(1+q^{\frac{n}{2}}\right)\left((1+a)-q^{-\frac{n}{2}}x\right)}{aq^{-\frac{5}{2}}\left(1+q^{\frac{n}{2}}\right)-q(1+a)\left(q^{\frac{n+3}{2}}(1+a)-q^{2-n}x\right)},$$

$$A_{2}(x) = \frac{-aq^{-\frac{7}{2}}(1-q^{n})-\left((1+a)q^{n}-x\right)\left(q^{\frac{3}{2}}(1+a)-q^{2-\frac{3n}{2}}x\right)}{\left(1-q^{\frac{n}{2}}\right)\left[aq^{-\frac{5}{2}}\left(1+q^{\frac{n}{2}}\right)-q(1+a)\left(q^{\frac{n+3}{2}}(1+a)-q^{2-n}x\right)\right]},$$

$$A_{3}(x) = \frac{a+q^{\frac{11}{2}-2n}x^{2}+q^{-\frac{n}{2}}\left(a-(1+a)q^{5}x\right)}{a(1-q)\left[aq^{n}+q^{\frac{3n}{2}}\left(a-(1+a)^{2}q^{5}\right)+(1+a)q^{\frac{11}{2}}x\right]}.$$

$$(37)$$

Then, the Al-Salam I q-polynomials satisfy the the following relation

$$A_1(x)\Delta^{(1)}U_{n-1}^{(a)}(x;q) + A_2(x)\Delta^{(1)}U_n^{(a)}(x;q) + A_3(x)U_{n+1}^{(a)}(x;q) = 0,$$
(38)

where the coefficients A_1 , A_2 and A_3 are given by (37).

Notice that the coefficients A_1 , A_2 and A_3 are rational functions on x. Therefore, multiplying (38) by and appropriate factor it becomes a linear relation with polynomials coefficients.

Alternative q-Charlier polynomials. In this case (see [1, table 6.6, p.209])

$$\sigma(x) = q^{-1}x(1-x), \quad \tau_n(x) = -\frac{q^{-\frac{n+1}{2}}}{1-q} \left(\left(1 + aq^{1+2n} \right) x - 1 \right),$$

$$\tau(x) = \tau_0(x), \quad \lambda_n = \frac{q^{\frac{1}{2}-n}(1-q^n)(1+aq^n)}{(1-q)^2},$$

and, for the monic case, $\alpha_n = 1$

$$\beta_n = \frac{q^n \left(1 + aq^{n-1} + aq^n - aq^{2n}\right)}{\left(1 + aq^{2n-1}\right)\left(1 + aq^{2n+1}\right)}, \quad \gamma_n = \frac{aq^{3n-2}\left(1 - q^n\right)\left(1 + aq^{n-1}\right)}{\left(1 + aq^{2n-2}\right)\left(1 + aq^{2n-1}\right)^2\left(1 + aq^{2n}\right)}.$$

The corresponding normalizing constant B_n is given by

$$B_n = \frac{(-1)^n q^{\frac{1}{4}n(3n-1)} (1-q)^n}{(-aq^n; q)_n}.$$

Following the same procedure as before we obtain the following relation for the alternative Charlier q-polynomials:

$$A_1(x)\Delta^{(1)}K_{n-1}(x;a;q) + A_2(x)\Delta^{(1)}K_n(x;a;q) + A_3(x)K_{n+1}(x;a;q) = 0,$$

with the coefficients

$$\begin{split} A_1(x) = & \frac{a \left(1 + aq^{\frac{n}{2}}\right) \left(\left(1 + aq^{2n+1}\right) x - q^{-\frac{n}{2}}\right) x}{q^2 \left(1 + aq^{2n-2}\right) \left(1 + aq^{2n-1}\right) \left(1 + aq^{2n}\right) \left(1 + aq^{2n+1}\right)} \,, \\ A_2(x) = & \frac{-q^{\frac{3n+1}{2}} \left(1 + aq^n\right) x + \left(1 + aq^{2n}\right) \left(q^{1 + \frac{n}{2}} \left(1 + aq^{2n+1}\right) + aq^{2n + \frac{1}{2}} \left(1 + q\right) + q^{\frac{3n}{2}} \left(1 - aq^{2n}\right)\right) x^2}{q^{3n} \left(1 + aq^n\right) \left(1 + aq^{2n}\right) \left(1 + aq^{2n+1}\right)} \\ & \frac{q^{\frac{3}{2}} \left(1 + aq^{2n-1}\right) \left(1 + aq^{2n+1}\right) x^3}{q^{3n} \left(1 + aq^n\right) \left(1 + aq^{2n}\right) \left(1 + aq^{2n+1}\right)} \,, \\ A_3(x) = & \frac{q^{\frac{n+1}{2}} + aq^{2n} \left(q^{\frac{n}{2}} + 1 + q^{\frac{1}{2}}\right) - q^{\frac{3}{2}} \left(1 - aq^{\frac{3n}{2}}\right) \left(1 + aq^{2n-1}\right) x}{q^{\frac{9n}{2}} \left(1 + aq^n\right)} \,. \end{split}$$

Big q-Jacobi polynomials. In this case (see [1, see table 6.2, p.204] or [11])

$$\sigma(x) = q^{-1}(x - aq)(x - cq), \lambda_n = -q^{\frac{1}{2} - n} \frac{\left(1 - abq^{1+n}\right)\left(1 - q^n\right)}{(1 - q)^2},$$

$$\tau_n(x) = \frac{q^{\frac{1-n}{2}}}{1 - q} \left(\frac{1 - abq^{2+2n}}{q}x + a(b+c)q^{1+n} - (a+c)\right), \tau(x) = \tau_0(x),$$

and, for the monic case $\alpha_n = 1$,

$$\beta_n = \frac{c + a^2 b q^n \left((1 + b + c) q^{1+n} - q - 1 \right) + a \left(1 + b + c - q^n \left(b (1 + q) + c \left(1 + q + b + b q - b q^{1+n} \right) \right) \right)}{q^{-1-n} (1 - abq^{2n}) (1 - abq^{2n+2})},$$

$$\gamma_n = -\frac{a \left(1 - q^n \right) \left(1 - aq^n \right) \left(1 - bq^n \right) \left(1 - cq^n \right) \left(c - abq^n \right)}{q^{-1-n} \left(1 - abq^{2n-1} \right) \left(1 - abq^{2n} \right)^2 \left(1 - abq^{2n+1} \right)}.$$

The corresponding normalizing constant is

$$B_n = \frac{(1-q)^n q^{\frac{1}{4}n(3n-1)}}{(abq^{1+n};q)_n}.$$

The big q-Jacobi polynomials satisfy the following relation

$$A_1(x)\Delta^{(1)}p_{n-1}(x;a,b,c;q) + A_2(x)\Delta^{(1)}p_n(x;a,b,c;q) + A_3(x)p_{n+1}(x;a,b,c;q) = 0,$$

with the coefficients A_1 , A_2 and A_3 given by

$$A_{1}(x) = \frac{aq^{-\frac{1}{2}+n}(1-abq^{n+1})(1-x)(c-bx)\left(c-(b+c)x+bx^{2}\right)}{1-abq^{2n-1}} \times \left\{ (1-q)q^{\frac{n}{2}}\left(1-abq^{2n+2}\right) \left[\frac{c+a\left(1+b+c+b(c+a(1+b+c))q^{2n+1}-(c+b(1+a+c))q^{n}(1+q)\right)}{q^{-(n+1)}\left(1-abq^{2n}\right)\left(1-abq^{2n+2}\right)} - x \right] D(x) - (1-q)q^{n} \left(1-abq^{2n}\right) \left[\left(1-abq^{2n}\right)\left(-c+a\left(-1+(b+c)q^{n+1}\right)\right) + q^{\frac{n}{2}}\left(c+a\left(1+b+c+b\left(c+a(1+b+c\right)\right)q^{2n+1}-\left(c+b(1+a+c)\right)q^{n}(1+q)\right)\right) N(x) \right] \right\},$$

$$A_{2}(x) = a(1-q)q^{n} \left(1-abq^{2n}\right)^{2} \left(1-abq^{2n+2}\right) (1-x)(c-bx)\left(c-(b+c)x+bx^{2}\right) N(x),$$

$$A_{3}(x) = \left(1-abq^{n+1}\right) \left(1-abq^{2n+2}\right) (1-x)(c-bx)D(x) + q^{-1-\frac{n}{2}}\left(1-q^{\frac{n}{2}}\right) \left(1+abq^{1+\frac{3n}{2}}\right) \left(1-abq^{2n}\right)^{2} \left(1-abq^{2n+2}\right) \left(c-(b+c)x+bx^{2}\right) N(x),$$

where the polynomials N(x) and D(x) are given by

$$N(x) = \frac{aq^{2}(1-q^{n})(1-aq^{n})(1-bq^{n})(1-cq^{n})(c-abq^{n})}{(1-abq^{2n})^{2}(1-aq^{2n+1})} - \left[\frac{q\left(-c+a(-1+(b+c)q^{n})\right)}{1-abq^{2n}} + q^{\frac{1-n}{2}}x\right] \times \left[\frac{c+a^{2}bq^{n}\left(-1-q+(1+b+c)q^{n+1}\right)+a\left(1-(b+c)\left(-1+q^{n}+q^{n+1}\right)-bcq^{n}\left(1+q-q^{n+1}\right)\right)}{q^{-n-1}(1-abq^{2n})(1-abq^{2n+2})} - x\right]$$

and

$$\begin{split} D(x) &= \frac{aq(1-q^n)(1-aq^n)(1-bq^n)(1-cq^n)(c-abq^n)}{1-abq^{2n+1}} + \frac{1-q^{\frac{n}{2}}}{1-abq^{2n+2}} \times \\ &\left\{ -c + a^2bq^{\frac{3n}{2}} \left(-1 - q + (b+c)q^{n+1} - q^{1+\frac{n}{2}} \right) + a \left[-1 + (b+c)\left(q^{\frac{n}{2}} + q^n + q^{n+1}\right) - bc\left(q^{\frac{3n}{2}} + q^{1+\frac{3n}{2}} + q^{2n+1}\right) \right] \left[(c+a)q^{1+\frac{n}{2}} - a(b+c)q^{1+\frac{3n}{2}} - q^{\frac{1}{2}} \left(1 - abq^{2n} \right) x \right] \right\}, \end{split}$$

respectively.

5.2. The second difference-recurrece relation. If we choose $\nu_1 = n - 1$, $\nu_2 = n$, $\nu_3 = n + 1$, $k_1 = 0$, $k_2 = 0$ and $k_3 = 1$ in Theorem 4.1, and proceeding as in the previous case one gets

$$A_1(x)P_{n-1}(x;q) + A_2(x)P_n(x;q) + A_3(x)\Delta^{(1)}P_{n+1}(x;q) = 0,$$
(39)

where the coefficients A_1 , A_2 and A_3 , satisfy the linear relation

$$A_{3}(x) \left[\left(q^{-\frac{n+1}{2}} - \frac{B_{n+1}}{\alpha_{n+1}\tau'_{n+1}B_{n+2}} \right) \left(x - \beta_{n+1} \right) + \left(q^{-\frac{n+1}{2}}\beta_{n+1} + \frac{\tau_{n+1}(0)}{\tau'_{n+1}} \right) \right] P_{n+1} + \left[A_{3}(x) \frac{B_{n+1}}{\alpha_{n+1}\tau'_{n+1}B_{n+2}} \gamma_{n+1} + \left(\sigma(x) + \tau(x)\Delta x \left(s - \frac{1}{2} \right) \right) \frac{[n+1]_{q}}{\lambda_{n+1}} A_{2}(x) \right] P_{n} + \left(\sigma(x) + \tau(x)\Delta x \left(s - \frac{1}{2} \right) \right) \frac{[n+1]_{q}}{\lambda_{n+1}} A_{1}(x) P_{n-1} = 0.$$

Comparing the above relation with the three-term recurrence relation (35) one can obtain the explicit expressions for the coefficients A_1 , A_2 and A_3 in (39).

5.2.1. Some examples.

Al-Salam and Carlitz I polynomials. Using the main data for the Al-Salam and Carlitz I polynomials we obtain the relation

$$A_1(x)U_{n-1}^{(a)}(x;q) + A_2(x)U_n^{(a)}(x;q) + A_3(x)\Delta^{(1)}U_{n+1}^{(a)}(x;q) = 0$$

where

$$A_1(x) = aq^{n-1} (1 - q^n) x, \quad A_2(x) = \left[a \left(1 + q^{\frac{n+1}{2}} \right) q^n - \left((1+a)q^n - x \right) x \right],$$

$$A_3(x) = -a \frac{1 - q}{1 - q^{\frac{n+1}{2}}} q^{\frac{3n+1}{2}}.$$

Alternative q-Charlier polynomials. In this case, one gets

$$A_1(x)K_{n-1}(x;a;q) + A_2(x)K_n(x;a;q) + A_3(x)\Delta^{(1)}K_{n+1}(x;a;q) = 0,$$

$$A_1(x) = \frac{a\left(1-q^n\right)\left(1+aq^{n-1}\right)\left\{aq^n\left(1-q^{n+1}\right)+q^{-\frac{n+1}{2}}\left(1+aq^{2n+1}\right)\left[\left(1+aq^{n+1}\right)-q^{-\frac{n+1}{2}}\left(1+aq^{2n+2}\right)\right]x\right\}}{q^{2-3n}\left(1+aq^{2n-2}\right)\left(1+aq^{2n-1}\right)\left(1+aq^{2n}\right)},$$

$$A_2(x) = -x\left\{aq^n\left(1-q^{n+1}\right)+q^{-\frac{n+1}{2}}\left(1+aq^{2n+1}\right)\left[\left(1+aq^{n+1}\right)-q^{-\frac{n+1}{2}}\left(1+aq^{2n+2}\right)\right]x\right\}+$$

$$\frac{a^2q^{3n-1}\left(1-q^n\right)\left(1-q^{n+1}\right)+q^{\frac{n-1}{2}}\left(1+aq^{n-1}+aq^{n}-aq^{2n}\right)\left(1+aq^{2n+1}\right)\left[\left(1+aq^{n+1}\right)-q^{-\frac{n+1}{2}}\left(1+aq^{2n+2}\right)\right]x}{\left(1+aq^{2n-1}\right)\left(1+aq^{2n+1}\right)},$$

$$A_3(x) = a(1-q)q^{\frac{n+1}{2}}(1+aq^{2n+1})x^2$$

Concluding remarks. In this paper we present a constructive approach for finding recurrence relations for the hypergeometric-type functions on the linear-type lattices, i.e., the solutions of the hypergeometric difference equation (1) on the linear-type lattices. Important instances of "discret" functions are the celebrated Askey-Wilson polynomials and q-Racah polynomials. Such functions are defined on the non-uniform lattice of the form $x(s) = c_1(q)q^s + c_2(q)q^{-s} + c_2(q)$ with $c_1c_2 \neq 0$, i.e., a non linear-type lattice and therefore they require a more detailed study (some preliminar general results can be found in [16]).

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