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Tidal circulation, sediment and pollutant transport in Cádiz Bay (SW Spain): A modelling study



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ABSTRACT

A numerical model to simulate the dispersion of particle-reactive tracers in Cádiz Bay (SW Spain) has been developed. It includes a hydrodynamic submodel to provide water currents, a sediment transport submodel, which provides suspended matter concentrations and sedimentation rates and the pollutant dispersion model. Pollutant exchanges between the liquid and solid phases are described in a dynamic way, using kinetic transfer coefficients. Results of the hydrodynamic and sediment transport models have been compared with observations. The contamination of sediments of the bay by fallout ¹³⁷Cs has been simulated. The existence of an accumulation area in the inner bay, found in field measurements, has been reproduced by the model. Dispersion experiments for Zn and ²²⁶Ra have also been carried out. Flushing-times of the inner bay have been finally determined through numerical experiments.

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1. Introduction

Cádiz Bay is located in the southwest of Spain. It is a semi-enclosed water body, landlocked except along its northwestern limit where it faces the Atlantic Ocean. It may be divided (Fig. 1) into two regions: the shallower inner bay and the deeper outer bay, open to the Atlantic Ocean, which are connected through a narrow channel. Anyway the whole bay is shallow, with maximum depth around 20 m in its connection to the open sea. The inner bay is characterized by the presence of extensive intertidal areas, which are covered by water during part of each tidal cycle. Tides are predominantly semidiurnal, with a M_2 amplitude of approximately 1 m. Water circulation inside the bay is mainly controlled by tidal dynamics (Araujo et al., 2009).

Several cities surround the bay: Rota, El Puerto de Santa María, Puerto Real, San Fernando and Cádiz; thus some 400 000 people live in the area. Main industrial activities are related to ship, offshore, car and aerospace manufacturing (Carrasco et al., 2003). As a result of this population and industrial activity contamination in water and seabed sediments of the bay has been reported (Araujo et al., 2009; Carrasco et al., 2003; Gómez-Parra et al., 1984; Ligeró et al., 2002).

Some interesting modelling works concerning tidal dynamics in the bay have already been carried out (Alvarez et al., 1999, 2003; Kagan et al., 2001; Alvarez, 1999). However, modelling studies on the dynamics of pollutants in Cádiz Bay have not been published before. These studies are useful, through the knowledge of water

and sediment dynamics, to predict the effects of planned authorized releases of pollutants, as well as to assess the impact of potential accidental spills in the system. Studies of similar nature have been recently reported for Algeciras Bay (south Spain), also affected by industrial and urban waste releases (Periañez, 2012), and for southern Spain coastal waters (Periañez, 2011).

Thus, the objective of this paper consists of presenting a model for the Cádiz Bay suitable for the description of the transport of non-conservative pollutants, which are those that do not remain dissolved in the water column, but are fixed to suspended particles and to seabed sediments. Actually, many pollutants as heavy metals and some radioisotopes are predominantly attached to these solid phases. A hydrodynamic model, which provides the currents that induce advective transport and turbulent mixing in the water column, as well as a sediment transport model is required for describing the transport processes of these particle-reactive pollutants. The sediment model will calculate suspended matter concentrations and sedimentation rates, which are relevant for describing adsorption of pollutants by suspended matter and pollutant exchanges between suspended particles in the water column and seabed sediments. The model is described in the next section. Later, results are presented and discussed.

2. Model description

2.1. Tidal simulations

As mentioned above, water circulation in the bay is essentially controlled by tides, thus our focus has been simulating tidal dynamics.

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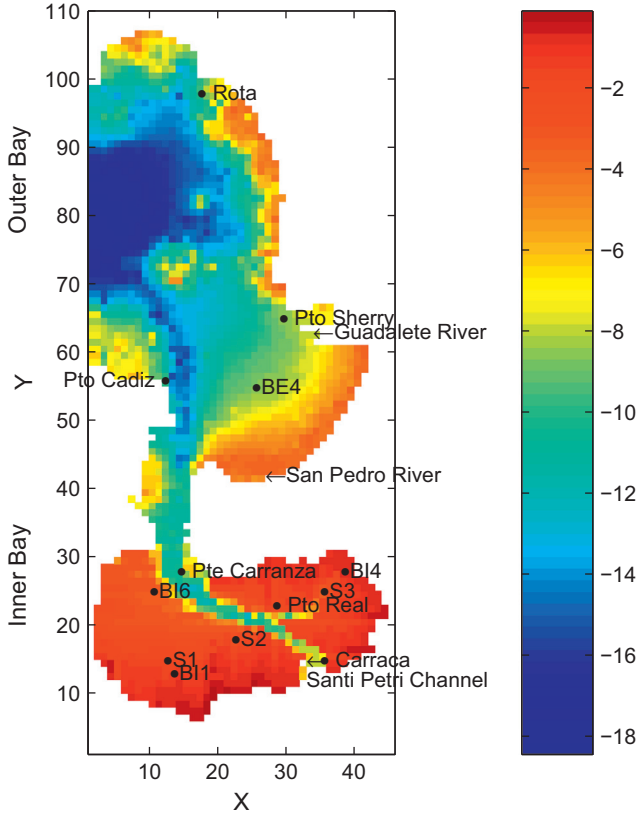


Fig. 1. Model domain. The color scale gives water depths in m. Locations mentioned in the text are displayed. Arrows indicate main rivers. Each unit in the axis gives the grid cell number. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this article.)

Some authors have stated that using a 2D model for simulating tides is a reasonable approach (Dyke, 2001; Yanagi, 1999). Also, it must be considered that the bay is rather shallow. For instance, mean depth in the inner bay is only about 3 m. Also, flows of channels and rivers discharging in the bay are extremely low and have been neglected in previous modelling studies of tidal dynamics (Alvarez, 1999). Significant horizontal or vertical density gradients are not expected. Consequently, depth-averaged models have already been successfully used to study the hydrodynamics of Cádiz Bay (Alvarez, 1999; Alvarez et al., 2003; Kagan et al., 2001). The influence of sediment load on tidal dynamics has also been studied with a depth-averaged model (Alvarez et al., 1999). Moreover, the theoretical analysis in Prandle et al. (1993) indicates that vertical homogeneity of substance concentrations is maintained for water depths lower than 50 m if tidal action ensures that the vertical eddy viscosity is at least of the order of $10^{-3} \text{ m}^2/\text{s}$, a value which may be considered representative of a relatively weak tide. Consequently, the use of a 2D depth averaged barotropic model is justified as a first approach. It is based on the following equations:

$$\frac{\partial \zeta}{\partial t} + \frac{\partial}{\partial x}(Hu) + \frac{\partial}{\partial y}(Hv) = 0; \quad (1)$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + g \frac{\partial \zeta}{\partial x} - fv + \frac{\tau_u}{\rho_w H} = A \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right); \quad (2)$$

$$\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + g \frac{\partial \zeta}{\partial y} + fu + \frac{\tau_v}{\rho_w H} = A \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right), \quad (3)$$

where u and v are the depth averaged water velocities along the x - and y -axes, h is the depth of water below the mean sea level, ζ is the displacement of the water surface above the mean sea level (positive upwards), $H = h + \zeta$ is the total water depth, f is the Coriolis

parameter ($f = 2\Omega \sin \lambda$, where Ω is the Earth rotational angular velocity and λ is latitude), g is acceleration due to gravity, ρ_w is water density and A is the horizontal eddy viscosity. Friction stresses τ_u and τ_v are written in terms of a quadratic law:

$$\begin{aligned} \tau_u &= k\rho_w u \sqrt{u^2 + v^2} \\ \tau_v &= k\rho_w v \sqrt{u^2 + v^2}, \end{aligned} \quad (4)$$

where k is the bed friction coefficient.

The solution of these equations provides the water currents at each point in the model domain and for each time step. Currents are treated through standard tidal analysis (Pugh, 1987, Chapter 4) and tidal constants are stored in files that will be used by the dispersion model to calculate the advective transport. The model includes the following tidal constituents, M_2 , S_2 , N_2 , M_1 , O_1 and M_4 . The hydrodynamic equations are solved for each constituent and tidal analysis is also carried out for each constituent separately. A residual transport cannot be produced by the pure harmonic currents given by the tidal analysis, thus tidal residuals have been calculated as well. Tidal residuals for each constituent are calculated from the equation:

$$\vec{q}_r = \frac{\langle H\vec{q}_t \rangle}{\langle H \rangle}, \quad (5)$$

which corresponds to the Eulerian residual transport velocity (Delhez, 1996). In this equation $\langle \rangle$ is the time averaging operator, \vec{q}_r is the tidal residual and \vec{q}_t is the instantaneous tidal current. Tidal constants and residual are stored in files which are later read by the sediment and pollutant transport models for a extremely fast calculation of current in each grid cell in the domain.

Due to the large intertidal areas existing in the inner bay, a flood-dry algorithm is required. The one reported in Kampf (2009) has been applied in this study. Wet grid cells are defined as those with a total water depth H larger than a threshold value H_{min} typically set as a few centimeters. Dry cells are defined as cells where $H \leq H_{min}$. Flooding and drying is implemented in the code via the calculation of the water velocity normal to the interface between wet and dry cells. The calculation is performed when the pressure gradient force is directed towards the dry cell. Otherwise velocity is set to zero at this point. In the case of a non-zero velocity, water level in the dry cell will increase and the cell turns into a wet one once the water depth is larger than H_{min} .

2.2. Sediment transport

Sediment dynamics is described in a very simple way, but still retaining its main aspects. This has proven to be enough to satisfactorily simulate reactive pollutant dispersion. The transport of sediments is described by a 2D advection–diffusion equation to which some terms are added. These are external sources of particles, terms describing particle deposition on the seabed and erosion from the bed to the water column. The formulation of these processes is based upon standard formulae. Thus, the erodability constant is used for the erosion term. Particle deposition is described using the settling velocity, which is obtained from Stokes' law. Critical erosion and deposition stresses are applied as usually (Perić, 2005a, 2005b; Liu et al., 2002; Lumborg and Windelin, 2003).

The equation for suspended sediment transport is

$$\begin{aligned} \frac{\partial(Hm)}{\partial t} + \frac{\partial(uHm)}{\partial x} + \frac{\partial(vHm)}{\partial y} &= \frac{\partial}{\partial x} \left(HK_h \frac{\partial m}{\partial x} \right) + \frac{\partial}{\partial y} \left(HK_h \frac{\partial m}{\partial y} \right) \\ &+ (ER - DEP) + S, \end{aligned} \quad (6)$$

where m is the suspended matter concentration, K_h is an effective horizontal diffusion coefficient, S is the external particle source and ER and DEP are the erosion and deposition terms, respectively.

The deposition term is written in the following form:

$$DEP = w_s m \left(1 - \frac{|\vec{\tau}|}{\tau_{cd}} \right), \quad (7)$$

where $\vec{\tau}$ is the bottom friction stress (whose components are given by Eq. (4)) and τ_{cd} is a critical deposition stress above which no deposition occurs since particles are kept in suspension by turbulence. The settling velocity of particles is obtained from Stokes' law as mentioned above:

$$w_s = \frac{\rho - \rho_w}{\rho_w} \frac{gD^2}{18\nu}, \quad (8)$$

where ρ and D are suspended particle density and diameter, respectively, and ν is the kinematic viscosity of water.

The erosion rate is written in terms of the erodability constant E :

$$ER = Ef_p \left(\frac{|\vec{\tau}|}{\tau_{ce}} - 1 \right), \quad (9)$$

where f_p gives the fraction of fine particles in the bed sediment and τ_{ce} is a critical erosion stress below which no erosion occurs. The model can also calculate sedimentation rates (SR) as the balance between the deposition and erosion terms.

2.3. Pollutant transport

Non-conservative pollutants are those which do not remain dissolved in the water column, but have a certain affinity to be fixed to particles. If the pollutant is introduced in the water column, it will be fixed to settling suspended particles and their deposition on the sea bed will contaminate the bottom sediment. Of course there are also advection–diffusion processes in water and direct adsorption of pollutants on the seabed. Three phases are then considered: dissolved, suspended matter particles and active bed sediments. These correspond to the fine sediments (particles with diameter $< 62.5 \mu\text{m}$) since metals and radionuclides are predominantly fixed to them (Eisma, 1993). The exchanges between the dissolved and solid phases may be described in terms of kinetic transfer coefficients. Thus, assuming that adsorption/release reactions are described by a single reversible reaction, a coefficient k_1 characterizes the transfer from the liquid to the solid phases and a coefficient k_2 characterizes the inverse process. Dimensions of these coefficients are $[T]^{-1}$.

The adsorption process is a surface phenomenon that depends on the surface of particles per water volume unit. This quantity has been denoted as the exchange surface (Periañez, 2008, 2009). Thus in general

$$k_1 = \chi(S_m + S_s) = k_1^{spm} + k_1^{sed}, \quad (10)$$

where S_m and S_s are the exchange surfaces for suspended matter and bottom sediments, respectively (dimensions $[L]^{-1}$). χ is a parameter with the dimensions of a velocity. It is denoted as the exchange velocity (Periañez, 2008, 2009).

Assuming spherical particles, the exchange surfaces are written as (see references cited above)

$$S_m = \frac{3m}{\rho R} \quad (11)$$

and

$$S_s = \frac{3Lf_p(1-p)\phi}{RH}, \quad (12)$$

where R is particle radius, p is sediment porosity, L is the sediment mixing depth (the distance to which the dissolved phase penetrates the sediment) and ϕ is a correction factor that takes into account that part of the sediment particle surface may be hidden

by other sediment particles. This formulation has been successfully used in all modelling works cited above. Real particles are not spheres, but with this approach it is possible to obtain an analytical expression for the exchange surface (Duursma and Carroll, 1996).

The equation that gives the temporal evolution of pollutant concentration in the dissolved phase, C_d , is

$$\begin{aligned} \frac{\partial(HC_d)}{\partial t} + \frac{\partial(uHC_d)}{\partial x} + \frac{\partial(vHC_d)}{\partial y} \\ = \frac{\partial}{\partial x} \left(HK_h \frac{\partial C_d}{\partial x} \right) + \frac{\partial}{\partial y} \left(HK_h \frac{\partial C_d}{\partial y} \right) \\ - H(k_1^{spm} + k_1^{sed})C_d + k_2 H m C_s + k_2 L \rho_s \phi f_p A_s, \end{aligned} \quad (13)$$

where A_s and C_s are concentrations in the active fraction of bed sediments and suspended matter, respectively. ρ_s is the sediment bulk density.

The temporal evolution of pollutant concentration in suspended particles is given by

$$\begin{aligned} \frac{\partial(HC_s m)}{\partial t} + \frac{\partial(uHC_s m)}{\partial x} + \frac{\partial(vHC_s m)}{\partial y} = \frac{\partial}{\partial x} \left(HK_h \frac{\partial(C_s m)}{\partial x} \right) \\ + \frac{\partial}{\partial y} \left(HK_h \frac{\partial(C_s m)}{\partial y} \right) \\ + k_1^{spm} C_d H - k_2 m C_s H + SED, \end{aligned} \quad (14)$$

where SED expresses the pollutant exchange between suspended particles and the bed sediment resulting from erosion/deposition:

$$SED = \begin{cases} -SRC_s, & SR > 0 \\ -SRA_s, & SR < 0. \end{cases} \quad (15)$$

The equation for the temporal evolution of concentration in the bed sediment is

$$\frac{\partial A_s}{\partial t} = k_1^{sed} \frac{C_d H}{L \rho_s f_p} - k_2 A_s \phi + SED, \quad (16)$$

where now the exchange due to erosion/deposition of suspended particles is written as

$$SED = \begin{cases} \frac{SRC_s}{L \rho_s f_p}, & SR > 0 \\ \frac{SRA_s}{L \rho_s f_p}, & SR < 0. \end{cases} \quad (17)$$

The total concentration of pollutants in the sediment, A_{tot} , is computed from

$$A_{tot} = f_p A_s. \quad (18)$$

2.4. Numerical solution

All equations are solved using explicit finite difference schemes (Kowalik and Murty, 1993) on a grid with resolution of approximately 200 m. Second order accuracy schemes are used for advective and diffusive terms. Time step to solve the hydrodynamic equations, limited by the CFL condition (Kowalik and Murty, 1993), is $\Delta t = 5$ s. Since the sediment transport and pollutant dispersion models are run off-line, time step in these calculations was increased to 30 s (stability conditions for hydrodynamic calculations are more restrictive than for the transport ones). The model domain may be seen in Fig. 1, where locations of interest are also shown.

Open boundary conditions consist of, in hydrodynamic calculations, specifying sea surface elevations compiled from observations (Alvarez, 1999). In transport calculations, zero gradient is assumed across the open boundary.

As commented before, the hydrodynamic model is run off-line. Once it has been calibrated and validated through comparisons of calculated tidal constants with observations, these stored tidal constants are applied to simulate sediment processes. Once realistic

results on sediment transport are obtained, the pollutant transport model may be simulated. The pollutant transport model is run coupled with the sediment transport one, since instantaneous values of suspended matter concentration and sedimentation rates are required in each point of the domain by the pollutant model.

It is worth commenting that the simple addition of tidal and residual currents for different constituents implies that non-linear interactions between flow components are neglected. However, this is a common practice in pollutant transport modelling (see for instance Proctor et al., 1994).

3. Results

Results from the hydrodynamic, sediment transport and pollutant dispersion models are presented separately in the following subsections.

3.1. Tide simulations

The model has been run for the M_2 , S_2 , N_2 , M_1 , O_1 and M_4 constituents separately. A calibration procedure has been carried out with the M_2 constituent to select the optimum values of the bed friction coefficient and of horizontal eddy viscosity. Once selected, the same values are used for the other constituents. They are $k=0.007$ and $A=10 \text{ m}^2/\text{s}$. As usual, the hydrodynamic model is started from rest and run until stable oscillations are obtained. Then tidal analysis is carried out.

As an example, maps showing the tide amplitude and current amplitude for the main constituent, M_2 , are presented in Fig. 2. Residual currents for this constituent are shown in Fig. 3. These maps are in agreement with those presented in Alvarez (1999). It may be seen that the tide amplitude increases by a few cm from the bay connection with the Atlantic towards the inner Bay. Maximum currents, as should be expected, are found in the narrow channel connecting the outer and inner bays. Tide residual

circulation presents a number of gyres, in agreement with those previously found (Alvarez, 1999).

A quantitative comparison between observed and calculated tidal constants for the considered semidiurnal constituents is presented in Table 1, and for the diurnal ones in Table 2. Finally, results for the M_4 constituent are given in Table 3. Generally speaking, there is a good agreement between the observed and calculated tidal constants for all considered constituents.

3.2. Sediment dynamics

The sediment transport model is started from a uniform suspended matter concentration over all the domain and, as in the case of tides, run until stable oscillations in suspended matter are obtained. One of the difficulties of modelling sediment transport is defining parameters appearing in the model. In the present application, the deposition and erosion threshold stresses are fixed, respectively, as 0.2 and 0.35 N/m^2 , and the erodability constant as $1.1 \times 10^{-4} \text{ kg}/\text{m}^2 \text{ s}$. They are within typical ranges found in literature (Tattersall et al., 2003). The characteristic diameter of suspended particles has been defined as $4 \mu\text{m}$. This value has been selected after model calibration, and is of the same order as the value used in a similar modelling study for the Albor n Sea (Perri nez, 2008), which was $8 \mu\text{m}$. A standard value of $2600 \text{ kg}/\text{m}^3$ has been used for particle density ρ . The effective diffusion coefficient has been set as $15 \text{ m}^2/\text{s}$. With this set of parameters, results in agreement with available data are obtained. Three suspended sediment sources are included: San Pedro River, Guadalete River and Santi Petri Channel (see Fig. 1). As stream flows are extremely low, input of particles is produced only during the ebb tide. It is known that sediments in the inner bay are predominantly muds, since this area is protected from winds and waves (Ligero et al., 2005). On the other hand, sediments are predominantly sandy in areas with strong tidal currents, as the entrance of the bay and the channel connecting the inner and outer bays (Araujo et al., 2009; Carrasco et al., 2003). Thus, in a schematic way, parameter f_p (which gives the fraction of muddy

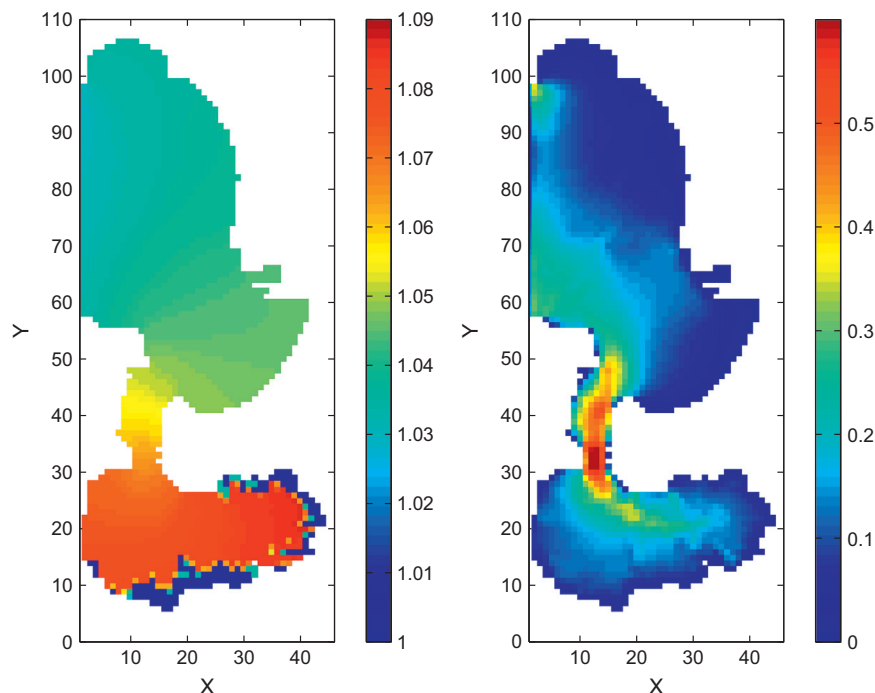


Fig. 2. Calculated M_2 amplitude (m, left) and current amplitude (m/s, right). Blues in the inner bay (left panel) correspond to intertidal areas. (For interpretation of the references to color in this figure caption, the reader is referred to the web version of this article.)

sediments) has been fixed as 0.2 in the channel, 0.45 in the outer bay and 1.0 in the inner bay. As will be shown, this schematic representation of the seabed composition is enough to obtain realistic results.

As an example, suspended particle concentrations at four instants during a tidal cycle (high water, ebb, low water and flood) are presented in Fig. 4. Minimum concentrations are observed

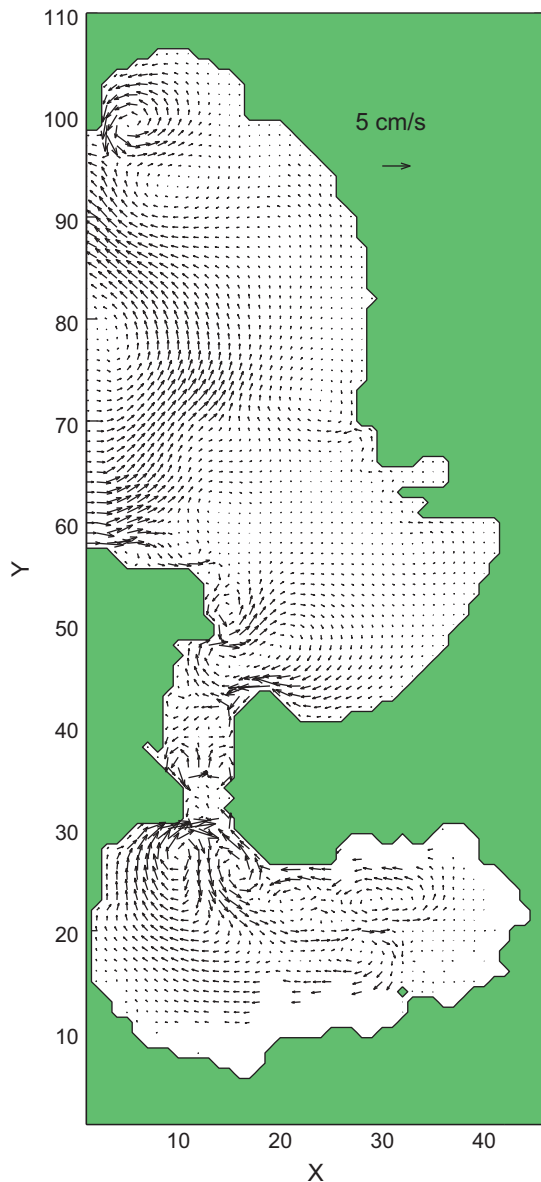


Fig. 3. Calculated M_2 residual currents.

during high water, since there are not riverine inputs and already present particles are diluted in a larger water volume. During the ebb and at low water, inputs from river are apparent. Finally, during the flood period the outer bay is relatively clear due to the input of water from the sea, while higher suspended matter concentrations are observed in the inner bay. Indeed, the flood tide will erode sediments of the channel connecting the inner and outer bays, and these resuspended sediments will be transported towards the inner bay, where they are concentrated.

Suspended matter concentrations have been compared with measurements in 4 points of the domain (see Fig. 1). Sampling started 1.5 h before high water and ended 1.5 h after high water, but the exact sampling time is not known for each sample. Thus, we have compared the range of suspended matter concentrations calculated in such time frame with measurements. Results are presented in Table 4. It may be seen that computed values are reasonable, similar to measured concentrations.

Sedimentation rates have also been compared with those obtained from measurements in sediment cores (Ligero et al., 2002). Location of sampling points is indicated in the map in Fig. 1 and a comparison between measurements and model results is presented in Table 5. Calculated results in Table 5 have been obtained averaging the instantaneous sedimentation rates provided by the model (each 30 s) over a period of 60 days. The model produces sedimentation rates of the order of 0.20 cm/y, similar to measured values.

Table 2

Observed, subscript *obs*, and computed, subscript *comp*, amplitudes (A , m) and phases (g , deg) of diurnal tidal elevations at several locations indicated in Fig. 1.

Station	K_1				O_1			
	A_{obs}	g_{obs}	A_{comp}	g_{comp}	A_{obs}	g_{obs}	A_{comp}	g_{comp}
Carraca	–	–	0.062	43	0.055	302	0.051	301
Pto Real	–	–	0.062	43	0.055	302	0.051	301
P. Carranza	0.063	46	0.062	43	0.053	303	0.051	301
Pto C�diz	0.067	41	0.061	42	0.052	302	0.050	300
Pto Sherry	0.064	42	0.061	42	0.054	300	0.050	300
Rota	0.064	38	0.062	42	0.058	302	0.051	301

Table 3

Observed, subscript *obs*, and computed, subscript *comp*, amplitudes (A , m) and phases (g , deg) of M_4 tidal elevations at several locations indicated in Fig. 1.

Station	A_{obs}	g_{obs}	A_{comp}	g_{comp}
Carraca	0.028	302	0.024	324
Pto Real	0.026	309	0.024	324
P. Carranza	0.028	309	0.023	323
Pto C�diz	0.021	315	0.020	319
Pto Sherry	0.020	308	0.020	320
Rota	0.018	312	0.019	320

Table 1

Observed, subscript *obs*, and computed, subscript *comp*, amplitudes (A , m) and phases (g , deg) of semidiurnal tidal elevations at several locations indicated in Fig. 1.

Station	M_2				S_2				N_2			
	A_{obs}	g_{obs}	A_{comp}	g_{comp}	A_{obs}	g_{obs}	A_{comp}	g_{comp}	A_{obs}	g_{obs}	A_{comp}	g_{comp}
Carraca	1.08	60	1.08	60	0.35	86	0.38	85	0.27	41	0.28	35
Pto Real	1.07	58	1.08	60	0.39	82	0.37	85	0.29	43	0.28	35
P. Carranza	1.07	58	1.07	58	0.36	87	0.37	84	0.27	38	0.27	35
Pto C�diz	1.03	55	1.04	54	0.36	84	0.36	82	0.26	37	0.26	33
Pto Sherry	1.03	53	1.04	54	0.35	81	0.36	82	0.26	33	0.26	33
Rota	1.01	56	1.04	54	0.33	81	0.35	82	0.25	48	0.26	34

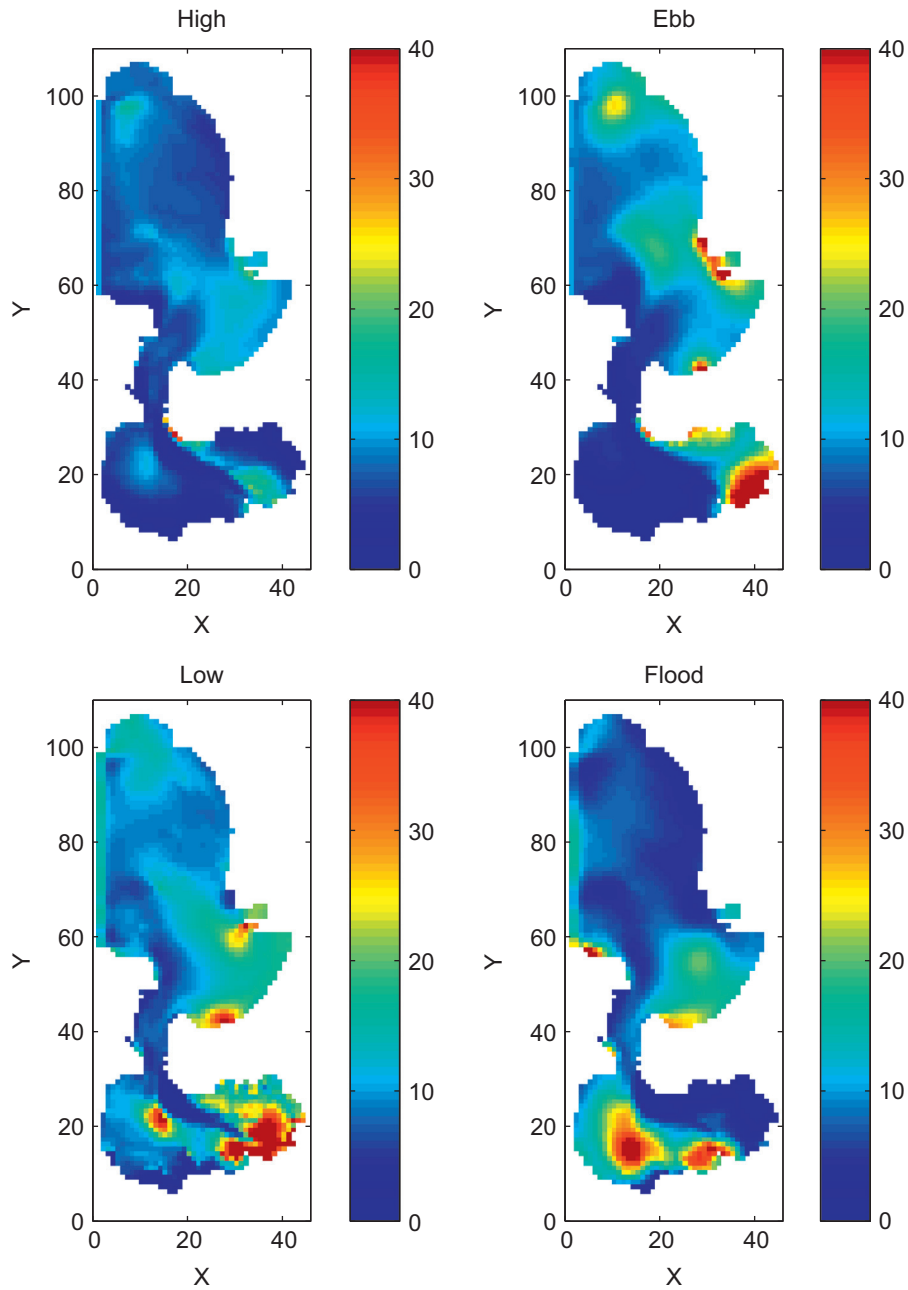


Fig. 4. Suspended matter concentration (mg/L) maps at 4 tidal states.

Table 4

Measured spm concentrations 1.5 h around high water and calculated range of values in the same time frame.

Station	Measured spm (mg/L)	Calculated range (mg/L)
BI4	7.4	1.0–5.6
BI1	8.0	1.7–12.3
BI6	12.2	6.0–11.8
BE4	16.0	9.0–13.4

Table 5

Measured and computed sedimentation rates in three points of the domain.

Station	Measured (cm/y)	Calculated (cm/y)
S1	0.16	0.12
S2	0.22	0.20
S3	0.27	0.24

Once it is known that the model produces realistic suspended matter concentrations and sedimentation rates, it may be applied to carry out particle reactive pollutant dispersion experiments.

3.3. Pollutant dispersion

3.3.1. Dispersion experiments

Some additional parameters have to be defined for the pollutant transport model. Following previous work (Periañez, 2000,

2003, 2009), $L=0.1$ m and $\phi=0.1$. Also, a constant porosity has been assumed for sediments, fixed as 0.6 from Ligeró et al. (2005).

Kinetic rates must be defined for each simulated pollutant. As widely described before (Periañez, 2005a, 2008, 2009), the same value may be used for k_2 for different elements since the main parameter controlling pollutant geochemical behavior is χ . Following the procedure described in such references, the exchange velocity may be deduced from k_2 and the pollutant equilibrium distribution coefficient k_d since the following relation holds (see the same references):

$$k_d = \frac{\chi}{k_2} \frac{3}{\rho R}. \quad (19)$$

The value $k_2 = 1.16 \times 10^{-5} \text{ s}^{-1}$ has been fixed. It has been successfully used in earlier simulations in the Odiel-Tinto estuary (Periañez et al., 2005) for Ra, in the Strait of Gibraltar-Alborán Sea (Periañez, 2008) for Cs and Pu, and in the Gulf of Cádiz (Periañez, 2009) for heavy metals. It was measured for Cs by Nyffeler et al. (1984).

A first numerical experiment has been carried out to simulate the adsorption of fallout ^{137}Cs (no other anthropogenic sources exist) on seabed sediments. A k_d of $2.0 \text{ m}^3/\text{kg}$ has been used since it agrees with the established value by IAEA (2004) for coastal waters and also was successfully applied in Periañez (2008). Since fallout has been occurring since several decades and with varying intensity, we have just made a qualitative comparison of the calculated ^{137}Cs inventories (total amounts) in the sediment mixed layer with the measured inventories in Ligeró et al. (2005). Thus, we have made a 3 month simulation with a constant and arbitrary atmospheric flux of ^{137}Cs . The calculated inventories in each grid cell have been normalized in such a way that the maximum calculated inventory is the same as the maximum measured one. Then, both spatial distributions are compared.

Results are presented in Fig. 5. Although measurements are available only in the inner bay, measured and calculated distributions are similar, indicating a significant accumulation of caesium in the zone of Carraca station (see Fig. 1). Model results are normalized, thus the significant aspect of this test is that the model predicts the accumulation of a pollutant introduced via atmospheric deposition in a given area of the bay, and this is in agreement with observations. Nevertheless, a second, but weaker, accumulation area in the western inner bay is not reproduced by the model. Calculated levels of ^{137}Cs in the outer bay are lower, due to the higher water exchange with the open sea and subsequent sediment cleaning.

As commented before, a number of pollutant sources exist in Cádiz Bay (industrial and urban wastewater discharges), whose magnitudes are not known. Thus, it makes no sense trying to compare calculated metal distributions with measurements. This would imply setting the magnitude of each source, by trial and error, until results in agreement with observations are obtained. Consequently, we have just carried out some simple numerical experiments to study the pollutant distributions resulting from some sources.

An experiment has been carried out with Zn, which is a much more reactive element than Cs. Its distribution coefficient is (IAEA, 2004) $70 \text{ m}^3/\text{kg}$, which has been used in previous simulations of the Gulf of Cádiz (Periañez, 2009). Carrasco et al. (2003) have found high Zn contents in waters of both the inner and outer bays, without significant differences between tidal conditions. These Zn concentrations were particularly high in the channel connecting the inner and outer bays, in the area of Guadalete River (see Fig. 6 for its location) and in Pto. Cádiz site (Fig. 1). A numerical experiment has been carried out introducing a constant Zn release, of arbitrarily defined magnitude, from the Guadalete River. Calculated Zn concentrations in water, suspended matter and bed sediments after 60 days of release are shown in Fig. 6. It seems

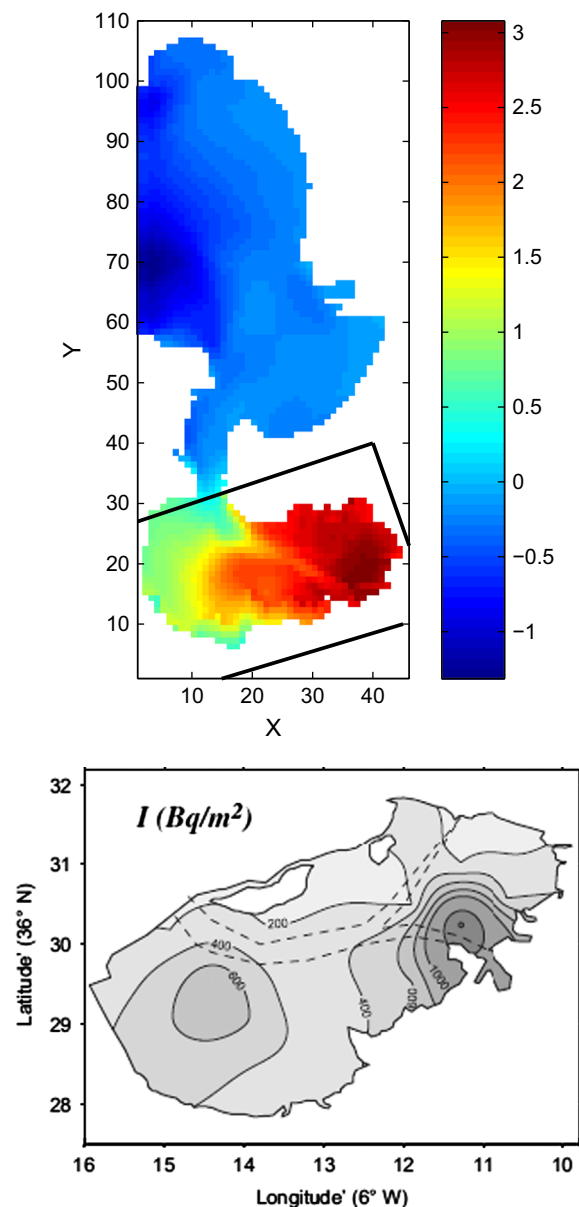


Fig. 5. Calculated (top, normalized logarithms) and measured (bottom, Ligeró et al., 2005) ^{137}Cs inventories in bed sediments (Bq/m^2). Note that the model grid is rotated 30.8° clockwise. The black rectangle on the top panel indicates the region where measurements indicated in the lower panel were carried out.

clear that contaminants are pumped into the inner bay, where they remain trapped, confirming the existence of an accumulation area in the zone of Carraca station, as already found with the Cs experiment. This accumulation area must be related to residual circulation in the inner bay, where several gyres (see Fig. 3) retain substances making them to recirculate. Thus any pollutant release, even if it occurs in the outer bay, may have a significant environmental impact on the inner bay and has to be planned with extreme care. In this sense, it is relevant to investigate flushing times of pollutants for the inner bay, which are analyzed in the following section.

Recently, an extensive sampling campaign has been carried out in Cádiz Bay (Casas-Ruiz et al., 2012), and ^{226}Ra concentrations have been measured in sediments. A map of the measured concentrations is presented in Fig. 7. A simulation has been carried out introducing a source of arbitrary magnitude from Guadalete River, where the maximum concentrations have been measured. The recommended k_d value for Ra by IAEA (2004) is $2 \text{ m}^3/\text{kg}$,

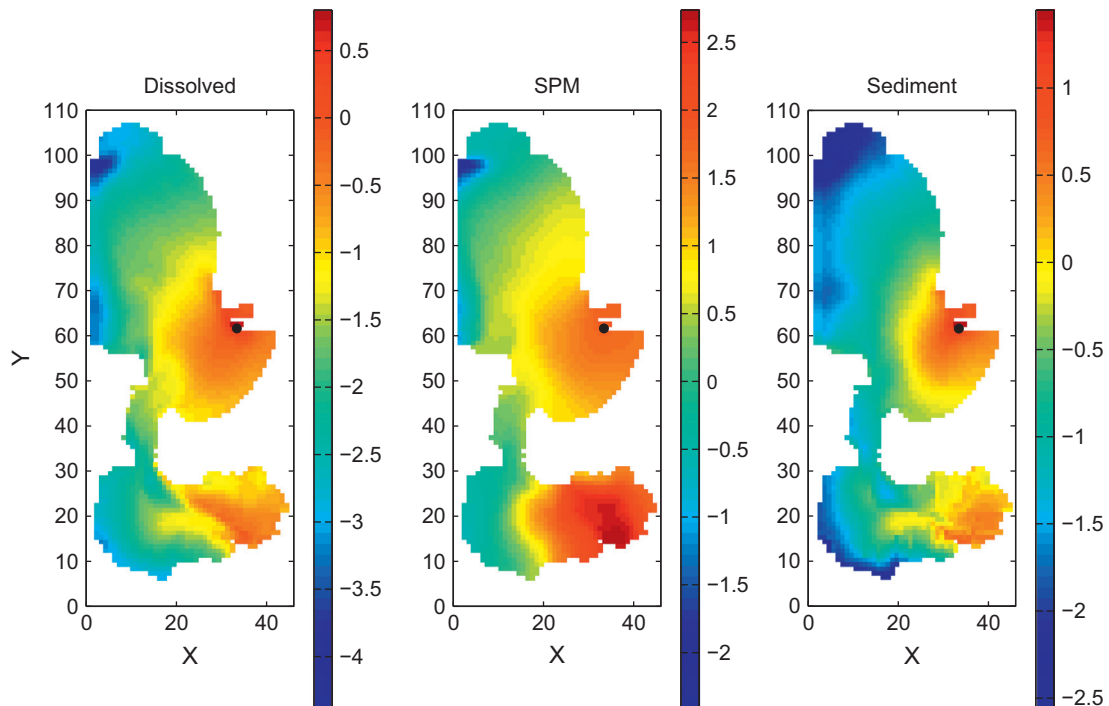


Fig. 6. Calculated Zn concentrations in water (units/m³), suspended matter and bed sediments (both in units/kg). The black dot indicates the Guadalete River mouth, where a constant release of arbitrary magnitude occurs. Logarithms of concentrations are drawn.

which has been used in our simulation. The computed ²²⁶Ra distribution in sediments is also presented in Fig. 7. Measured and computed distributions are somewhat similar, with lower concentrations in the channel connecting the inner and outer bays and higher values again in the inner bay. As in the case of the ¹³⁷Cs experiment, the calculated concentrations have been normalized. Even so, calculated concentrations cannot be compared with the measured ones (note for instance that calculated concentrations in the inner bay are lower than measured ones). This is due to the fact that ²²⁶Ra is a natural radionuclide, also having a geological (thus diffuse) origin (Casas-Ruiz et al., 2012). We are not including this diffuse source, but only a point source in Guadalete River. It is not surprising that calculated concentrations are lower than measured ones. Nevertheless, the overall calculated ²²⁶Ra distribution seems to be reasonable.

3.3.2. Flushing time of the inner bay

Flushing time is defined as the time in which the tracer inventory in the water column (total amount of tracer in the considered region) decreases in a factor e (Prandle, 1984). Flushing time is relevant for the water quality of a system, and it is important to know the time scale for a pollutant discharged into a water body to be transported out of the system (Shen and Haas, 2004). Thus, flushing times have been recently determined for a number of water bodies using numerical experiments (Shen and Haas, 2004; Choi and Lee, 2004; Patgaonkar et al., 2012; Periañez, 2012).

Flushing time has been determined considering a perfectly conservative (dissolved) contaminant. Model initial conditions consist of setting a 100 units/m³ concentration at every grid cell inside the inner bay and zero concentration elsewhere. The model is integrated and the time evolution of the pollutant inventory inside the bay written to an output file. Fitting to exponential decay curves provides the flushing time. Thus, the system-wide (as defined in Choi and Lee, 2004) flushing time has actually been

determined. These authors found that this parameter is useful for determining the long-term water quality of a system.

Choi and Lee (2004) have found that the system-wide flushing time is better obtained from a double exponential decay curve rather than a single exponential decay. The temporal evolution of the mass within the system, $M(t)$, is thus given by the following equation:

$$\frac{M(t)}{M_0} = (1 + \beta)e^{-\alpha_1 t} - \beta e^{-\alpha_2 t}, \quad (20)$$

where M_0 is the initial mass within the system. If the three parameters β , α_1 and α_2 are determined from numerical fitting, the system-wide flushing time is given by (Choi and Lee, 2004)

$$T_f = \frac{1 + \beta}{\alpha_1} - \frac{\beta}{\alpha_2}. \quad (21)$$

As an example, the temporal evolution of the normalized tracer inventory, M/M_0 in the inner bay may be seen in Fig. 8. Numerical fitting to the double exponential provides a flushing time of 63.1 days, with a determination coefficient for the fitting $r^2 = 0.9966$. This time is considerably longer than in Algeciras Bay, for instance (Periañez, 2012). This is not surprising since, although tides are stronger in Cádiz Bay than in Algeciras Bay, the inner Cádiz Bay is a much closed water body. Also, pollutants are trapped by residual circulation as commented above. Indeed, an additional numerical experiment has been carried out in which advective terms in the dispersion equation have been removed, i.e., like there is not any water circulation at all and the only transport process is diffusion. The computed flushing time is of the same order as that obtained with all tidal constituents, confirming the extremely low capacity of tides to flush pollutants off the inner bay.

4. Conclusions

A numerical model to simulate the dispersion of particle-reactive contaminants in Cádiz Bay (SW Spain) has been developed.

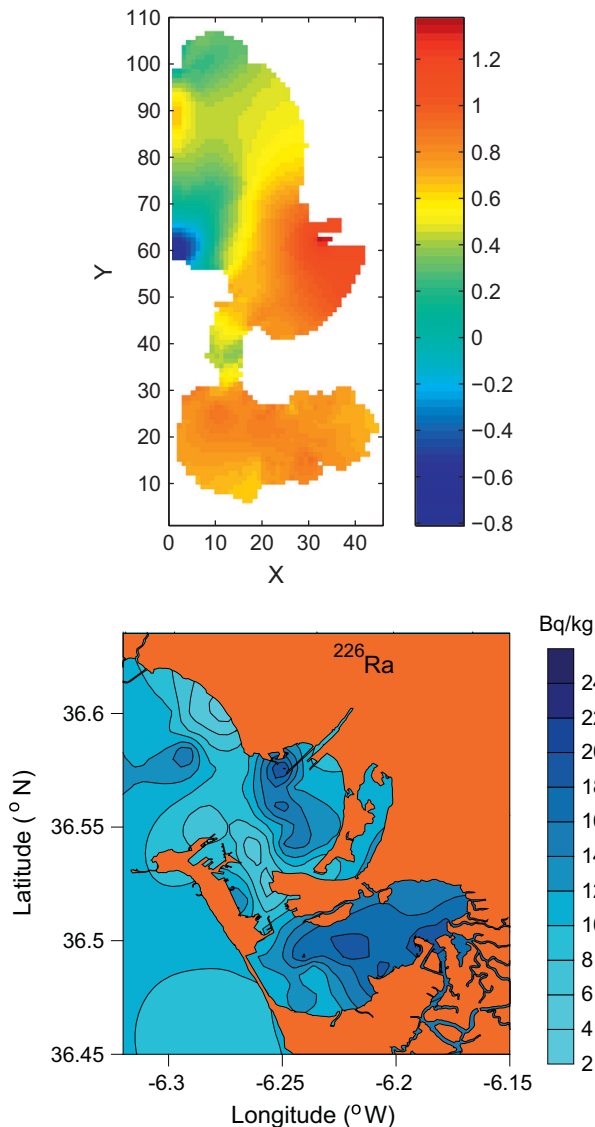


Fig. 7. Calculated (top) ^{226}Ra concentrations in sediments (Bq/kg), and measured (bottom) concentrations (Bq/kg). Logarithms of calculated concentrations are drawn.

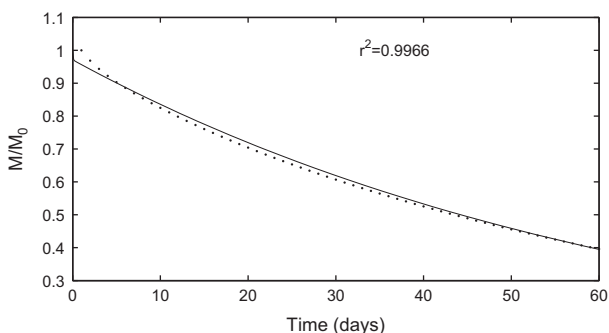


Fig. 8. Time evolution of the normalized inventory of a dissolved contaminant in the inner bay. Fitting to the double exponential in Eq. (20) is also shown (solid line).

It consists of a hydrodynamic model, a sediment transport model and a pollutant transport model. Hydrodynamics is calculated in advance, and the sediment and pollutant models are run off-line.

Calculated water circulation and suspended sediment dynamics is generally reproduced by the model. The contamination of seabed

sediments with fallout ^{137}Cs has been studied through a numerical experiment. It has been found that there is an accumulation area in the inner bay. This was already found in field measurements. Thus, the model is correctly describing transport processes inside the bay. If a continuous source of pollutants is present, even if it is in the outer bay, there is a pumping of the contaminant to the inner bay, where they are accumulated in the sediments of the area already found with the ^{137}Cs experiment. The existence of such accumulation area must be attributed to water circulation in the inner bay.

Finally, flushing time of the inner bay has been determined. A value of the order of 60 days has been obtained. This time is significantly longer than in other environments with a weaker tidal forcing (for instance Algeciras Bay), since this is a much closed environment. Also, residual circulation produces an accumulation area of pollutants in the inner bay, where they remain trapped. This increases the flushing time.

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